

# A formula for the number of Gelfand-Zetlin patterns

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## Abstract

In this article, we give a formula for the number of Gelfand-Zetlin patterns, using dimensions of the symmetry classes of tensors.

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## 1 Introduction

Suppose  $\Lambda = (a_1, a_2, \dots, a_{n-1})$  is a sequence of decreasing non-negative integers. A Gelfand-Zetlin pattern based on  $\Lambda$  is an array of integers:

$$\begin{array}{cccccc} a_1 & a_2 & a_3 & \cdots & a_{n-2} & a_{n-1} \\ & b_1 & b_2 & b_3 & \cdots & b_{n-2} \\ & & c_1 & c_2 & \cdots & c_{n-3} \\ & & & & \cdots & \end{array}$$

such that for all  $i$ ,

$$\begin{aligned} a_i &\geq b_i \geq a_{i+1} \\ b_i &\geq c_i \geq b_{i+1} \\ &\vdots \end{aligned}$$

We denote the set of all Gelfand-Zetlin patterns based on  $\Lambda$  by  $\Gamma_\Lambda$ . The set  $\Gamma_\Lambda$  has an important role in the representation theory of general (equivalently, special) and orthogonal linear Lie algebras. For example, let  $\Lambda$  be a dominant weight for the Lie algebra  $\mathfrak{sl}_n(\mathbb{C})$  and suppose  $L(\Lambda)$  is the corresponding irreducible representation with the highest weight  $\Lambda$ . In [5], Gelfand and Zetlin have proved that the set  $\Gamma_\Lambda$  can be viewed as a basis for  $L(\Lambda)$ . For the matrix representations of the elements of the Chevalley basis of  $\mathfrak{sl}_n(\mathbb{C})$  with respect to  $\Gamma_\Lambda$ , see [1] or [3]. The set  $\Gamma_\Lambda$  is also important from the point of view of *branching rules*. Branching rules are descriptions of the reduction of irreducible representations upon restriction to a subalgebra (subgroup). The first branching rule discovered is possibly the well-known branching rule of representations of the symmetric group  $S_m$  in the early twenties.

Since then, it was an exciting job to discover other kinds of branching rules for finite groups, Lie groups, and Lie algebras. We can employ the set  $\Gamma_\Lambda$  to describe the branching rule for type  $\mathfrak{sl}_n(\mathbb{C}) \rightarrow \mathfrak{sl}_r(\mathbb{C})$ . Suppose we like to restrict the representation  $L(\Lambda)$  to  $\mathfrak{sl}_r(\mathbb{C})$ . For any Gelfand-Zetlin pattern  $M \in \Gamma_\Lambda$ , let  $M^i$  be the  $i$ -th row of  $M$ . Then, the restriction of  $L(\Lambda)$  to  $\mathfrak{sl}_r(\mathbb{C})$  is equal to the following direct sum decomposition:

$$\bigoplus_{M \in \Gamma_\Lambda} \frac{1}{|\Gamma_{M^{n-r+1}}|} L(M^{n-r+1}).$$

The aim of this article is to compute the number of elements of  $\Gamma_\Lambda$ . Although, one can use the well-known dimension formula of Weyl, but our formula is an alternative one, which uses the irreducible characters of the symmetric group. To give a survey of our main result, suppose

$$m = a_1 + 2a_2 + 3a_3 + \cdots + (n-1)a_{n-1}.$$

We consider a partition  $\pi$  of  $m$  with the parts:

$$\pi_i = a_i + a_{i+1} + \cdots + a_{n-1}.$$

Let  $\chi_\pi$  be the irreducible character of the symmetric group  $S_m$  corresponding to  $\pi$ , (for standard terms about partitions and characters of  $S_m$ , see [11]). Also, for any permutation  $\sigma \in S_m$ , let  $c(\sigma)$  be the number of disjoint cycles in the cycle decomposition of  $\sigma$ . It is clear that the function  $\xi_n(\sigma) = n^{c(\sigma)}$  is a character of  $S_m$ , (for its irreducible constituents, see [12]). Our main result will be

$$|\Gamma_\Lambda| = [\chi_\pi, \xi_n],$$

where  $[\cdot, \cdot]$  is the inner product of characters in  $S_m$ . In the other words, we will see that

$$|\Gamma_\Lambda| = \frac{1}{m!} \sum_{\sigma \in S_m} \chi_\pi(\sigma) n^{c(\sigma)}.$$

## 2 Symmetry classes of tensors

In this section, we are going to review the notion of a symmetry class of tensors. The reader interested in the subject can find a detailed introduction in [9] or [10].

Let  $V$  be an  $n$ -dimensional complex inner product space and let  $G$  be a subgroup of the full symmetric group  $S_m$ . Let  $V^{\otimes m}$  denote the tensor product of  $m$  copies of  $V$  and for any  $\sigma \in G$ , define the *permutation operator*:

$$P_\sigma : V^{\otimes m} \longrightarrow V^{\otimes m}$$

by

$$P_\sigma(v_1 \otimes v_2 \otimes \cdots \otimes v_m) = v_{\sigma^{-1}(1)} \otimes v_{\sigma^{-1}(2)} \otimes \cdots \otimes v_{\sigma^{-1}(m)}.$$

Suppose that  $\chi$  is a complex irreducible character of  $G$  and define the *symmetrizer*:

$$S_\chi = \frac{\chi(1)}{|G|} \sum_{\sigma \in G} \chi(\sigma) P_\sigma.$$

The *symmetry class of tensors* associated with  $G$  and  $\chi$  is the image of  $S_\chi$  and it is denoted by  $V_\chi(G)$ . So

$$V_\chi(G) = S_\chi(V^{\otimes m}).$$

For example, if we let  $G = S_m$  and  $\chi = \varepsilon$ , the alternating character, then we get  $\wedge^m V$ , the  $m$ -th Grassman space over  $V$  and if  $G = S_m$  and  $\chi = 1$ , the principal character, then we obtain  $V^{(m)}$ , the  $m$ -th symmetric power of  $V$ , as symmetry classes of tensors.

Several monographs and articles have been published on symmetry classes of tensors during the last decades, see for example [9, 10].

Let  $v_1, \dots, v_m$  be arbitrary vectors in  $V$  and define the *decomposable symmetrized tensor*:

$$v_1 * v_2 * \dots * v_m = S_\chi(v_1 \otimes v_2 \otimes \dots \otimes v_m).$$

Let  $\{e_1, \dots, e_n\}$  be a basis of  $V$ , and suppose that  $\Gamma_n^m$  is the set of all  $m$ -tuples of integers  $\alpha = (\alpha_1, \dots, \alpha_m)$  with  $1 \leq \alpha_i \leq n$ . For  $\alpha = (\alpha_1, \dots, \alpha_m) \in \Gamma_n^m$ , we use the notation  $e_\alpha^*$  for decomposable symmetrized tensor  $e_{\alpha_1} * \dots * e_{\alpha_m}$ . It is clear that  $V_\chi(G)$  is generated by all  $e_\alpha^*$ ;  $\alpha \in \Gamma_n^m$ . We define an action of  $G$  on  $\Gamma_n^m$  by

$$\alpha^\sigma = (\alpha_{\sigma^{-1}(1)}, \dots, \alpha_{\sigma^{-1}(m)})$$

for any  $\sigma \in G$  and  $\alpha \in \Gamma_n^m$ . Given two elements  $\alpha, \beta \in \Gamma_n^m$ , we say that  $\alpha \sim \beta$  if and only if  $\alpha$  and  $\beta$  lie in the same orbit. Suppose that  $\Delta$  is a set of representatives of orbits of this action and let  $G_\alpha$  denote the stabilizer subgroup of  $\alpha$ . Define

$$\Omega = \{\alpha \in \Gamma_n^m : [\chi, 1_{G_\alpha}] \neq 0\},$$

where  $[\ , \ ]$  denotes the inner product of characters (see [7]). It is well known that  $e_\alpha^* \neq 0$ , if and only if  $\alpha \in \Omega$ , see for example [10]. Suppose  $\bar{\Delta} = \Delta \cap \Omega$ . For any  $\alpha \in \bar{\Delta}$ , we have the cyclic subspace:

$$V_\alpha^* = \langle e_{\alpha^\sigma}^* : \sigma \in G \rangle.$$

It is proved that we have the direct sum decomposition:

$$V_\chi(G) = \sum_{\alpha \in \bar{\Delta}} V_\alpha^*,$$

see [10] for a proof. It is also proved that

$$s_\alpha := \dim V_\alpha^* = \chi(1)[\chi, 1_{G_\alpha}],$$

and in particular, if  $\chi$  is linear then  $s_\alpha = 1$  and so the set:

$$\{e_\alpha^* : \alpha \in \bar{\Delta}\}$$

is an orthogonal basis of  $V_\chi(G)$ . Also in the case of linear character  $\chi$ , we have  $e_{\alpha^\sigma}^* = \chi(\sigma^{-1})e_\alpha^*$ . In the general case, let  $\alpha \in \bar{\Delta}$  and suppose

$$e_{\alpha^\sigma_1}^*, e_{\alpha^\sigma_2}^*, \dots, e_{\alpha^\sigma_t}^*$$

is a basis of  $V_\alpha^*$  with  $\sigma_1 = 1$ . Let

$$A_\alpha = \{\alpha^{\sigma_1}, \alpha^{\sigma_2}, \dots, \alpha^{\sigma_t}\}.$$

Then, we define  $\hat{\Delta} = \bigcup_{\alpha \in \bar{\Delta}} A_\alpha$ . It is clear that

$$\bar{\Delta} \subseteq \hat{\Delta} \subseteq \Omega,$$

and the set:

$$\{e_\alpha^* : \alpha \in \hat{\Delta}\}$$

is a basis of  $V_\chi(G)$ . Finally, we remember a formula for dimension of symmetry classes. We have

$$\dim V_\chi(G) = \frac{\chi(1)}{|G|} \sum_{\sigma \in G} \chi(\sigma) n^{c(\sigma)},$$

where  $c(\sigma)$  denotes the number of disjoint cycles (including cycles of length one) in cycle decomposition of  $\sigma$ .

### 3 Symmetry classes as $\mathfrak{sl}_n(\mathbb{C})$ -modules

In this section, we define a Lie module structure on  $V_\chi(G)$ , so let  $L$  be a complex Lie algebra and suppose that  $V$  is an  $L$ -module. For any  $x \in L$ , define

$$D(x) : V^{\otimes m} \longrightarrow V^{\otimes m}$$

by

$$D(x)(v_1 \otimes v_2 \otimes \cdots \otimes v_m) = \sum_{i=1}^m v_1 \otimes \cdots \otimes xv_i \otimes \cdots \otimes v_m.$$

We know that  $D(x)S_\chi = S_\chi D(x)$  and so  $V_\chi(G)$  is invariant under  $D(x)$ . Suppose

$$D^*(x) = D(x) \downarrow_{V_\chi(G)},$$

where the down arrow denotes restriction.

**Definition 3.1.** Define an action of Lie algebra  $L$  on  $V_\chi(G)$  by

$$\begin{aligned} x(v_1 * \cdots * v_m) &= D^*(x)(v_1 * \cdots * v_m) \\ &= \sum_{i=1}^m v_1 * \cdots * xv_i * \cdots * v_m. \end{aligned}$$

Then,  $V_\chi(G)$  becomes an  $L$ -module. In what follows, we will assume that  $L = \mathfrak{sl}_n(\mathbb{C})$  and  $V = \mathbb{C}^n$ , the standard module for  $L$ . Then,  $V_\chi(G)$  becomes an  $L$ -module. In [8], the irreducible constituents of  $V_\chi(G)$  are determined. In this section, we give a summary of result of [8]. We also assume that  $G = S_m$  and  $\chi = \chi_\pi$ , the irreducible character of  $S_m$  corresponding a partition  $\pi$ . For simplicity, we denote the symmetry class of tensors by

$V_\pi(S_m)$ . To describe the irreducible constituents of  $V_\pi(S_m)$ , it is necessary to introduce some notations.

A Cartan subalgebra for  $L$  is

$$H = \{ \text{diag}(h_1, h_2, \dots, h_n) : h_1 + h_2 + \dots + h_n = 0 \}.$$

For any  $1 \leq i \leq n$ , define a linear functional:

$$\mu_i : H \longrightarrow \mathbb{C}$$

by

$$\mu_i(h) = h_i,$$

where  $h = \text{diag}(h_1, h_2, \dots, h_n)$ , so we have

$$\mu_1 + \mu_2 + \dots + \mu_n = 0,$$

and hence  $\mu_1, \mu_2, \dots, \mu_{n-1}$  is a basis for  $H^*$ .

Now let  $\Lambda_1, \Lambda_2, \dots, \Lambda_{n-1}$  be the fundamental weights corresponding to  $H$ . It is easy to see that for any  $k$ :

$$\Lambda_k = \mu_1 + \mu_2 + \dots + \mu_k.$$

Let  $\alpha \in \Gamma_n^m$ . We define a composition of  $m$  by  $m(\alpha) = (m_1, m_2, \dots, m_n)$ , where  $m_i$  is the multiplicity of  $i$  in  $\alpha$ . Suppose

$$\mu_\alpha = \mu_{\alpha_1} + \mu_{\alpha_2} + \dots + \mu_{\alpha_m}.$$

So we have

$$\mu_\alpha = m_1\mu_1 + m_2\mu_2 + \dots + m_n\mu_n.$$

Also we can see that

$$\mu_\alpha = (m_1 - m_2)\Lambda_1 + (m_2 - m_3)\Lambda_2 + \dots + (m_{n-1} - m_n)\Lambda_{n-1}.$$

It is easy to prove that  $\mu_\alpha = \mu_\beta$ , if and only if  $m(\alpha) = m(\beta)$ . So, for any  $\alpha \in \Gamma_n^m$ , we introduce a partition  $M(\alpha)$ , which is just the multiplicity composition  $m(\alpha)$  with a descending arrangement of entries. In fact, any partition of  $m$ , with height at most  $n$ , is of the form  $M(\alpha)$ , where  $\alpha \in \hat{\Delta}$ . For any  $h \in H$  and  $\alpha \in \hat{\Delta}$ , we have

$$h \cdot e_\alpha^* = \left( \sum_{i=1}^m \mu_{\alpha_i} \right) (h) e_\alpha^*.$$

In other words, we have

$$h \cdot e_\alpha^* = \mu_\alpha(h) e_\alpha^*.$$

So the set of weights of  $V_\pi(S_m)$  is  $\{\mu_\alpha : \alpha \in \hat{\Delta}\}$ . We also can see by an easy argument that the weight  $\mu_\alpha$  is dominant, if and only if  $M(\alpha) = m(\alpha)$ , i.e.  $m(\alpha)$  is a partition. For two dominant weights  $\mu_\alpha$  and  $\mu_\beta$ , it is routine to check that  $\mu_\beta$  appears in  $L(\mu_\alpha)$  as a weight, if and only if  $m(\alpha)$  majorizes  $m(\beta)$ , (for definition of the majorization, see [11]). We are now ready to compute irreducible constituents of  $V_\pi(S_m)$ . Although, the following theorem is proved in [8] in a more general framework, we prove it here again because what we need is only this special case.

**Theorem 3.2.** *We have*

$$V_\pi(S_m) = L(\mu_\alpha)^{\chi_\pi(1)},$$

where  $\alpha \in \bar{\Delta}$  is any element with the property  $M(\alpha) = \pi$ .

**Proof.** First of all, note that

$$\bar{\Delta} = \{\alpha \in G_n^m : M(\alpha) \preceq \pi\},$$

where  $G_n^m$  denotes the set of all increasing sequences in  $\Gamma_n^m$  and  $\preceq$  denotes majorization. Suppose that  $\eta_1, \dots, \eta_s$  is the set of all dominant weights of  $V_\pi(S_m)$ , ordered in such a way that  $\eta_j \preceq \eta_i$  implies  $i \leq j$  (we say  $\eta_j \preceq \eta_i$  iff  $\eta_i - \eta_j$  is a sum of positive roots). As we saw above, for any  $i$ , there is  $\alpha \in \bar{\Delta}$  such that  $\eta_i = \mu_\alpha$ . Suppose that  $r_i$  equals number of  $\beta \in \hat{\Delta}$  such that  $m(\alpha) = m(\beta)$ . Let  $m_{ij}$  be the multiplicity of  $\eta_i$  in  $\eta_j$ . Define a sequence of integers as follows:

$$\begin{aligned} c_1 &= r_1, \\ c_i &= r_i - \sum_{j=1}^{i-1} m_{ij} c_j. \end{aligned}$$

Now we have

$$V_\pi(S_m) = \sum_i L(\eta_i)^{c_i},$$

so we must show that  $c_1 = \chi_\pi(1)$  and  $c_i = 0$  for  $i \neq 1$ . First, note that if  $\alpha \in \bar{\Delta}$  has the property  $\mu_\alpha = \eta_1$ , then  $m(\alpha) = \pi$  and hence we have

$$\begin{aligned} c_1 &= r_1 \\ &= |\{\beta \in \hat{\Delta} : m(\alpha) = m(\beta)\}| \\ &= |\{\beta \in \hat{\Delta} : \alpha \sim \beta\}| \\ &= s_\alpha \\ &= \chi_\pi(1) [1_{(S_m)_\alpha}, \chi_\pi] \\ &= \chi_\pi K_{\pi, m(\alpha)} \\ &= \chi_\pi K_{\pi, \pi} \\ &= \chi_\pi(1). \end{aligned}$$

Note that  $K_{*,*}$  denote the well-known Kostka numbers, see [11] or [4] for definition. Now, we compute  $c_2$ ; we have  $c_2 = r_2 - m_{2,1}c_1$ . As in the above case, we easily see that  $r_2 = \chi_\pi(1)K_{\pi, m(\alpha')}$ , where  $\alpha'$  corresponds to  $\eta_2$ . It is proved that (see [4])

$$K_{\pi, m(\alpha')} = m_{2,1}.$$

Hence,

$$\begin{aligned} c_2 &= \chi_\pi K_{\pi, m(\alpha')} - K_{\pi, m(\alpha')} \chi_\pi(1) \\ &= 0. \end{aligned}$$

By a similar argument, we see that  $c_i = 0$  for other “ $i$ ”s and the result follows.  $\square$

Note that this theorem affords a new method of constructing of all irreducible  $\mathfrak{sl}_n(\mathbb{C})$ -modules, namely, let  $\Lambda = \mu_\pi$  be any integral dominant weight of  $\mathfrak{sl}_n(\mathbb{C})$ . As in [10], we have

$$V_\pi(S_m) = \sum_{i=1}^{\chi_\pi(1)} V_\pi^i(S_m),$$

where  $V_\pi^i(S_m)$  is defined as follows: let

$$F : S_m \longrightarrow GL_{\chi_\pi(1)}(\mathbb{C})$$

be the corresponding representation of  $\chi_\pi$  with  $F(\sigma) = [a_{ij}(\sigma)]$ . We introduce the partial symmetrizer  $S_\pi^i$  by

$$S_\pi^i = \frac{\chi_\pi(1)}{m!} \sum_{\sigma \in S_m} a_{ii}(\sigma) P_\sigma.$$

Now,  $V_\pi^i(S_m)$  is precisely the image of  $S_\pi^i$ . So we have

$$L(\Lambda) = V_\pi^i(S_m)$$

for all  $1 \leq i \leq \chi_\pi(1)$ . One of the most important consequences of this construction is the following dimension formula, which is different from Weyl's one.

**Corollary 3.3.** *Let  $L(\Lambda)$  be an irreducible  $\mathfrak{sl}_n(\mathbb{C})$ -module with the highest weight  $\Lambda$  and define the corresponding number  $m$  and the partition  $\pi$  as in Section 1. Then,*

$$\dim L(\Lambda) = \frac{1}{m!} \sum_{\sigma \in S_m} \chi_\pi(\sigma) n^{c(\sigma)}.$$

## 4 The number of the Gelfand-Zetlin patterns

We are ready now to give the interesting relation between dimension of the symmetry classes of tensors  $V_\pi(S_m)$  and the number of Gelfand-Zetlin patterns based on  $\Lambda$ . Note that we have the following relations between the weight  $\Lambda$  and the partition  $\pi$ :

$$m = a_1 + 2a_2 + 3a_3 + \cdots + (n-1)a_{n-1},$$

$$\pi_i = a_i + a_{i+1} + \cdots + a_{n-1},$$

$$a_i = \pi_i - \pi_{i+1}.$$

**Main theorem.** *The number of Gelfand-Zetlin patterns based on  $\Lambda$  is equal to the inner product  $[\chi_\pi, \xi_n]$ . Equivalently, we have*

$$\dim V_\pi(S_m) = \chi_\pi(1) |\Gamma_\Lambda|.$$

**Remark 4.1.** In [12], the inner product  $[\chi_\pi, \xi_n]$  is expressed in term of Kostka numbers. Suppose that  $\rho = [\rho_1, \dots, \rho_s]$  is a partition of  $m$  with distinct parts  $b_1, \dots, b_l$ . Suppose that the multiplicity of  $b_i$  in  $\rho$  is  $r_i$ . If  $s \leq n$ , define

$$f(n, \rho) = \frac{n!}{(n-s)! r_1! r_2! \cdots r_l!},$$

and let  $f(n, \rho) = 0$  for  $s > n$ . Then the multiplicity of  $\chi_\pi$  in  $\xi_n$  is equal to

$$\sum_{\rho} f(n, \rho) K_{\pi, \rho},$$

where  $K_{\pi, \rho}$  is the Kostka number.

**Remark 4.2.** Note that we can *normalize*  $\Lambda$  in such a way that we have  $a_{n-1} = 1$ . To do this, let  $d = a_{n-1} - 1$ . Define

$$\Lambda^* = (a_1 - d, a_2 - d, \dots, a_{n-2} - d, 1).$$

Although in general  $L(\Lambda)$  and  $L(\Lambda^*)$  are non-isomorphic representations, it is clear that  $|\Gamma_\Lambda| = |\Gamma_{\Lambda^*}|$ . For  $\Lambda^*$ , we have

$$m^* = m - \frac{n(n-1)}{2}(a_{n-1} - 1),$$

and also the corresponding partition  $\pi^*$  has the parts:

$$\pi_i^* = \pi_i - (n-1)(a_{n-1} + 1).$$

Hence, we have also the following normalized formula for the number of Gelfand-Zetlin patterns:

$$|\Gamma_\Lambda| = [\chi_{\pi^*}, \xi_n^*].$$

## References

- [1] A. O. Barut and R. Raczka. *Theory of Group Representations and Applications*. 2nd ed. PWN-Polish Scientific Publishers, Warszawa, 1980.
- [2] R. Carter. *Lie Algebras of Finite and Affine Type*. Cambridge Studies in Advanced Mathematics, **96**, Cambridge University Press, Cambridge, 2005.
- [3] L. Frappat, A. Sciarrino and P. Sorba. *Dictionary on Lie Algebras and Superalgebras*. Academic Press, San Diego, CA, 2000.
- [4] W. Fulton and J. Harris. *Representation Theory. A First Course*. Graduate Texts in Mathematics, **129**, Springer-Verlag, New York, 1991.
- [5] I. M. Gelfand and M. L. Zetlin. Finite dimensional representations of a group of unimodular matrices. *Dukl. Akad. Nauk. SSSR*, **71** (1950), 825–828.
- [6] J. Humphreys. *Introduction to Lie Algebras and Representation Theory*. Graduate Texts in Mathematics, **9**, Springer-Verlag, New York, 1972.
- [7] M. Isaacs. *Character Theory of Finite Groups*. Academic Press, New York, 1976.
- [8] A. R. Madadi and M. Shahryari. Symmetry classes of tensors as  $\mathfrak{sl}_n(\mathbb{C})$ -modules. *Linear and Multilinear Algebra*, **56** (2008), 517–541.
- [9] M. Marcus. *Finite Dimensional Multilinear Algebra, Part I*. Pure and Applied Mathematics, **23**, Marcel Dekker, New York, 1973.
- [10] R. Merris. *Multilinear Algebra*. Algebra, Logic and Applications, **8**, Gordon and Breach Science Publishers, Amsterdam, 1997.
- [11] B. Sagan. *The Symmetric Group. Representation, Combinatorial Algorithms, and Symmetric Functions*. The Wadsworth & Brooks/Cole Mathematics Series, Wadsworth & Brooks/Cole Advanced Books & Software, Pacific Grove, CA, 1991.
- [12] M. Shahryari and M. A. Shahabi. On a permutation character of  $S_m$ . *Linear and Multilinear Algebra*, **44** (1998), 45–52.

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