

Evolution of Self-organization and gelation in Laponite Ionogels

Nidhi Joshi¹ and Kamla Rawat^{2,3*}

¹Polymer and Biophysics Laboratory, School of Physical Sciences, JNU, New Delhi 110067, India

²Special Center for Nanosciences, JNU, New Delhi 110067, India

³Inter University Accelerator Centre (IUAC), New Delhi 110067, India

Abstract

We report formation of colloidal gel network in aqueous laponite dispersions containing ionic liquid (IL). With the passage of time, these flowing dispersions reorganized and attained a solid-like rigid structure. Structural reorganization of clay platelets in the IL environment caused reduction in the inter platelet repulsion due the IL-double layer screening and facilitated formation of stronger ionogels for samples with [IL] < 0.03% (w/v). Dynamic structure factor data revealed transition from ergodic to non-ergodic state is faster at higher IL concentration ([IL] > 0.03% (w/v)). The structure factor comprised two relaxation modes diffusion fast mode followed by an anomalous slow mode. The slow mode was frozen at all IL concentrations whereas the fast mode contained two diffusion coefficients both strongly dependent on IL concentration. In summary, laponite gels could be systematically customized in IL solutions to generate a range of soft materials with properties solely dependent on IL concentration which may not exceed 0.05% (w/v).

Keywords: Laponite gels; Ionic liquid; Dynamic light scattering; Nonergodicity; Self-assembly; Colloidal gels

Introduction

A rich phase diagram containing a hierarchy of structures manifested at different length and time scales is the signature of colloidal gel which is a unique state of soft matter [1-9]. Gelation and their arrested phases have been well studied in the literature in the past [1-13]. Multiple arrested phase states can be best followed, and hence probed using colloidal clay dispersions as model systems. The anisotropy associated with the geometrical shape and charge distribution in clays like laponite attributes unique properties to their dispersions. For instance, such systems contain both attractive and repulsive potentials in their interactions that allow either the system to form gel or to form a glassy phase. An interesting feature of clay suspension is the existence of interactions at different length and time scales which lead to a variety of structural and dynamical phenomena occurring as function of clay concentration and of the ionic strength. Recent studies performed on laponite dispersions have unveiled clear understanding of aggregation and glass formation which has further led to the unambiguous description of dynamical arrest and viscoelastic phase separation [9,10]. The gelation of dispersions of clay particles is both industrially important, as well as scientifically challenging [11,12]. In general, gels have found applications in drilling fluids, paints, varnishes and in cosmetic and pharmaceutical formulations. In addition, clay-polymer nanocomposites with customized properties could be successfully made [13]. Thus, it is imperative to consider new clay based soft matter systems that can widen the potential of their applications to newer domains.

Clays like laponite are discotic platelets with fixed aspect ratio (30, diameter \approx 30 nm thickness \approx 1 nm) and surface charge density (zeta potential \approx -35 mV), and their dispersions exhibit an array of soft matter phases that evolve with waiting time as well as solid concentration. Laponite is known to exhibit significant age dependent dynamics in aqueous dispersions [7-12]. Similarly, the salt dependent phase states of this clay have been reported extensively in the literature. In the recent times there has been a lot interest in room temperature ionic liquids. Ionic liquids (ILs) are salts composed of poorly coordinating organic cations and inorganic anions which are liquid below the boiling point of water. They possess appealing features such as low melting point, negligible vapor pressure, non-volatility, no-flammability, wide

electrochemical windows and so on [14-26]. ILs are often referred to as “designed solvents” because their physical properties could be tailored by adjusting the structures and species of the constituting cations and/or anions [19]. The variety of possible combinations of ILs with laponite can, in principle, generate a wide spectrum of designer soft materials. These can be used as filling material and matrices to tailor requirement based personal care, pharmaceutical and other industrial products.

Gels prepared in ionic liquid environment (ionogels) display unusual properties [27-36]. In ionic liquid (IL) solutions, one introduces mobile ions (in the form of mobile anions and comparatively heavy cations), hydrophobicity (alkyl chain length) and hydrogen bonding (N---H) possibility simultaneously into the continuous phase. These liquids, comprising of organic ions, are associated with negligible vapor pressure, high thermal, chemical and electrochemical stability that enable these to be treated as green solvents [37]. There exist several examples of ionogels depicting unusual properties. Preparation and microscopic dynamics of ionogels exhibiting two glass transition temperatures, which drastically decrease as the ionic liquid content increased due to increased plasticization of the polymer, was studied and has been reported [31,32].

In this work, we explore the gelation kinetics of laponite dispersions made in ionic liquid solutions in a systematic and comprehensive manner in a bid to establish their structure-property relationship. Further, the study was extended to probe the temporal relaxation dynamics of concentration fluctuations through the measurement of dynamic structure factor. The following protocol was used to achieve the above mentioned objective. The laponite dispersions were prepared and ionic liquid was added to this preparation to evaluate their

***Corresponding author:** Kamla Rawat, Special Center for Nanosciences, JNU, New Delhi 110067, India, Tel: +91 11 2670 4699, Fax: +91 11 2674 1837; E-mail: kamla.jnu@gmail.com

Received: November 10, 2014; **Accepted:** January 07, 2015; **Published:** January 09, 2015

Citation: Joshi N, Rawat K (2015) Evolution of self-organization and gelation in Laponite ionogels. J Phys Chem Biophys 5: 171. doi:10.4172/2161-0398.1000171

Copyright: © 2015 Joshi N, et al. This is an open-access article distributed under the terms of the Creative Commons Attribution License, which permits unrestricted use, distribution, and reproduction in any medium, provided the original author and source are credited.

influence on temporal effects of re-organisation self assembly, gelation kinetics and relaxation dynamics.

We have performed a systematic and comprehensive study to establish the time dependent formation of gel phase of laponite dispersions in ionic liquid (IL), 1-Ethyl-3-methylimidazolium chloride, solutions at room temperature. Self-assembly of IL coated laponite platelets preceded gelation. Considering the fact that this particular nanoclay, in its gel form, has found applications in pharmacy, food, and cosmetic industry etc. modifying its thermal and rheological properties is called for.

Materials and Methods

The IL, 1-ethyl-3-methyl imidazolium chloride, was bought from Sigma-Aldrich and was used as received. Deionized water from Organo Biotech Laboratories, India, was used to prepare the solutions. Laponite RD clay was purchased from Southern Clay Products, USA. We dried the laponite powder for 4 hrs at 120°C to remove its moisture content. IL solution was prepared by dissolving IL in double distilled deionized water for an hour. Laponite dispersions (2% w/v) were prepared by dispersing Laponite into double distilled deionized water and stir for 2 hours [19]. The desired IL concentration was achieved by adding a concentrated stock of IL solution. The pH of Laponite dispersion was adjusted to 10. All dispersions appeared optically clear and transparent. The concentration of laponite was maintained fixed at 2% (w/v) and the IL concentration was varied from [IL] = 0.01 to 0.05% (w/v). All preparations were made at room temperature 20° C and the relative humidity in the room was less than 50%. All the concentrations are in % (w/v) unless otherwise stated.

Dynamic light scattering (DLS) experiments were performed on a 256-channel digital correlator (Photocor Instruments, USA) that was operated in the multi-tau mode (logarithmically spaced channels). The time scale spanned over eight decades, i.e. from 0.5 μs to 10s. This instrument used a 35mW linearly polarized He-Ne laser. A Glan Thomson analyzer was used in front of the photomultiplier tube to enable collection of polarized components of scattered light. The scattered intensity values were measured for 20° < θ < 150°. The probe length scale is defined by the inverse of the modulus of the scattering wave vector q where the wave vector $q = (4\pi n/\lambda) \sin(\theta/2)$, the medium refractive index is n , excitation wavelength is λ ($\lambda = 654$ nm) and θ is scattering angle. Further details on dynamic light scattering can be obtained from ref [20].

Results and Discussion

Light scattering analysis

Dynamic light scattering (DLS) probing was used to study time dependent fluctuations of light scattered from laponite dispersion containing Ionic liquid (IL). In order to observe the effect of ionic liquid (IL) concentration on the aging dynamics of 2% laponite dispersion we performed dynamic light scattering measurements. Here, we measured the polarized (VV) intensity correlation functions independently at different scattering angles 20° < θ < 150° that allowed probing the sample over various length scales. The polarized autocorrelation function $g_2(q,t)$ gives information on the translational diffusion coefficient. The dynamic structure factor or electric field correlation function $g_1(q,t)$ could be described as sum of exponential term and Kohlrausch-Williams-Watts (KWW) stretched exponential function given as [38]

$$g_2(q,t) - 1 = A \exp\left(-\frac{t}{\tau_1}\right) + (1 - A) \exp\left(-\left(\frac{t}{\tau_2}\right)^\beta\right) \quad (1)$$

where the exponential term describes fast relaxation with characteristic time τ_1 whereas the stretched exponential term which is governed by characteristic time τ_2 and stretching parameter β describes slow relaxation. A and $(1-A)$ are the weights of the two contributions. For an ergodic scatterers, i.e. $0 \leq g_1(q,t) \leq 1$ and at $t \rightarrow \infty$, the correlation function relaxes to a value equal to zero (baseline). In non-ergodic systems, the scatterer is dynamically arrested and is unable to probe the entire phase space. As a consequence, the $g_1(q,t)$ at $t \rightarrow \infty$ cannot relax to its limiting value, but remains frozen at value greater than 0. The deviation of the baseline from the fully relaxed state is a measure of non ergodicity present in the system (defined as h_w).

Figure 1 illustrate the polarized correlation function $g_1(q,t)$ data for measurements pertaining to $\theta = 90^\circ$ for samples containing various amount of IL in the dispersion. The effect of aging is clearly visualised in this data. The observation can be clearly made from this data. The correlation function did not fully relax with aging time of the samples implying the emergence of non-ergodicity in the samples.

Figure 1 shows the waiting time evolution of polarized autocorrelation function at scattering angle 90° and indicate that all the system undergo transition from ergodic to non ergodic state. The system undergo transition from sol to gel state can be visually confirmed by titling the vials filled of samples and observing that the samples doesn't flow after a certain time and second through the decay of correlation function which doesn't undergo fully decay as system is undergoing transition from sol to gel state. The measurement times was varied from 9 hr to 18hr depending on the concentration so that accurate correlations functions are measured during the sol to gel transition process. For pure laponite, we can see that with increase in waiting time, $g_1(t)$ did not completely relax which indicate the confinement of laponite platelets in arrested phase. On introducing IL into the Laponite dispersion the system undergo earlier transition as compared to pure Laponite. On comparing the decay of correlation function of Laponite-IL dispersion we observethat the samples exhibited non-ergodicity much earlier above 0.03% as compared to low concentration. Thus, we conclude that with the increase in concentration of [IL] in laponite dispersion, the system developed non ergodic phase rather quickly as compared to pure laponite dispersion. Thus, introducing IL into the Laponite dispersion enhances the transition to non-ergodic state.

From the correlation function also we classified our samples into two gels. The decay of correlation function for higher IL concentration is earlier and also the difference is much less. The decay of correlation function varies from 1 to approximately 0.88 which led us to conclude that above 0.03% (w/v) IL concentration the gels formed are fragile. From time dependence behaviour of Laponite-IL we can say that above 0.03% (w/v) the gels formed are soft as compared to below this IL concentration.

We now have a look at the behavior of stretching exponent (β) which characterized the broadness of the distribution of slow relaxation mode contributing to the correlation function. In figure 2 the behavior of stretching parameter (β) which characterized the broadness of the distribution of slow mode contributing to the correlation function as the function of waiting time (t_w) is shown for samples with [IL] concentration 0% 0.01%, 0.03% and 0.05% (w/v) and it was found that the value of stretching parameter decreased with time from a value of ~0.7 to an extremely low value ~ 0.01 in the case of 0%, 0.01% and 0.03% (w/v) IL samples which implies that at higher waiting times the correlation functions become more and more stretched. However, decrease in case of 0.05% (w/v) is also seen but it is very small because

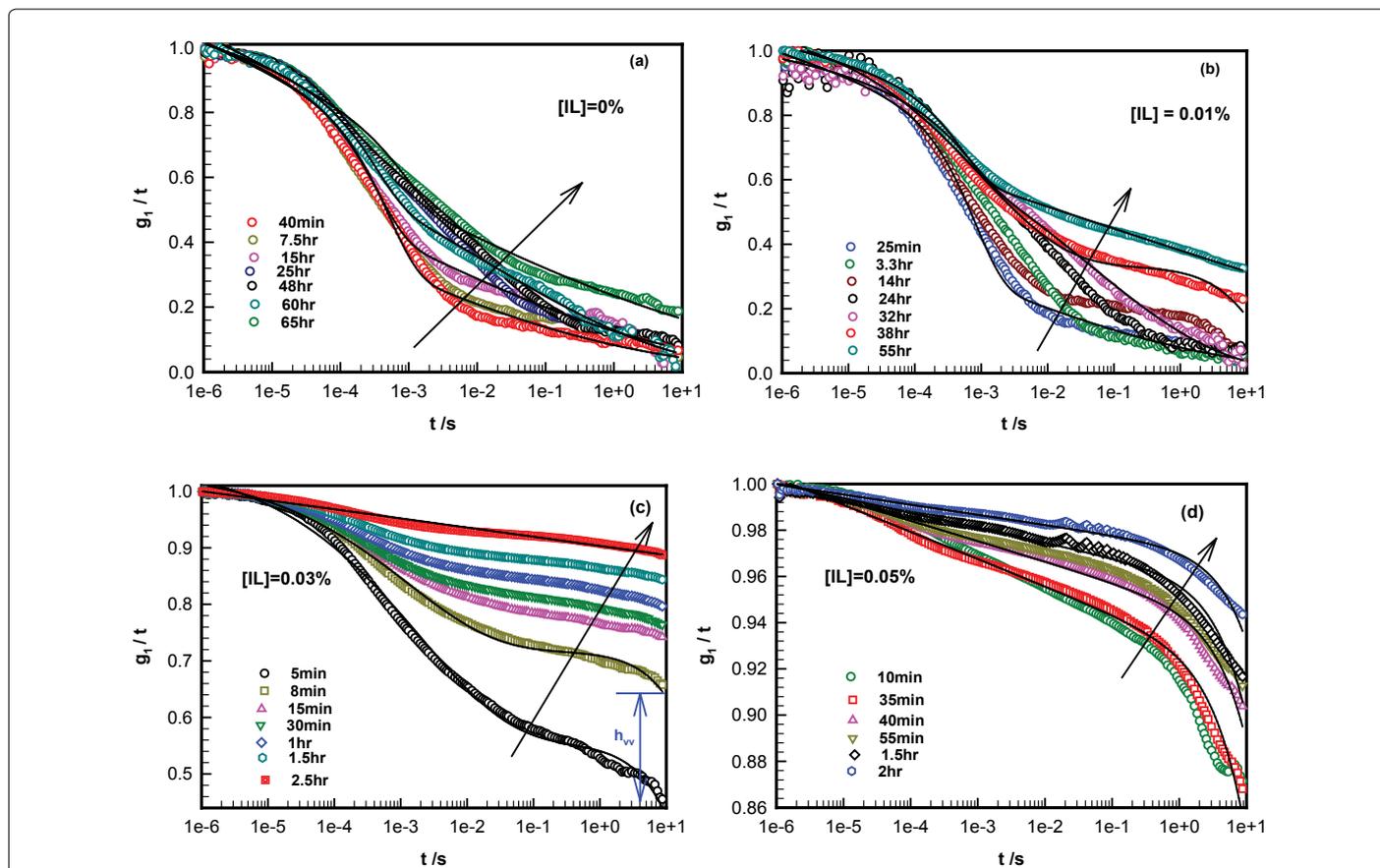


Figure 1: Waiting time evolution of normalized polarized autocorrelation function measured at an angle 90° for 2% (w/v) laponite with (a) 0% (b) 0.01%(c) 0.03% and (d) 0.05% (w/v) [IL]. The solid lines represent fits to equation (1).

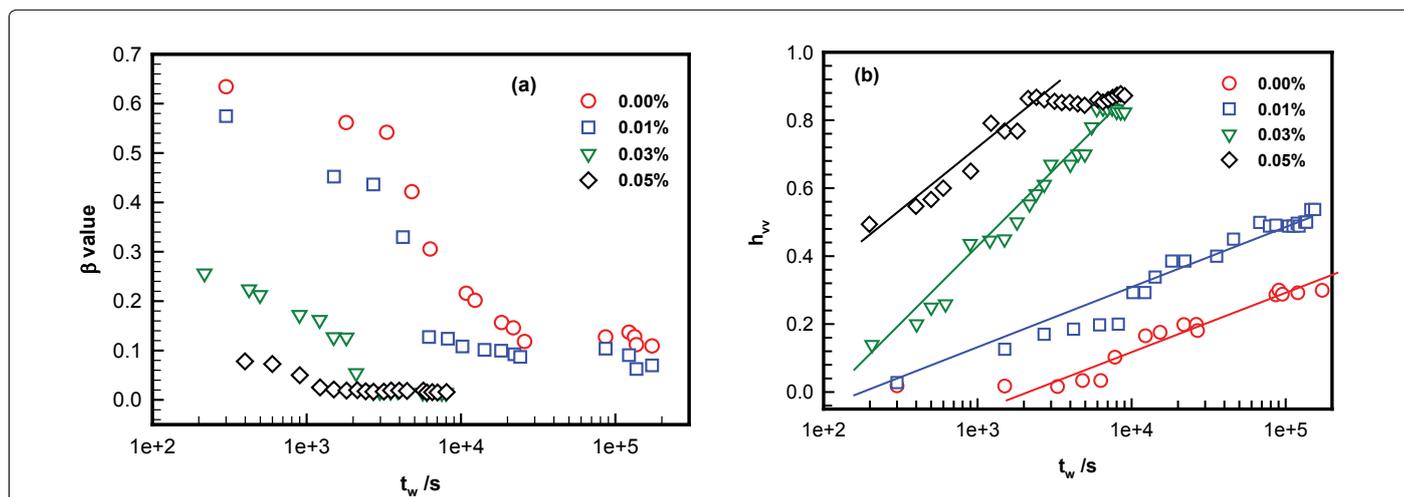


Figure 2: (a) A plot of stretching exponent as a function of waiting time for different angles with 2% laponite and different concentration of [IL]. (b) Value of $g_{v1}(t)$ at long delay time say h_w as a function of waiting time for laponite with different concentration of [IL].

above 0.03% the samples undergo much earlier transition as compared to others. The decrease in stretching parameter with waiting time implies that as the waiting time increases, the system is reaching to arrested state i.e transition from ergodic to non ergodic takes place. This decrease in stretching parameter shows that the dynamical properties of the sample change with waiting time and the structural rearrangement inside the sample during the transition from ergodic

to an arrested non ergodic state has been observed. However, the variation of stretch exponent β remains independent of wave vector (q) for different IL concentration.

It is impossible to determine transition from ergodic to non ergodic in the system because of low signal modulation, thus the value of $g_1(t)$ at long delay time ($t \rightarrow \infty$) abbreviated as h_w for polarized mode.

This h_{vv} value provides quantitative estimation of ergodicity and non ergodicity. This value of $g_1(q,t)$ at $t \rightarrow \infty$ which defines the degree of non-ergodicity (h_{vv} parameter) increases with aging time for all the samples. For fully ergodic systems it assumes a value of 0 whereas for a completely non ergodic system it has a value approximately equal to 1. Figure 2(b) shows variation of h_{vv} as a function of waiting time for different concentration of [C2mim][Cl] i.e. 0, 0.01, 0.03 and 0.05% (w/v), respectively, in which increase in h_{vv} value is observed.

The sharp change in value of the stretching exponent (β) and h_{vv} value with waiting time gives us the ergodicity breaking time (t_{EB}). The variation of ergodicity breaking time calculated from both the parameter of Figure 2 is shown in Figure 3. The aging time at which there is sharp change correspond to ergodic –nonergodic transition. Higher concentration of IL i.e. >0.03% (w/v) take less time as compared to low concentration to get arrested. From Figure 3 also we can demarcate two regions one below 0.03% where structured gels are formed and other above 0.03% IL where fragile gels are formed.

Next we study the angular dependence of the scattering function $g_1(q,t)$ shown in Figure 1. The correlation function were measured at different angles changing the scattering wave vector in the range $0.006 < q < 0.0231 \text{ nm}^{-1}$ ($30^\circ < \theta < 130^\circ$). The relaxation time τ_1 and τ_2 were obtained from the analysis of correlation functions.

However, observing their angular dependence we found that τ_1 varies as q^{-2} seen from Figure 4 (a). The q^{-2} dependence of τ_1 led us to conclude the diffusive nature of fast mode relaxation time. On observing the behaviour of τ_2 , it was also found to scale as q^{-2} . Thus the variation of relaxation time as q^{-2} confirms the diffusive nature in our system (Figure 4b).

To further investigate the characteristics of the slow mode, the time dependence of the relaxation rate of the fast and slow mode was determined. The behavior of slow mode relaxation time τ_2 for polarized correlation function is shown in Figure 5 as a function of waiting time in a logarithmic plot for our [IL] samples. From these we observed increase in slow mode relaxation time with increase in waiting time. The slow mode decay time of correlation function provide information for the time a particle needs to forget its initial position. From the figure representing the time dependence of slow relaxation rate, two distinct regions are evident where τ_2 progressively increased with waiting time initially and after certain time it become nearly constant. The increase in relaxation time τ_2 with waiting time implies that very rapidly the aging freezes in a certain degree of system. The increase shows the exponential behaviour in one region and linear in second region.

The relaxation time τ_1 correspond to fast relaxation mode was found to increase with aging time. The selective binding of polar head groups of monomeric and oligomeric IL molecules to negatively charged faces of laponite platelets dominates over the other usual interactions. The degree of charge neutralization thus achieved combined with the favourable hydrophobic interactions facilitated the gelation and arrest kinetics. The matured arrested phase was found to be gels in responses to dilution. The coherent picture that emerged was that IL molecules due to their large polar head groups and short hydrophobic tails are capable of exfoliating the colloidal dispersions and hence making better gels. ILs can produce interesting soft matter systems, both potential-wise and scientifically.

Conclusions

We have performed a systematic and comprehensive study to

establish the gelation kinetics of ionogels of laponite platelets in ionic liquid (IL), 1-ethyl-3-methyl imidazolium chloride, solutions at room temperature. The phenomenon of self-assembly of IL coated laponite platelets preceded gelation transition. In our earlier work, we discussed the two categories of ionogels produced one of which is a homogeneous and structured gel having high gel strength and the other one is a heterogeneous, fragile gel with low gel strength [39]. The binding was strongly dominated by electrostatic forces in low IL concentration samples whereas at higher IL concentration the network was held together by weak edge-face associations. Incidentally, the relaxations of density fluctuations were found to be governed by a fast and a slow mode. The fast mode contained two distinguishable components, one manifesting at low length scale and another operating at high length scale. Both these modes were diffusive and were IL concentration dependent. The slow mode, on the other hand, was observed to be frozen in space.

Ionic liquids are surfactant-like molecules with considerable surface activity. The complexity arises from the fact that the IL molecules interact with water molecules and laponite platelets simultaneously through electrostatic, hydrophobic and hydrogen bonding forces. In aqueous laponite dispersion (2% w/v) in presence of 1-methyl-3-ethylimidazolium chloride ionic liquid (IL, [C2mim][Cl], concentration 0.01-0.05% (w/v). With aging, the time-dependent relative viscosity increased and the time required for the system to get arrested decreases as we increase IL concentration 0 to 0.03% (w/v). Hence, the arrest kinetics slows down as we increase IL concentration [39]. This time dependent behaviour was measured from dynamic light scattering studies revealing formation of IL mediated finite size colloidal networks on a time scale $\approx 4 \times 10^3$ s whereas the dispersion developed a large viscosity a decade in time later ($\approx 4 \times 10^4$ s). Dynamic light scattering exhibited bimodal relaxation time of concentration fluctuations.

The behavior of laponite dispersions in ionic liquid (IL), 1-ethyl-3-methyl imidazolium, was studied by dynamic light scattering (DLS) in this manuscript. The results of this study suggest that the interaction between laponite dispersions may increase in a time-dependent manner. Hence, we conclude that the presence of ionic liquid greatly promotes the formation of gel network in aqueous laponite dispersions.

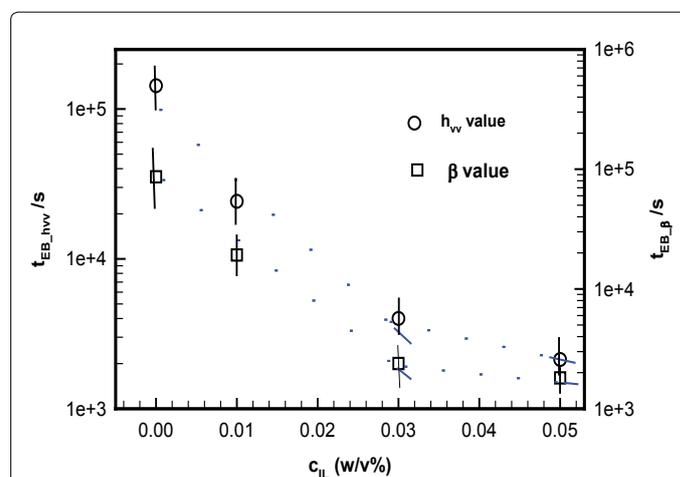


Figure 3: Value of ergodicity breaking time (tEB) as a function of [IL] concentration from h_{vv} and stretching exponent analysis. The vertical lines represent the error bars.

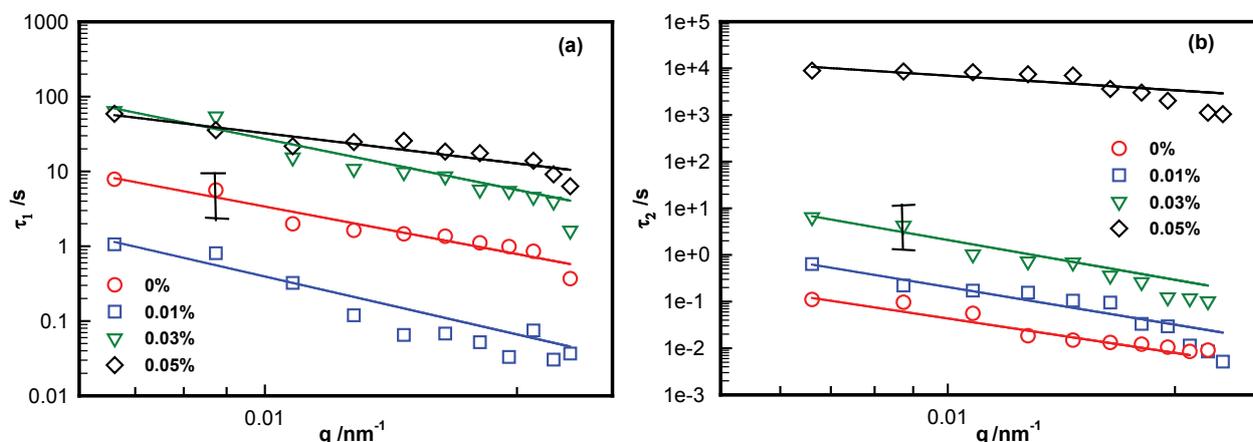


Figure 4: Plot of q dependence of (a) fast relaxation time (τ_1) and (b) slow relaxation time (τ_2) for 2% (w/v) laponite with different [IL] concentration ($t = 0$). Representative error is shown.

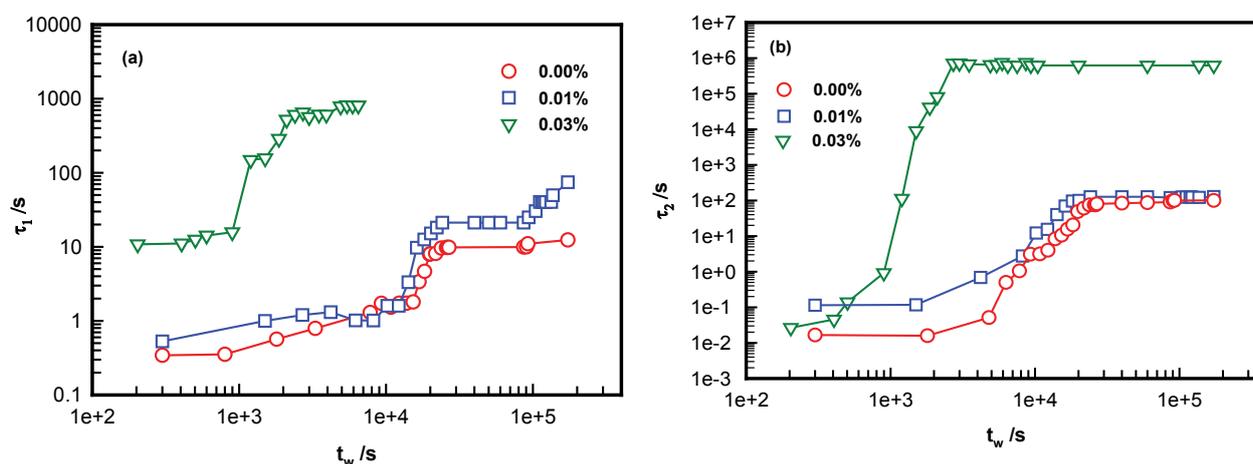


Figure 5: Evolution of the fast and slow relaxation times as function of waiting time, t_w for samples with 2% (w/v) laponite and different IL concentration.

Acknowledgment

NJ acknowledges University Grants Commission, Government of India for a research fellowship. KR is thankful to Department of Science and Technology, Government of India-Inspire Faculty Award. This work was supported by a research grant received from Department of Science and Technology, Government of India.

References

- Dijkstra M, Hansen JP, Madden PA (1995) Gelation of a clay colloid suspension. Phys Rev Lett 75: 2236-2239.
- Bonn D, Tanaka H, Wegdam G, Kellay H, Meunier J (1998) Aging of a colloidal "Wigner" glass. J Europhys Lett 45: 52-57.
- Coussot P, Nguyen QD, Huynh HT, Bonn D (2002) Avalanche behavior in yield stress fluids. Phys Rev Lett 88: 175501.
- Ruzicka B, Zulian L, Ruocco G (2004) Routes to gelation in a clay suspension. Phys Rev Lett 93: 258301.
- Abou B, Gallet F (2004) Probing a nonequilibrium einstein relation in an aging colloidal glass. Phys Rev Lett 93: 160603.
- Cummins HZ (2007) Liquid, glass, gel: The phases of colloidal Laponite. J Non-Cryst Solids 353: 3891-3905.
- Ruzicka B, Zaccarelli E, Zulian L, Angelini R, Sztucki M, et al. (2011) Observation of empty liquids and equilibrium gels in a colloidal clay. Nat Mater 10: 56-60.
- Taki B, John S, Arakawa, Okamoto M (2013) Structure and rheology of nanocomposite hydrogels composed of DNA and clay. Eur Polym J 49: 923-931.
- Shalkevich A, Stradner A, Bhat SK, Muller F, Schurtenberger P (2007) Cluster, glass, and gel formation and viscoelastic phase separation in aqueous clay suspensions. Langmuir 23: 3570-3580.
- Leger D, Levesque D (2002) Monte Carlo simulations of charged platelet systems. J Chem Phys 116: 2251-2260.
- Jönsson B, Labbez C, Cabane B (2008) Interaction of nanometric clay platelets. Langmuir 24: 11406-11413.
- Huang X, Netravali AN (2006) Characterization of nano-clay reinforced phytigel-modified soy protein concentrate resin. Biomacromolecules 7: 2783-2789.
- Marsh KN, Brennecke JF, Chirico RD, Frenkel M, Heintz A, et al. (2009) Thermodynamic and thermophysical properties of the reference ionic liquid: 1-Hexyl-3-methylimidazolium bis[(trifluoromethyl)sulfonyl]amide (including mixtures). Part 1. Experimental methods and results. Pure Appl Chem 81: 781-828.
- Wasserscheid P, Keim W (2000) Ionic Liquids-New "Solutions" for Transition Metal Catalysis. Angew Chem Int Ed Engl 39: 3772-3789.
- Martins MAP, Frizzo CP, Moreira DN, Nilo Z, Bonaccorso HG (2008) Ionic Liquids in Heterocyclic Synthesis. Chem Rev 108: 2015-2050.
- Hagiwara R, Ito Y (2000) Room temperature ionic liquids of alkyimidazolium cations and fluoroanions. J Fluor Chem 105: 221-227.

17. Ohno H, Fukumoto K (2007) Amino acid ionic liquids. *Acc Chem Res* 40: 1122-1129.
18. Earle MJ, Seddon KR (2000) Ionic liquids. Green solvents for the future. *Pure Appl Chem* 72: 1391-1398.
19. Ruzicka B, Zaccarelli E (2011) A fresh look at the Laponite phase diagram. *Soft Matter* 7: 1268-1286.
20. Pujala RK, Pawar N, Bohidar HB (2011) Universal sol state behavior and gelation kinetics in mixed clay dispersions. *Langmuir* 27: 5193-5203.
21. Pujala RK, Bohidar HB (2012) Ergodicity breaking and aging dynamics in Laponite–Montmorillonite mixed clay dispersions. *Soft Matter* 8: 6120-6127.
22. Mongondry P, Tassin JF, Nicolai T (2005) Revised state diagram of Laponite dispersions. *J Colloid Interface Sci* 283: 397-405.
23. Bala P, Samantaray BK, Srivastava SK (2000) Dehydration transformation in Ca-montmorillonite. *Bull Mater Sci* 23: 61-67.
24. Tanaka H, Meunier J, Bonn D (2004) Nonergodic state of charged colloidal suspensions: repulsive and attractive glasses and gels. *Phys Rev E* 69: 031404-1-031404-6.
25. Ruzicka B, Zulian L, Ruocco G (2006) More on the phase diagram of Laponite. *Langmuir* 22: 1106-1111.
26. Glotzer SC, Solomon MJ (2007) Anisotropy of building blocks and their assembly into complex structures. *Nat Mater* 6: 557-562.
27. Winterton N (2006) Solubilization of polymers by ionic liquids. *J Mater Chem* 16: 4281-4293.
28. Le Bideau J, Viau L, Vioux A (2011) Ionogels, ionic liquid based hybrid materials. *Chem Soc Rev* 40: 907-925.
29. He Y, Lodge TP (2007) A thermoreversible ion gel by triblock copolymer self-assembly in an ionic liquid. *Chem Commun (Camb)* : 2732-2734.
30. Lei Y, Lodge TP (2012) Effects of component molecular weight on the viscoelastic properties of thermoreversible supramolecular ion gels via hydrogen bonding. *Soft Matter* 8: 2110–2120.
31. Kofu M, Someya T, Tatsumi S, Ueno K, Ueki T, et al. (2008) Microscopic insights into ion gel dynamics using neutron spectroscopy. *Soft Matter* 8: 7888-7897.
32. Matsumoto K, Endo T (2008) Confinement of Ionic Liquid by Networked Polymers Based on Multifunctional Epoxy Resins. *Macromolecules* 41: 6981-6986.
33. Stanzione JF 3rd, Jensen RE, Costanzo PJ, Palmese GR (2012) Synthesis and characterization of ionic polymer networks in a room-temperature ionic liquid. *ACS Appl Mater Interfaces* 4: 6142-6150.
34. Fuller J, Breda AC, Carlin RT (1997) Ionic Liquid-Polymer Gel Electrolytes. *J Electrochem Soc* 144: L67-L70.
35. Ribot JC, Guerrero-Sanchez C, Hoogenboom R, Schubert US (2010) Aqueous gelation of ionic liquids: reverse thermoresponsive ion gels. *Chem Commun (Camb)* 46: 6971-6973.
36. Wasserscheid P, Welton T, Wiley-VCH, Weinheim, (2003) *Ionic Liquids in Synthesis*.
37. Rogers RD, Seddon KR (2003) *American Chemical Society? Washington, DC*.
38. Abou B, Bonn D, Meunier J (2001) Aging dynamics in a colloidal glass. *Phys Rev E Stat Nonlin Soft Matter Phys* 64: 021510.
39. Joshi N, Rawat K, Bohidar HB (2014) *Colloids and Surfaces A: Physicochemical and Engineering Aspects*. Elsevier 461: 66-75.