Numerical Study of Falling Film Thickness on Horizontal Circular Tube - A CFD Approach

Arun Kumar
National Institute of Technology Kurukshetra, Haryana, India

*Corresponding author: Arun Kumar, National Institute of Technology Kurukshetra, Haryana, India, Tel: +91 9728430195; E-mail: arun.mit2006@gmail.com

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Introduction

Evaporation of thin liquid falling films from horizontal circular tubes is of fundamental interest and widely used in refrigeration equipment, chemical and food industry such as heat pipes and capillary pumped loops, spray and jet impingement cooling, falling film evaporation and so on [1–4] as it provides high heat transfer rate and small use of liquid. For a fully developed laminar flow, heat transfer across the film is approximately evaluated by one dimensional heat conduction. The circumferential thickness of a falling liquid film is inversely proportional to the local heat transfer coefficient. Hence, the local film thickness is an important property to characterize the falling film and heat transfer phenomena. This in turn indicates that deeper insight into the film thickness distributions is of critical importance for many applications.

Since the characteristics of falling films have drawn great attention, the researches on the film thickness have been conducted by both theoretical and experimental approaches in the past years. Nusselt [5-7] first analyzed falling films theoretically, assuming a continuous sheet flow from tube to tube and momentum effects on the falling films to be negligible and derived the following expression for the film thickness:

$$\delta = \left( \frac{3H \Gamma}{\rho_L (\rho_L - \rho_G) g \sin \beta} \right)^{1/3}$$

Where $\beta$ is the circumferential angle measured from the top of a horizontal tube and $\Gamma$ is the liquid flow rate on one side of the tube. With this definition, the film Reynolds number is given below

$$Re = \frac{4\Gamma}{\mu_L}$$

Zhang et al. [8-14] presented theoretical and experimental investigations of the film thickness and temperature profile on a vertical heated/cooled plate. Xu et al. [15] studied the flowing state in liquid films outside horizontal tubes. Numerical methods based on the Computed Fluid Dynamic were also applied to explore the construction and instability of the film [16,17].

Adams et al. [8] reported falling film for heat exchangers and suggested three advantages of the falling film: the pressure drop of the falling film is negligible, the quantity of liquid required is small, and the heat transfer coefficients are high. Kocamustafaogullari et al. [9] have simulated the hydrodynamics and heat transfer of the liquid falling film on a tube to obtain the film thickness and local average heat transfer coefficients for both constant heat flux and isothermal boundary conditions. Rogers et al. [10] solved the motion and energy equations for the laminar liquid film flow by considering following: film thickness as a function of Reynolds number, Archimedes number and the angular position. Mitrovic [11] studied experimentally that the heat transfer and the mechanism of the fluid flow from a horizontal heated tube to a sub cooled liquid falling film. The experimental results indicated that the flow pattern of the falling liquid depends on both Reynolds number and the tube spacing [12].

However, despite above numerous theoretical and experimental studies, there has been little information available for film flow characteristics on horizontal tubes. In this paper, we present a detail numerical simulation of the flow field, which could not be observed experimentally, on horizontal circular tubes with different tube spacing and Reynolds numbers. First, the experimental data was numerically simulated and validated with the help of Ansys fluent using volume of fluid model. Then, the thickness of liquid film on tube were investigated. Furthermore, the effect of tube spacing and Reynolds numbers on film thickness was considered. Finally, numerical results were compared to theoretical values by Nusselt expression and the reported values.

Numerical Approach

Analysis geometry and physical condition

Ansys fluent has been employed to model and mesh the domain of the geometry. A tri pave mesh has been generated for the 2D domain. The region around the cylinder wall was carefully meshed by using a
boundary layer technique to capture details of the liquid flow around
the cylinder and the tracking of a liquid-gas interface for 2D flow field
simulations which is shown in Figure 1a and Figure 1b respectively.
The domain consists of three phases, where gas represents the
surrounding air, liquid represents water and the solid phase is the wall
for the cylinders. A size function technique has been employed which
enables starting from a minimal mesh size where flow details are
important. It was assumed here that the falling film flow over the
horizontal tube is incompressible, low-velocity and viscous flow
derived by gravity downward as shown in Figure 2.

The properties of the surrounding gas phase, air, and the liquid
phase, water, were considered constant under the normal condition
with temperature of 25°C and pressure of 101.325 kPa. Considering its
symmetrical structure, the shaded region was chosen as the solution
domain to minimize the computing time. The solution domain was
discretized by quadrilateral elements, as shown in Figure 1b. The total
number of mesh elements and the time step were kept 70,147 and 0.1
ms respectively. The grid size increases with distance from the cylinder
surface, as flow gradients become sever. The size function was chosen
here 0.003 m with growth rate of 1.1 to a size limit of 1 (Table 1).

Table 1: Fluent 2D simulation settings.

<table>
<thead>
<tr>
<th>S No.</th>
<th>Settings</th>
<th>Choice</th>
</tr>
</thead>
<tbody>
<tr>
<td>1</td>
<td>Simulation</td>
<td>2D</td>
</tr>
<tr>
<td>2</td>
<td>Solver</td>
<td>Segregated implicit</td>
</tr>
<tr>
<td>3</td>
<td>Model</td>
<td>VOF</td>
</tr>
<tr>
<td>4</td>
<td>Material</td>
<td>Air – primary and water – secondary</td>
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<tr>
<td>5</td>
<td>Gravitational acceleration</td>
<td>9.81</td>
</tr>
<tr>
<td>6</td>
<td>Volume fraction</td>
<td>1 for water and 0 for air</td>
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<tr>
<td>7</td>
<td>Pressure Velocity Coupling</td>
<td>PISO</td>
</tr>
<tr>
<td>8</td>
<td>Discretization Pressure</td>
<td>PRESTO</td>
</tr>
<tr>
<td>9</td>
<td>Discretization Momentum</td>
<td>1 Order Upwind</td>
</tr>
<tr>
<td>10</td>
<td>Top boundary</td>
<td>Velocity inlet and pressure inlet</td>
</tr>
<tr>
<td>11</td>
<td>Left side boundary</td>
<td>Pressure outlet</td>
</tr>
<tr>
<td>12</td>
<td>Right side boundaries</td>
<td>Symmetry</td>
</tr>
<tr>
<td>13</td>
<td>Bottom side boundary</td>
<td>Pressure outlet</td>
</tr>
<tr>
<td>14</td>
<td>Wall boundary</td>
<td>Cylinder’s edge</td>
</tr>
</tbody>
</table>

Numerical Results and Discussion

Verification of computational method

The comparison of the calculated results from Nusselt correlation
[19], the experimental results by Gstoehl et al. [20] and simulated
results of this paper are plotted against the circumferential angles and
film thickness shown in Figure 3a and b. It is seen that the relative
error between the simulation results of this paper and the calculated
results from Nusselt correlation is within the range of 20%
approximately and the variation trend of the simulation results of this
paper well agrees with the experimental data by Gstoehl et al. [20]. It
should be noticed that the thickness of falling liquid film between 0
and 40 by Gstoehl et al. [20] does not vary monotonously. This is
casted due the measurement error and should not be considered as a
normal phenomenon. The comparison results shows that the VOF

Governing equations

To track the interface between two phases, volume of fluid (VOF)
model is used to compute two separate phase flow [18] whereas Fluent
solves conservation equations for mass and momentum and volume
fraction function [17].

Mass equation
$$\frac{\partial}{\partial t}(\rho \mathbf{v}) + \nabla \cdot (\rho \mathbf{v} \mathbf{v}) = 0$$

Momentum equation
$$\frac{\partial}{\partial t}(\rho \mathbf{v} \mathbf{v}) + \nabla \cdot (\rho \mathbf{v} \mathbf{v}) = -\nabla p + \nabla \cdot (\mu(\nabla \mathbf{v} + \nabla \mathbf{v}^T)) + \rho \mathbf{g} + \mathbf{f}$$

Volume fraction
$$\frac{\partial}{\partial t}(\alpha_q) + \nabla \cdot \mathbf{v} \alpha_q = 0$$
method is capable of capturing the characteristics of falling film thickness on horizontal tube under adiabatic condition. Therefore, this method is recommended for predicting the film thickness characteristics over horizontal tube.

The effect of Reynolds number (Re) and tube spacing(s)

When the surface of tube is fully covered by the liquid film, the following heat and mass transfer during the process is considered as steady state with constant boundary conditions. The effect of Reynolds number on the distribution of film thickness in the circumferential direction of a horizontal circular tube with tube spacing (s), 19.4 mm is shown in Figure 3b and the circumferential distributions of falling film thickness for all three different Reynolds numbers are similar to each other. It means with the increment of circumferential angle, the film thickness decreases first and then increases after the minimum value of film thickness which appears in the circumferential range between 100° to 140°. However, the Reynolds number affects the amount of the film thickness and this result is consistent with the experimental result by Hou et al. [21], who has verified the effect laws of Reynolds number on film thickness by using displacement micrometre.

![Figure 3: Film thickness distribution versus circumferential angle for (a) comparison of results of this paper and result by other researchers and (b) the effects of Reynolds numbers.](image)

The tube surface which is covered by liquid film has two typical positions in the circumferential direction one is the upper stagnation point, 0° and second is the lower stagnation point, 180° where the circumferential velocity of liquid film equals to zero. The stagnation film thickness increases as Reynolds number increases. Apparently the centrifugal force induced due the circular motion of liquid film that is directly related to the liquid film velocity which has overcome the adhesive force between the tube surface and the liquid film. The most of the liquid is thrown off the tube surface even the lower stagnation point is not reached. Compared to this part of liquid, the liquid film which stays on the tube surface has very low velocity. The corresponding centrifugal force is not big enough to overcome the adhesive force. Therefore, it continues to flow along the tube surface till the lower stagnation point is reached that means the liquid stays longer over the tube surface when the Reynolds number is small. The similar phenomenon near the lower stagnation point was also observed by Min and Choi [22] when the Reynolds number is greater than 40. The film thickness on the top of the tube tends to decrease with increased tube spacing, which is logical because of the added momentum of the falling liquid by Gstoehl et al. [20].

Conclusion

The falling liquid film is in steady state after it fully covers the surface of tube under the constant boundary conditions. The circumferential distribution of falling film thickness over the tube surface is not uniform which decreases from the top of the horizontal tube and then increases after it reaches to the minimum value of approximately which appears within the circumferential angle range between 90° to 140°. The asymmetric distribution, contradictory to the Nusslet theory is supported by the results of various experimental researches. The circumferential distribution of falling film thickness along the tube surface corresponds to the variation of liquid film velocity which indicates that the momentum effect on the flow cannot be neglected.

The falling film thickness increases with the increase in Reynolds number which is directly related to the mass flow rate of liquid. The sizeable zone that is lack of liquid is formed near the lower stagnation point and the zone size is bigger for higher Reynolds number.

The film thickness on the top of the tube tends to decrease with increased tube spacing, which is logical because of the added momentum of the falling liquid.

References


