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# On the First aff(1)-Relative Cohomology of the Lie Algebra of Vector Fields and Differential Operators 

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#### Abstract

Let $\operatorname{Vect}\left(\mathbb{R P}^{1}\right)$ be the Lie algebra of smooth vector fields on $\mathbb{R} \mathbb{P}^{1}$. In this paper, we classify $\mathfrak{a f f}(1)$-invariant linear differential operators from $\operatorname{Vect}\left(\mathbb{R} \mathbb{P}^{1}\right)$ to $\mathcal{D}_{\lambda, \mu, v}$ vanishing on $\mathfrak{a f f}(1)$, where $\mathcal{D}_{\lambda, \mu ; j}:=\operatorname{Homdiff}\left(\mathcal{F}_{\lambda} \otimes \mathcal{F}_{\mu} ; \mathcal{F}_{v}\right)$ is the space of bilinear differential operators acting on weighted densities. This result allows us to compute the first differential $\mathfrak{a f f}(1)$-relative cohomology of $\operatorname{Vect}\left(\mathbb{R}^{1}\right)$ with coefficients in $\mathcal{D}_{\lambda, \mu ; j}$.


Keywords: Differential operators; Transvectants; Lie algebra; Cohomology

## Introduction

Let $\mathfrak{g}$ be a Lie algebra and let $\mathcal{M}$ and $\mathcal{N}$ be two $\mathfrak{g}$-modules. It is wellknown that nontrivial extensions of $\mathfrak{g}$-modules:

$$
0 \rightarrow \mathcal{M} \rightarrow . \rightarrow \mathcal{N} \rightarrow 0
$$

are classified by the first cohomology group $\mathrm{H}^{1}(\mathfrak{g} ; \operatorname{Hom}(\mathcal{N}, \mathcal{M}))$ [1]. Any 1-cocycle $\mathcal{L}$ generates a new action on $\mathcal{M} \oplus \mathcal{N}$ as follows: for all $g \in \mathfrak{g}$ and for all $(a, b) \in \mathcal{M} \oplus \mathcal{N}$, we define $g^{*}(a, b):=\left(g^{*} a+C^{t t} \mathcal{L}(b), g^{*} b\right)$. For the space of tensor density of weight $\lambda, \mathcal{F}_{\lambda}$, viewed as a module over the Lie algebra of smooth vector fields Vect $\left(\mathbb{R P}^{1}\right)$, the classification of nontrivial extensions

$$
0 \rightarrow \mathcal{F}_{\mu} \rightarrow . \rightarrow \mathcal{F}_{\lambda} \rightarrow 0,
$$

leads Feigin and Fuks [2] to compute the cohomology group $\mathrm{H}^{1}\left(\operatorname{Vect}\left(\mathbb{R P}^{1}\right) ; \operatorname{Hom}\left(\mathcal{F}_{\lambda}, \mathcal{F}\right)\right)$. Later, Ovsienko and Bouarroudj [3] have computed the corresponding relative cohomology group with respect to $\mathfrak{s l}(2, \mathbb{R})$, namely
$\mathrm{H}^{1}\left(\operatorname{Vect}\left(\mathbb{R P}^{1}\right), \mathfrak{s l}(2, \mathbb{R}) ; \operatorname{Hom}\left(\mathcal{F}_{\lambda}, \mathcal{F}_{\mu}\right)\right)$.
In this paper, we will compute the first cohomology group
$\mathrm{H}^{1}\left(\operatorname{Vect}\left(\mathbb{R P}^{1}\right), \mathfrak{a f f}(1) ; \operatorname{Hom}\left(\mathcal{F}_{\lambda} \otimes, \mathcal{F}_{\mu}, \mathcal{F}_{\nu}\right)\right)$.

## $\operatorname{Vect}(\mathbb{R})$-Module Structures on the Space of Bilinear Differential Operators

Consider the standard (local) action of $\mathfrak{a f f}(1)$ on $\mathbb{R}$ by linearfractional transformations. Although the action is local, it generates global vector fields

$$
\left\{\frac{d}{d x}, x \frac{d}{d x}\right\},
$$

that form a Lie subalgebra of $\operatorname{Vect}(\mathbb{R})$ isomorphic to the Lie algebra $\mathfrak{a f f}(1)$. This realization of $\mathfrak{a f f}(1)$ is understood throughout this paper.

## The space of tensor densities on $\mathbb{R} \mathbb{P}^{1}$

The space of tensor densities of weight $\lambda$ (or $\lambda$-densities) on $\mathbb{R P}^{1}$, denoted by:

$$
\mathcal{F}_{\lambda}=\left\{f(d x)^{\lambda} \mid f \in C^{\infty}(\mathbb{R})\right\}, \lambda \in \mathbb{R},
$$

is the space of sections of the line bundle $\left(T^{*} \mathbb{R} \mathbb{P}^{1}\right)^{\otimes^{\lambda}}$. This space coincides with the space of functions and differential forms for $\lambda=0$ and for $\lambda=1$, respectively. The Lie algebra $\operatorname{Vect}\left(\mathbb{R P}^{1}\right)$ acts on $\mathcal{F}_{\lambda}$ by the Lie derivative. For all $X \in \operatorname{Vect}\left(\mathbb{R P}^{1}\right)$ and for all $\varphi \in \mathcal{F}_{\lambda}$ :

$$
\begin{equation*}
L_{X}\left(\varphi(d x)^{\lambda}\right)=X \varphi^{\prime}+\lambda \varphi X^{\prime}, \tag{1}
\end{equation*}
$$

where the superscript ' stands for $d / d x$.
The space of bilinear differential operators as a $\operatorname{Vect}\left(\mathbb{R}^{1} \mathbb{P}^{1}\right)$ module

We are interested in defining a three-parameter family of $\operatorname{Vect}\left(\mathbb{R} \mathbb{P}^{1}\right)$-modules on the space of bilinear differential operators. The counterpart $\operatorname{Vect}\left(\mathbb{R}^{1} \mathbb{P}^{1}\right)$-modules of the space of linear differential operators is a classical object [4].

Consider bilinear differential operators that act on tensor densities:

$$
\begin{equation*}
A: \mathcal{F}_{\lambda} \otimes \mathcal{F}_{\mu} \rightarrow \mathcal{F}_{v} \tag{2}
\end{equation*}
$$

The Generalized Lie algebra Vect $\left(\mathbb{R P}^{1}\right)$ acts on the space of bilinear differential operators as follows. For all $\varphi \in \mathcal{F}_{\lambda}$ and for all $\psi \in \mathcal{F}_{\mu}$ :

$$
\begin{equation*}
L_{X}^{\lambda, \mu, \nu}(A)(\phi, \psi)=L_{X}^{\nu} \circ A(\phi, \psi)-A\left(L_{X}^{\lambda} \phi, \psi\right)-A\left(\phi, L_{X}^{\mu} \psi\right) \tag{3}
\end{equation*}
$$

where $L_{X}^{\lambda}$ is the action (1). We denote by $\mathcal{D}_{\lambda, 4,4 \nu}$ the space of bilinear differential operators (2) endowed with the defined $\operatorname{Vect}\left(\mathbb{R}^{1} \mathbb{P}^{1}\right)$-module structure (3).

## Relative Cohomology

Let us first recall some fundamental concepts from cohomology theory [1]. Let $g$ be a Lie algebra acting on a vector space $V$ and let $h$ be a sub- algebra of $g$. (If $h$ is omitted it assumed to be $\{0\}$.) The space of $h$-relative $n$-cochains of $g$ with values in $V$ is the $g$-module

$$
C^{n}(g, h ; V):=\operatorname{Hom}_{h}\left(\Lambda^{n}(g / h) ; V\right)
$$

The coboundary operator $\delta_{n}: C^{n}(g, h ; V) \rightarrow C^{n+1}(g, h ; V)$ is a $g$-map satisfying ${ }_{n} \circ \delta_{n 1}=0$. The kernel of $\delta_{n}^{n}$, denoted $Z^{n}(\mathfrak{g}, h ; V)$, is the space of $h$-relative $n$-cocycles, among them, the elements in the range of $\delta_{n-1}$ are called $h$-relative $n$-coboundaries. We denote $B^{n}(\mathfrak{g}, h ; V)$ the space of $n$-coboundaries.

[^0]By definition, the $n^{\text {th }} h$-relative cohomolgy space is the quotient space

## $\mathrm{H}^{n}(\mathfrak{g}, h ; V)=Z^{n}(\mathfrak{g}, h ; V) / B^{n}(\mathfrak{g}, h ; V)$.

We will only need the formula of $\delta_{n}$ (which will be simply denoted $\delta$ ) in degrees 0,1 and 2 : for $v \in C^{0}(g, h ; V)=V^{h}, \delta v(g):=(-1)^{\mid g\| \|} g . v$, where
$V^{h}=\{v \in V \mid h . v=0$ for all $h \in h\}$,
and for $\Upsilon \in C^{1}(g, h ; V)$,
$\delta(\Upsilon)(x, y):=x \cdot \Upsilon(y)-y \cdot \Upsilon(x)-\Upsilon([x, y])$ for any $x, y \in g$.

## $\mathfrak{a f f}(\mathbf{1})$-Invariant Differential Operators

The following steps to compute the relative cohomology has intensively been used in refs. [3,5-8]. First, we classify $\mathfrak{a f f}(1)$-invariant differential operators, then we isolate among them those that are 1-cocycles. To do that, we need the following Lemma.

## Lemma 4.1

Any 1 -cocycle vanishing on the subalgebra $\mathfrak{a f f}(1)$ of $\operatorname{Vect}(\mathbb{R})$ is $\mathfrak{a f f}(1)$-invariant.

The 1-cocycle condition of $\Upsilon$ reads:

$$
\begin{equation*}
X \cdot \Upsilon(Y)-Y \cdot \Upsilon(X)-\Upsilon([X, Y])=0 \tag{4}
\end{equation*}
$$

where $X, Y \in \operatorname{Vect}\left(\mathbb{R P}^{1}\right)$. Thus, if $\Upsilon(X)=0$ for all $X \in \operatorname{aff}(1)$, eqn. (4) becomes

$$
\Upsilon([X, Y])=X \cdot \Upsilon(Y)
$$

expressing the $\mathfrak{a f f}(1)$-invariance property of $\Upsilon$.
As our 1-cocycles vanish on aff(1), we will investigate aff(1)invariant linear differential operators that vanish on aff( 1 ).

Proposition 4.2: There exist $\mathfrak{a f f}(1)$-invariant bilinear differential operators $J_{k}^{\lambda, \mu}: \mathcal{F}_{\lambda} \otimes \mathcal{F}_{\mu} \rightarrow \mathcal{F}_{\lambda+\mu+k}$ given by:

$$
\begin{equation*}
J_{k}^{\lambda, \mu}\left(\varphi d x^{\lambda}, \phi d x^{\mu}\right)=\sum_{i+j=k} \gamma_{i, j} \varphi^{(i)} \phi^{(j)} d x^{\lambda+\mu+k} \tag{5}
\end{equation*}
$$

where $k \in \mathbb{N}$ and the coefficients $\gamma_{i, j}$ are constants.
Proof. Any differential operator $J_{k}^{\lambda, \mu}: \mathcal{F}_{\lambda} \otimes \mathcal{F}_{\mu} \rightarrow \mathcal{F}_{v}$ is of the form

$$
J_{k}^{\lambda, \mu}\left(f d x^{\lambda}, g d x^{\mu}\right)=\sum_{n=0}^{m} \sum_{i+j=k} \gamma_{i, j} f^{(i)} g^{(j)} d x^{v}, \quad m \in \mathbb{N}
$$

The $\operatorname{osp}(1 \mid 2)$-invariant property of the operators $J_{k}^{\lambda, \mu}$ with respect to the vector field $X=x \frac{d}{d x}$ yields:
$\frac{d}{d x} \gamma_{i, j}=0 \quad$ and $\quad v-\lambda-\mu=k \quad$ with $\quad k=i+j$.
So, we see that the corresponding operator can be expressed as (5).
Proposition 4.3: There exist $\mathfrak{a f f}(1)$-invariant trilinear differential operators $K_{k}^{\tau, \lambda, \mu}: \mathcal{F}_{\tau} \otimes \mathcal{F}_{\lambda} \otimes \mathcal{F}_{\mu} \rightarrow \mathcal{F}_{\tau+\lambda+\mu+k}$ given by:

$$
\begin{equation*}
K_{k}^{\tau, \lambda, \mu}(\varphi, \phi, \psi)=\sum_{i+j+l=k} \gamma_{i, j, l} \varphi^{(i)} \phi^{(j)} \psi^{(l)} \tag{6}
\end{equation*}
$$

where $i+j+l=k$ and the coefficients $\gamma_{i, j l}$ are constants.
If $\tau, \lambda$ and $\mu$ are generic, then the space of solutions is $\frac{1}{2}(k+1)(k+2)$ -dimensional.

Proposition 4.4: There exist $\mathfrak{a f f}(1)$-invariant trilinear differential
operators $K_{k}^{\lambda, \mu}: \operatorname{Vect}\left(\mathbb{R P}^{1}\right) \otimes \mathcal{F}_{\lambda} \otimes \mathcal{F}_{\mu} \rightarrow \mathcal{F}_{\lambda+\mu+k-1}$ that vanishe on $\mathfrak{a f f}(1)$ given by:

$$
\begin{equation*}
K_{k}^{\lambda, \mu}(X, \phi, \psi)=\sum_{i+j+l=k} \gamma_{i, j, l} X^{(i)} \phi^{(j)} \psi^{(l)} \tag{7}
\end{equation*}
$$

where $i+j+l=k$ and the coefficients $\gamma_{i, j, l}$ are constants but $\gamma_{0, j, k-j}=\gamma_{1, j, k-}$
 all $\lambda$ and $\mu$.

Proof of Proposition 4.3 and 4.4: We are going to prove Proposition 4.3 and 4.4 simultaneously. Any differential operator $K_{k}^{\tau, \lambda, \mu}: \mathcal{F}_{\tau} \otimes \mathcal{F}_{\lambda} \otimes \mathcal{F}_{\mu} \rightarrow \mathcal{F}_{\tau+\lambda+\mu+k}$ is of the form

$$
\begin{equation*}
K_{k}^{\tau, \lambda, \mu}(\varphi, \phi, \psi)=\sum_{i+j+l=k} \gamma_{i, j, l} \varphi^{(i)} \phi^{(j)} \psi^{(l)} . \tag{8}
\end{equation*}
$$

where $\gamma_{K_{i, j}, l}$ are functions. The $\mathfrak{a f f}(1)$-invariant property of the operators $K_{k}^{\tau, \lambda, \mu}$ reads as follows.

$$
L_{x}^{\nu} K_{k}^{\tau, \lambda, \mu}(\phi, \varphi, \psi,)=K_{k}^{\tau,, \mu, \mu}\left(L_{x}^{\tau} \phi, \varphi, \psi\right)+K_{k}^{\tau, \lambda, \mu}\left(\phi, L_{x}^{\lambda} \varphi, \psi\right)+K_{k}^{\tau, \lambda, \mu}\left(\phi, \varphi, L_{x}^{\mu} \psi\right) .
$$

The invariant property with respect to the vector field $X=\frac{d}{d x}$ implies that $\gamma_{i, j, l}^{\prime}=0$. On the other hand, the invariant propetty with respect to the vector fields $X=x \frac{d}{d x}$ implies that $v=\tau+\lambda+\mu+k$. If $\tau, \lambda$ and $\mu$ are generic, then the space of solutions is $\frac{1}{2}(k+1)(k+2)$
-dimensional, spanned by -dimensional, spanned by

$$
\begin{align*}
& \gamma_{0,0, k}, \gamma_{0,1, k-1}, \cdots, \gamma_{0, k, 0} \\
& \gamma_{1,0, k-1}, \gamma_{1,1, k-2}, \cdots, \gamma_{1, k-1,0} \\
& \quad \vdots  \tag{10}\\
& \gamma_{k-1,0,1}, \gamma_{k-1,1,0} \\
& \gamma_{k, 0,0}
\end{align*}
$$

Now, the proof of Proposition 4.4 follows as above by putting $\tau-1$. In this case, the space of solutions is $\frac{1}{2} k(k-1)$-dimensional, spanned by

$$
\begin{gather*}
\gamma_{2,0, k-2}, \gamma_{2,1, k-3}, \cdots, \gamma_{2, k-2,0} \\
\gamma_{3,0, k-3}, \gamma_{3,1, k-4}, \cdots, \gamma_{3, k-3,0} \\
\vdots \tag{11}
\end{gather*}
$$

$$
\begin{aligned}
& \gamma_{k-1,0,1}, \gamma_{k-1,1,0} \\
& \gamma_{k, 0,0}
\end{aligned}
$$

## Cohomology of $\operatorname{Vect}\left(\mathbb{R P}^{1}\right)$ acting on $\mathcal{D}_{\lambda, \mu ; v}$

In this section, we will compute the first cohomology group of $\operatorname{Vect}\left(\mathbb{R}^{11}\right)$ with values in $\mathcal{D}_{\lambda, \mu ; j}$, vanishing on $\mathfrak{a f f}(1)$. Our main result is the following:

## Theorem 5.1

(i) For $v-\mu-\lambda \leq 11$, the space $\mathrm{H}_{\text {diff }}^{1}\left(\operatorname{Vect}\left(\mathbb{R} \mathbb{P}^{1}\right), \mathfrak{a f f}(1) ; \mathcal{D}_{\lambda, \mu ; \nu}\right)$ has the following structure:
(1) If $v-\mu-\lambda=1$, then
$\mathrm{H}_{\text {diff }}^{1}\left(\operatorname{Vect}\left(\mathbb{R P}^{\mathrm{P}}\right), \mathfrak{a f f}(1) ; \mathcal{D}_{\lambda, \mu ; ;}\right) \simeq\left\{\begin{array}{cc}\mathbb{R} & \text { if } 0.2 \mathrm{~cm}(\lambda, \mu)=(0,0), \\ 0 & \text { otherwise. }\end{array}\right.$
(2) If $v-\mu-\lambda=2$, then
$H_{\text {diff }}^{\prime}\left(\operatorname{Vect}\left(\mathbb{R} \mathbb{R}^{P}\right)\right.$, aff $\left.(1) ; \mathcal{D}_{\lambda, \mu, t)}\right) \simeq\left\{\begin{array}{ccc}0.2 \mathrm{~cm} & \mathbb{R} \text { if }(\lambda, \mu) \in\left\{(0,0),\left(0,-\frac{1}{2}\right),\left(-\frac{1}{2}, 0\right)\right\}, 0.2 \mathrm{~cm} \\ 0 & \text { otherwise. }\end{array}\right.$
(3) If $v-\mu-\lambda=3$, then
$H_{\text {diff }}^{1}\left(\operatorname{Vect}\left(\mathbb{R} \mathbb{P}^{\prime}\right), a \operatorname{aff}(1) ; \mathcal{D}_{\lambda, \mu, i}\right) \simeq\left\{\begin{array}{cl}0.2 \mathrm{~cm} & \mathbb{R} \\ \text { if }(\lambda, \mu) \in \begin{cases}(0,-1),(-1,0),\left(0,-\frac{1}{3}\right), \\ \left(-\frac{1}{3}, 0\right),\left(-\frac{1}{2},-\frac{1}{2}\right)\end{cases} \\ 0 & \text { otherwise. }\end{array}\right\}, 0.2 \mathrm{~cm}$
(4) If $v-\mu-\lambda=4$, then
$\mathrm{H}_{\text {diff }}^{1}\left(\operatorname{Vect}\left(\mathbb{R}^{\mathrm{P}}\right)\right.$, aff $\left.(1) ; \mathcal{D}_{\lambda, \mu, \psi}\right) \simeq\left\{\begin{array}{c}0.2 \mathrm{~cm} \quad \mathbb{R} \quad \text { if }(\lambda, \mu) \in\left\{\begin{array}{l}\left(-\frac{2}{3}, 0\right),\left(0,-\frac{2}{3}\right),\left(-\frac{3}{2}, 0\right),\left(0,-\frac{3}{2}\right), \\ \left(-\frac{1}{2},-1\right),\left(-1,-\frac{1}{2}\right),\left(-\frac{1}{3},-\frac{1}{3}\right)\end{array}\right\} \text { otherwise. }\end{array}\right\}$,
(5) If $v-\mu-\lambda=5$, then

(6) If $v-\mu-\lambda=6$, then

(7) If $v-\mu-\lambda=7$, then

(8) If $v-\mu-\lambda=8$, then

(9) If $v-\mu-\lambda=9$, then

(10) If $v-\mu-\lambda=10$, then

(11) If $v-\mu-\lambda=11$, then

(ii) If $\nu-\mu-\lambda$ is semi-integer but $\lambda$ and $\mu$ are generic then,

$$
\mathrm{H}_{\mathrm{diff}}^{1}\left(\operatorname{Vect}\left(\mathbb{R} \mathbb{P}^{1}\right), \mathfrak{a f f}(1) ; \mathcal{D}_{\lambda, \mu ; ;}\right)=0 .
$$

Proof of Theorem 5.1: To proof Theorem (5.1) we proceed bye following the three steps:

- We will investigate the dimension of the space of operators that satisfy the 1 -cocycle condition. By Proposition (4.4), its dimension is at most $\frac{1}{2} k(k-1)$, where $k=v-\mu-\lambda+1$, since any 1 -cocycle that vanishes on $\mathfrak{a f f}(1)$ is certainly $\mathfrak{a f f}(1)$-invariant.
- We will study all trivial 1-cocycles, namely, operators of the form $L_{x} B$,
where $B$ is a bilinear operator. As our 1 -cocycles vanish on the Lie algebra $\mathfrak{a f f}(1)$, it follows that the operator $B$ coincides with the transvectant $J_{k}^{\lambda, \mu}$.
- By taking into account Part 1 and Part 2 and depending on $\lambda$ and $\mu$ the dimension of the cohomology group $\mathrm{H}_{\text {diff }}^{1}\left(\operatorname{Vect}\left(\mathbb{R} \mathbb{P}^{1}\right)\right.$,aff $\left.(1) ; \mathcal{D}_{\lambda, \mu ; \nu}\right)$ will be equal to
$\operatorname{dim}($ operators that are $1-$ cocycles $)-\operatorname{dim}\left(\right.$ operators of the form $\left.L_{X} J_{k}^{\lambda, \mu}\right)$.
Now, clearly the coboundary $L_{X} J_{k}^{\lambda, \mu}$ has the following form:

$$
\begin{equation*}
L_{X} J_{k}^{\lambda, \mu}(X, \phi, \psi)=\sum_{i+j+l=k+1} \beta_{i, j, l} X^{(i)} \phi^{(j)} \psi^{(l)} \tag{23}
\end{equation*}
$$

where

$$
\beta_{0, j, l}=\beta_{1, j, l}=0
$$

The following Lemma is proved directly which will be useful in the proof of Theorem 5.1.

## Lemma 5.2

For $\lambda, \mu \in \mathbb{R}$

$$
\left.\left.\left.\beta_{\alpha, \beta, k-\alpha-\beta+1}=-\left({ }_{\alpha}^{\alpha+\beta-1}\right)+\lambda\left({ }_{\alpha-1}^{\alpha+\beta-1}\right)\right) \gamma_{\alpha+\beta-1, k-\alpha-\beta+1}-\left(\begin{array}{c}
\alpha_{\alpha}^{k-\beta}
\end{array}\right)+\mu_{\alpha-1}^{k-\beta}\right)\right) \gamma_{\beta, k-k-\beta},
$$

where $\alpha \geq 2$ and $\beta \geq 0$.
We need also the following Lemma.

## Lemma 5.3

Every 1-cocycle on $\operatorname{Vect}\left(\mathbb{R P}^{1}\right)$ with values in $\left.\mathcal{D}_{\lambda, \mu ; j}\right)$ is differentiable Proof [7].
Now we are in position to prove Theorem (5.1). By Lemma (5.3), any 1-cocycle on $\operatorname{Vect}\left(\mathbb{R}^{1}\right)$ should retains the following general form:

$$
\begin{equation*}
C(X, \phi, \psi)=\sum_{i+j+l=k} c_{i, j, l} X^{(i)} \phi^{(j)} \psi^{(l)}, \tag{24}
\end{equation*}
$$

where $c_{i, j, 1}$ are constants. The fact that this 1-cocycle vanishes on $\mathfrak{a f f}(1)$ implies that

$$
c_{0, j, l}=c_{1, j, l}=0 .
$$

The 1-cocycle condition reads as follows: for all $\varphi \in \mathcal{F}_{\lambda}$, for all $\psi \in \mathcal{F}_{\mu}$ and for all $X \in \operatorname{Vect}\left(\mathbb{R P}^{1}\right)$, one has

$$
c([X, Y], \phi, \psi)-L_{X}^{\lambda, \mu ; \nu} B(Y, \phi, \psi)+L_{Y}^{\lambda, \mu ; v} B(X, \phi, \psi)=0 .
$$

The case where $\boldsymbol{v}-\boldsymbol{\mu}-\lambda=\mathbf{1}$ : In this case, according to Proposition 4.4, the 1 -cocycle (24) can be expressed as follows:

$$
\Upsilon(X, \phi, \psi)=c_{2,0,0} X^{\prime \prime} \phi \psi .
$$

By a direct computation, we can see that the 1-cocycle condition is always satisfied. Let us study the triviality of this 1-cocycle. A direct computation proves that

$$
L_{X} J_{1}^{\lambda, \mu}=\beta_{2,0,0} X^{\prime \prime} \phi \psi=-\left(\lambda \gamma_{1,0}+\mu \gamma_{0,1}\right) X^{\prime \prime} \phi \psi
$$

So, for $(\lambda, \mu)=(0,0)$, the coeffcient $c_{2,0,0}$ cannot be eliminated by adding a coboundary. Hence, the cohomology space is one-dimensional. While for $(\lambda, \mu) \neq(0,0)$, we can see that the coeffcient $c_{2,0,0}$ can be eliminated because $\beta_{2,0,0} \neq 0$. Hence, the cohomology is zero-dimensional.

The case where $\boldsymbol{v}-\mu-\lambda=2$ : In this case, according to Proposition 4.4, the 1-cocycle (24) can be expressed as follows:

$$
\Upsilon(X, \phi, \psi)=c_{3,0,0} X^{\prime \prime \prime} \phi \psi+c_{2,1,0} X^{\prime \prime} \phi^{\prime} \psi+c_{2,0,1} X^{\prime \prime} \phi \psi^{\prime} .
$$

By a direct computation, we can see that the 1-cocycle condition is always satisfied. Let us study the triviality of this 1-cocycle. A direct computation proves that

$$
L_{X} J_{2}^{\lambda, \mu}=\beta_{3,0,0} X^{\prime \prime \prime} \phi \psi+\beta_{2,1,0} X^{\prime \prime} \phi^{\prime} \psi+\beta_{2,0,1} X^{\prime \prime} \phi \psi^{\prime} .
$$

where
$\beta_{3,0,0}=-\lambda \gamma_{2,0}-\mu \gamma_{0,2} ; \beta_{2,1,0}=-(2 \lambda+1) \gamma_{2,0}-\mu \gamma_{1,1}$ and $\beta_{2,0,1}=-\lambda \gamma_{1,1}-(2 \mu+1) \gamma_{0,2}$.
So, for $(\lambda, \mu)=(0,0),\left(-\frac{1}{2}, 0\right),\left(0,-\frac{1}{2}\right)$, the cohomology space is one-dimensional, since only one of the coefficients $c_{3,0,0}, c_{2,1,0}$ or $c_{2,0,1}$ cannot be eliminated by adding a coboundary. While for $(\lambda, \mu) \neq(0,0),\left(-\frac{1}{2}, 0\right),\left(0,-\frac{1}{2}\right)$, the coeffcient $c_{3,0,0}, c_{2,1,0}$ and $c_{2,0,1}$ can be eliminated because $\beta_{3,0,0}, \beta_{2,1,0}$ and $\beta_{2,0,1}$ are nonzero. Hence, the cohomology space is zero-dimensional.

The case where $v-\mu-\lambda \geq 3$ : In this case, the 1 -cocycle condition is equivalent to the system:

$$
\begin{align*}
& \left(\binom{\alpha+\beta-1}{\alpha}-\binom{\alpha+\beta-1}{\alpha-1}\right) c_{\alpha+\beta-1, \gamma, a}+\left(\binom{\alpha+\gamma-1}{\alpha}+\lambda\binom{\alpha+\gamma-1}{\alpha-1}\right) c_{\beta, \alpha+\gamma-1, a} \\
& -\left(\binom{\beta+\gamma-1}{\beta}+\lambda\binom{\beta+\gamma-1}{\beta-1}\right) c_{\alpha, \beta+\gamma-1, a}+\left(\binom{\alpha+a-1}{\alpha}+\mu\binom{\alpha+a-1}{\alpha-1}\right) c_{\beta, \gamma, \alpha+\alpha-1}  \tag{25}\\
& -\left(\binom{\beta+a-1}{\beta}+\lambda\binom{\beta+a-1}{\beta-1}\right) c_{\alpha,,, \beta+\alpha-1}=0,
\end{align*}
$$

where $\alpha+\beta+\gamma+a=k+1, \alpha>\beta \geq 2, \alpha>\gamma$ and $\alpha>a$, obtained from the coefficient of $X^{(\alpha)} Y^{(\beta)(\gamma)(a)}$.

This system can be deduced by a simple computation. Of course, such a system has at least one solution in which the solutions $c_{i, j, l}$ are just the coefficients $\beta_{i, j, l}$ of the coboundaries (23).

The case where $\boldsymbol{v}-\boldsymbol{\mu}-\lambda=3$ : In this case, according to Proposition 4.4, the space of solutions is spanned by:

$$
c_{4,0,0,} c_{3,1,0,} c_{3,0,1,} c_{2,2,0,} c_{2,1,1,1} c_{2,0,2}
$$

Moreover, by formula (25), we readily obtain:
$-2 c_{2,0,0}+\lambda c_{2,0,0} \lambda c_{2,0,0}+\mu c_{2,0,0}-\mu c_{2,0,0}=0$.
Thus, we have just proved that the coefficients of every 1-cocycle is expressed in terms of

$$
c_{3,1,0,} c_{3,0,1,} c_{2,2,0,} c_{2,1,1,1} c_{2,0,2}
$$

A direct computation proves that

$$
\begin{aligned}
L_{X} J_{3}^{\lambda, \mu} & =\beta_{2,1,1} X^{\prime \prime} \phi^{\prime} \psi^{\prime}+\beta_{2,2,0}\left(\frac{\lambda}{2} X^{4} \phi \psi+X^{\prime \prime} \phi^{\prime \prime} \psi\right)+\beta_{2,0,2}\left(\frac{\mu}{2} X^{(4)} \phi \psi+X^{\prime \prime} \phi \psi^{\prime \prime}\right) \\
& +\beta_{3,1,0}\left(-\frac{\lambda}{2} X^{(4)} \phi \psi+X^{\prime \prime \prime} \phi^{\prime} \psi\right)+\beta_{3,0,1}\left(-\frac{\mu}{2} X^{(4)} \phi \psi+X^{\prime \prime \prime} \phi \psi^{\prime}\right) .
\end{aligned}
$$

where

$$
\begin{array}{lll}
\beta_{3,1,0}=-(3 \lambda+1) \gamma_{3,0}-\mu \gamma_{1,2} & , & \beta_{3,0,1}=-\lambda \gamma_{2,1}-(3 \mu+1) \gamma_{0,3} \\
\beta_{2,2,0}=-3(\lambda+1) \gamma_{3,0}-\mu \gamma_{2,1} & , & \beta_{2,0,2}=-\lambda \gamma_{1,2}-3(\mu+1) \gamma_{0,3} \\
\beta_{2,1,1}=-(2 \lambda+1) \gamma_{2,1}-(2 \mu+1) \gamma_{1,2} &
\end{array}
$$

So, for $(\lambda, \mu)=\left(-\frac{1}{3}, 0\right),\left(0,-\frac{1}{3}\right),(-1,0),(0,-1),\left(-\frac{1}{2},-\frac{1}{2}\right)$, the cohomology space is one-dimensional, since only one of the coefficients $c_{3,1,0}$, $c_{3,0,1}, c_{2,2,0}, c_{2,1,1}$ or $c_{2,0,2}$ cannot be eliminated by adding a coboundary. While for $(\lambda, \mu) \neq\left(-\frac{1}{3}, 0\right),\left(0,-\frac{1}{3}\right),(-1,0),(0,-1),\left(-\frac{1}{2},-\frac{1}{2}\right)$, the coeffcient $c_{3,1,0}, c_{3,0,1}, c_{2,2,0}, c_{2,1,1}$ and $c_{2,0,2}$ can be eliminated because $\beta_{3,1,0}, \beta_{3,0,1}, \beta_{2,2,0,}$, $\beta_{2,1,1}^{3,1}$ and $\beta_{2,0,2}^{2,2,0}$ are nonzero. Hence, the cohomology space is zerodimensional.

The case where $\boldsymbol{v}-\boldsymbol{\mu}-\boldsymbol{\lambda}=4$ : In this case, according to Proposition 4.4, the space of solutions is spanned by:

$$
c_{5,0,0}, c_{4,1,0}, c_{4,0,1}, c_{3,2,0}, c_{3,1,1}, c_{3,0,2}, c_{2,3,0}, c_{2,2,1}, c_{2,1,2}, c_{2,0,3}
$$

Moreover, by formula (25), we readily obtain:

$$
\begin{aligned}
& -2 c_{4,1,0}+(3 \lambda+1) c_{2,3,0}-(2 \lambda+1) c_{3,2,0}+\mu c_{2,1,2}-\mu c_{3,1,1}=0 \\
& -2 c_{4,0,1}+\lambda c_{2,2,1}-\lambda c_{3,1,1}+(3 \mu+1) c_{2,0,3}-(2 \mu+1) c_{3,0,2}=0 \\
& -5 c_{5,0,0}+\lambda c_{2,3,0}-\lambda c_{4,1,0}+\mu c_{2,0,3}-\mu c_{4,0,1}=0
\end{aligned}
$$

Thus, we have just proved that the coefficients of every 1-cocycle is expressed in terms of

$$
c_{3,2,0}, c_{3,1,1}, c_{3,0,2}, c_{2,3,0}, c_{2,2,1}, c_{2,1,2}, c_{2,0,3}
$$

A direct computation confirms that, the coefficients of $L_{X} J_{4}^{\lambda, \mu}$ are expressed in terms of:

$$
\begin{array}{lll}
\beta_{2,2,1}=-3(\lambda+1) \gamma_{3,1}-(2 \mu+1) \gamma_{2,2} & , & \beta_{2,1,2}=-(2 \lambda+1) \gamma_{2,2}-3(\mu+1) \gamma_{1,3} \\
\beta_{3,2,0}=-2(3 \lambda+2) \gamma_{4,0}-\mu \gamma_{2,2} & , & \beta_{3,0,2}=-\lambda \gamma_{2,2}-2(3 \mu+2) \gamma_{0,4}, \\
\beta_{2,3,0}=-2(2 \lambda+3) \gamma_{4,0}-\mu \gamma_{3,1} & , & \beta_{2,0,3}=-\lambda \gamma_{1,3}-2(2 \mu+3) \gamma_{0,4} \\
\beta_{3,1,1}=-(3 \lambda+1) \gamma_{3,1}-(3 \mu+1) \gamma_{1,3} . & &
\end{array}
$$

So, for $\quad(\lambda, \mu)=\left(-\frac{2}{3}, 0\right),\left(0,-\frac{2}{3}\right),\left(-\frac{3}{2}, 0\right),\left(0,-\frac{3}{2}\right),\left(-\frac{1}{2},-1\right),\left(-1,-\frac{1}{2}\right),\left(-\frac{1}{3},-\frac{1}{3}\right)$, the cohomology space is one-dimensional, since only one of the coefficients $c_{3,2,0}, c_{3,1,1}, c_{3,0,2}, c_{2,3,0}, c_{2,2,1}, c_{2,1,2}$ or $c_{2,0,3}$ cannot be eliminated by adding a coboundary. While for $(\lambda, \mu) \neq\left(-\frac{2}{3}, 0\right),\left(0,-\frac{2}{3}\right),\left(-\frac{3}{2}, 0\right), \quad\left(0,-\frac{3}{2}\right),\left(-\frac{1}{2},-1\right),\left(-1,-\frac{1}{2}\right),\left(-\frac{1}{3},-\frac{1}{3}\right), \quad$ the coeffcient $c_{3,2,0}, c_{3,1,1}, c_{3,0,2}, c_{2,3,0}, c_{2,2,1}, c_{2,1,2}$ and $c_{2,0,3}$ can be eliminated because $\beta_{3,2,0}, \beta_{3,1,1}, \beta_{3,0,2}, \beta_{2,3,0}, \beta_{2,2,1}, \beta_{2,1,2}$ and $\beta_{2,0,3}$ are nonzero. Hence, the cohomology space is zero-dimensional.

The case where $\boldsymbol{v}-\boldsymbol{\mu}-\boldsymbol{\lambda}=5$ : In this case, according to Proposition 4.4, the space of solutions is spanned by:

$$
\begin{aligned}
& c_{6,0,0}, c_{5,1,0}, c_{5,0,1}, c_{4,2,0}, c_{4,1,1}, c_{4,0,2}, c_{3,3,0}, c_{3,2,1} \\
& c_{3,1,2}, c_{3,0,3}, c_{2,4,0}, c_{2,3,1}, c_{2,2,2}, c_{2,1,3}, c_{2,0,4}
\end{aligned}
$$

Moreover, by formula (25), we readily obtain:
$-2 c_{4,1,1}+(3 \lambda+1) c_{2,3,1}-(2 \lambda+1) c_{3,2,1}+(3 \mu+1) c_{2,1,3}-(2 \mu+1) c_{3,1,2}=0$,
$-5 c_{5,1,0}+(4 \lambda+1) c_{2,4,0}-(2 \lambda+1) c_{4,2,0}+\mu c_{2,1,3}-\mu c_{4,1,1}=0$,
$-5 c_{5,0,1}+\lambda c_{2,3,1}-\lambda c_{4,1,1}+(4 \mu+1) c_{2,0,4}-(2 \mu+1) c_{4,0,2}=0$,
$-2 c_{4,2,0}+2(3 \lambda+2) c_{2,4,0}-3(\lambda+1) c_{3,3,0}+\mu c_{2,2,2}-\mu c_{3,2,1}=0$,
$-2 c_{4,0,2}+\lambda c_{2,2,2}-\lambda c_{3,1,2}+2(3 \mu+2) c_{2,0,4}-3(\mu+1) c_{3,0,3}=0$,
$-9 c_{6,0,0}+\lambda c_{2,4,0}-\lambda c_{5,1,0}+\mu c_{2,0,4}-\mu c_{5,0,1}=0$,
$-5 c_{6,0,0}+\lambda c_{3,3,0}-\lambda c_{4,2,0}+\mu c_{3,0,3}-\mu c_{4,0,2}=0$.

Thus, we have just proved that the coefficients of every 1 -cocycle is expressed in terms of

$$
c_{2,4,0}, c_{2,0,4}, c_{3,3,0}, c_{3,0,3}, c_{3,2,1}, c_{3,1,2}, c_{2,3,1}, c_{2,2,2} .
$$

A direct computation confirms that, the coefficients of $L_{X} J_{5}^{\lambda, \mu}$ are expressed in terms of:

$$
\begin{array}{ll}
\beta_{3,2,1}=-2(3 \lambda+2) \gamma_{4,1}-(3 \mu+1) \gamma_{2,3} & , \quad \beta_{3,1,2}=-(3 \lambda+1) \gamma_{3,2}-2(3 \mu+2) \gamma_{1,4} \\
\beta_{2,3,1}=-2(2 \lambda+3) \gamma_{4,1}-(2 \mu+1) \gamma_{3,2} & , \quad \beta_{2,2,2}=-3(\lambda+1) \gamma_{3,2}-3(\mu+1) \gamma_{2,3} \\
\beta_{3,3,0}=-10(\lambda+1) \gamma_{5,0}-\mu \gamma_{3,2} & , \quad \beta_{3,0,3}=-\lambda \gamma_{2,3}-10(\mu+1) \gamma_{0,5} \\
\beta_{2,4,0}=-5(\lambda+2) \gamma_{5,0}-\mu \gamma_{4,1} & , \quad \beta_{2,0,4}=-\lambda \gamma_{1,4}-5(\mu+2) \gamma_{0,5}
\end{array}
$$

So, for $\quad(\lambda, \mu)=\left(-\frac{2}{3},-\frac{1}{3}\right),\left(-\frac{1}{3},-\frac{2}{3}\right),\left(-\frac{3}{2},-\frac{1}{2}\right),(-1,-1),(-1,0),(0,-1),(-2,0),(0,-2)$, the cohomology space is one-dimensional, since only one of the coefficients $\quad c_{2,4,0}, c_{2,0,4}, c_{3,3,0}, c_{3,0,3}, c_{3,2,1}, c_{3,1,2}, c_{2,3,1}$ or $c_{2,2,2}$ cannot be eliminated by adding a coboundary. While for $(\lambda, \mu) \neq\left(-\frac{2}{3},-\frac{1}{3}\right),\left(-\frac{1}{3},-\frac{2}{3}\right),\left(-\frac{3}{2},-\frac{1}{2}\right),(-1,-1),(-1,0),(0,-1),(-2,0),(0,-2)$, the coeffcient $c_{2,4,0}, c_{2,0,4}, c_{3,3,0}, c_{3,0,3}, c_{3,2,1}, c_{3,1,2}, c_{2,3,1}$ and $c_{2,2,2}$ can be eliminatedbecause $\beta_{2,4,0}, \beta_{2,0,4}, \beta_{3,3,0}, \beta_{3,0,3}, \beta_{3,2,1}, \beta_{3,1,2}, \beta_{2,3,1}$ and $\beta_{2,2,2}$ are nonzero. Hence, the cohomology space is zero-dimensional.

The case where $\boldsymbol{v}-\boldsymbol{\mu}-\boldsymbol{\lambda}=\mathbf{6}$ : In this case, according to Proposition 4.4 together with formulas (25), we check that the coefficients of every 1-cocycle are expressed in terms of

$$
c_{4,2,1}, c_{4,0,3}, c_{3,2,2}, c_{3,1,3}, c_{3,0,4}, c_{2,5,0}, c_{2,4,1}, c_{2,3,2}, c_{2,1,4}
$$

A direct computation confirms that, the coefficients of $L_{X} J_{6}^{\lambda, \mu}$ are expressed in terms of:

$$
\beta_{4,2,1}, \beta_{4,0,3}, \beta_{3,2,2}, \beta_{3,1,3}, \beta_{3,0,4}, \beta_{2,5,0}, \beta_{2,4,1}, \beta_{2,3,2}, \beta_{2,1,4}
$$

So, in the same way as before, by Lemma 5.2 , we can see, with the help of the maple, that the cohomology space is given as in (17).

The case where $\boldsymbol{v}-\boldsymbol{\mu}-\lambda=7$ : In this case, according to Proposition 4.4 together with formulas (25), we check that the coefficients of every 1-cocycle are expressed in terms of

$$
c_{8,0,0}, c_{7,1,0}, c_{7,0,1}, c_{5,3,0}, c_{5,0,3}, c_{4,0,4}, c_{2,4,2}, c_{2,1,5}, c_{2,0,6} .
$$

A direct computation confirms that, the coefficients of $L_{X} J_{7}^{\lambda, \mu}$ are expressed in terms of:

$$
\beta_{8,0,0}, \beta_{7,1,0}, \beta_{7,0,1}, \beta_{5,3,0}, \beta_{5,0,3}, \beta_{4,0,4}, \beta_{2,4,2}, \beta_{2,1,5}, \beta_{2,0,6}
$$

So, in the same way as before, by Lemma 5.2 , we can see, with the help of the maple, that the cohomology space is given as in (18).

The case where $\boldsymbol{v}-\boldsymbol{\mu}-\lambda=8$ : In this case, according to Proposition 4.4 together with formulas (25), we check that the coefficients of every 1-cocycle are expressed in terms of:

$$
c_{9,0,0}, c_{8,0,1}, c_{7,2,0}, c_{7,0,2}, c_{6,3,0}, c_{6,0,3}, c_{5,3,1}, c_{3,0,6}, c_{2,0,7}
$$

A direct computation confirms that, the coefficients of $L_{X} J_{8}^{\lambda, \mu}$ are expressed in terms of:

$$
\beta_{9,0,0}, \beta_{8,0,1}, \beta_{7,2,0}, \beta_{7,0,2}, \beta_{6,3,0}, \beta_{6,0,3}, \beta_{5,3,1}, \beta_{3,0,6}, \beta_{2,0,7}
$$

So, in the same way as before, by Lemma 5.2, we can see, with the help of the maple, that the cohomology space is given as in (19).

The case where $\boldsymbol{v}-\boldsymbol{\mu}-\lambda=\mathbf{9}$ : In this case, according to Proposition 4.4 together with formulas (25), we check that the coefficients of every 1-cocycle are expressed in terms of

## $c_{10,0,0}, c_{9,0,1}, c_{8,2,0}, c_{6,4,0}, c_{6,0,4}, c_{5,3,2}, c_{3,0,7}, c_{2,0,8}$.

A direct computation confirms that, the coefficients of $L_{X} J_{9}^{\lambda, \mu}$ are expressed in terms of:

$$
\beta_{10,0,0}, \beta_{9,0,1}, \beta_{8,2,0}, \beta_{6,4,0}, \beta_{6,0,4}, \beta_{5,3,2}, \beta_{3,0,7}, \beta_{2,0,8} .
$$

So, in the same way as before, by Lemma 5.2 , we can see, with the help of the maple, that the cohomology space is given as in (20).

The case where $\boldsymbol{v}-\boldsymbol{\mu}-\boldsymbol{\lambda = 9}$ : In this case, according to Proposition 4.4 together with formulas (25), we check that the coefficients of every 1-cocycle are expressed in terms of

$$
c_{11,0,0}, c_{10,0,1}, c_{9,2,0}, c_{6,5,0}, c_{6,0,5}, c_{5,3,3}, c_{3,0,8}, c_{2,0,9}
$$

A direct computation confirms that, the coefficients of $L_{X} J_{10}^{\lambda, \mu}$ are expressed in terms of:

$$
\beta_{11,0,0}, \beta_{10,0,1}, \beta_{9,2,0}, \beta_{6,5,0}, \beta_{6,0,5}, \beta_{5,3,3}, \beta_{3,0,8}, \beta_{2,0,9}
$$

So, in the same way as before, by Lemma 5.2, we can see, with the help of the maple, that the cohomology space is given as in (21).

The case where $\boldsymbol{v}-\boldsymbol{\mu}-\boldsymbol{\lambda}=11$ : In this case, according to Proposition 4.4 together with formulas (25), we check that the coefficients of every 1-cocycle are expressed in terms of

$$
c_{12,0,0}, c_{11,0,1}, c_{10,0,2}, c_{9,0,3}, c_{7,0,5}, c_{5,7,0}, c_{5,0,7}, c_{2,1,9}
$$

A direct computation confirms that, the coefficients of $L_{X} J_{11}^{\lambda, \mu}$ are expressed in terms of:

$$
\beta_{12,0,0}, \beta_{11,0,1}, \beta_{10,0,2}, \beta_{9,0,3}, \beta_{7,0,5}, \beta_{5,7,0}, \beta_{5,0,7}, \beta_{2,1,9} .
$$

So, in the same way as before, by Lemma 5.2, we can see, with the help of the maple, that the cohomology space is given as in (22). This completes the proof.

## Conjecture 5.1

For $v-\mu-\lambda \in \mathbb{N}+12, \lambda$ and $\mu$ are generic, one hase

$$
\mathrm{H}_{\mathrm{diff}}^{1}\left(\operatorname{Vect}\left(\mathbb{R P}^{1}\right), \mathfrak{a f f}(1) ; \mathcal{D}_{\lambda, \mu ; \nu}\right)=0 .
$$

## Conclusion

In this paper, we classify $\mathfrak{a f f}(1)$-invariant linear differential operators from $\operatorname{Vect}\left(\mathbb{R P}^{1}\right)$ to $\mathcal{D}_{\mu, j v}$ vanishing on $\mathfrak{a f f}(1)$, where $\mathcal{D}_{, \mu ; j}:=\operatorname{Homdiff}\left(\mathcal{F}_{\lambda} \otimes \mathcal{F} ; \mathcal{F}_{\nu}\right)$ is the space of bilinear differential operators acting on weighted densities. This result allows us to compute the first differential $\mathfrak{a f f}(1)$-relative cohomology of $\operatorname{Vect}\left(\mathbb{R} \mathbb{P}^{1}\right)$ with coefficients in $\mathcal{D}_{\lambda, \mu, j,}$.

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