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#### **Research Article**

# On the First $\mathfrak{aff}(1)\text{-Relative}$ Cohomology of the Lie Algebra of Vector Fields and Differential Operators

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#### Abstract

Let  $Vect(\mathbb{RP}^1)$  be the Lie algebra of smooth vector fields on  $\mathbb{RP}^1$ . In this paper, we classify  $\mathfrak{aff}(1)$ -invariant linear differential operators from  $Vect(\mathbb{RP}^1)$  to  $\mathcal{D}_{\lambda,\mu\nu}$  vanishing on  $\mathfrak{aff}(1)$ , where  $\mathcal{D}_{\lambda,\mu\nu}$ :=Homdiff( $\mathcal{F}_{\lambda} \otimes \mathcal{F}_{\mu}; \mathcal{F}_{\nu}$ ) is the space of bilinear differential operators acting on weighted densities. This result allows us to compute the first differential  $\mathfrak{aff}(1)$ -relative cohomology of  $Vect(\mathbb{RP}^1)$  with coefficients in  $\mathcal{D}_{\lambda,\mu\nu}$ .

Keywords: Differential operators; Transvectants; Lie algebra; Cohomology

#### Introduction

Let  $\mathfrak{g}$  be a Lie algebra and let  $\mathcal M$  and  $\mathcal N$  be two  $\mathfrak{g}\text{-modules}.$  It is well-known that nontrivial extensions of  $\mathfrak{g}\text{-modules}:$ 

 $0{\rightarrow}\mathcal{M}{\rightarrow}.{\rightarrow}\mathcal{N}{\rightarrow}0$ 

are classified by the first cohomology group H<sup>1</sup>( $\mathfrak{g}$ ; Hom( $\mathcal{N},\mathcal{M}$ )) [1]. Any 1-cocycle  $\mathcal{L}$  generates a new action on  $\mathcal{M}\oplus\mathcal{N}$  as follows: for all  $g\in\mathfrak{g}$  and for all  $(a,b)\in\mathcal{M}\oplus\mathcal{N}$ , we define  $g^{*}(a,b):=(g^{*}a+C^{*t}\mathcal{L}(b),g^{*}b)$ . For the space of tensor density of weight  $\lambda$ ,  $\mathcal{F}_{\lambda}$ , viewed as a module over the Lie algebra of smooth vector fields Vect( $\mathbb{RP}^{1}$ ), the classification of nontrivial extensions

 $0 \rightarrow \mathcal{F}_{\mu} \rightarrow . \rightarrow \mathcal{F}_{\lambda} \rightarrow 0,$ 

leads Feigin and Fuks [2] to compute the cohomology group  $H^1(\text{Vect}(\mathbb{RP}^1); \text{Hom}(\mathcal{F}_{\lambda}, \mathcal{F}_{\mu}))$ . Later, Ovsienko and Bouarroudj [3] have computed the corresponding relative cohomology group with respect to  $\mathfrak{sl}(2, \mathbb{R})$ , namely

H<sup>1</sup>(Vect( $\mathbb{RP}^1$ ),  $\mathfrak{sl}(2, \mathbb{R})$ ;Hom( $\mathcal{F}_{\lambda}, \mathcal{F}_{\mu}$ )).

In this paper, we will compute the first cohomology group

H<sup>1</sup>(Vect( $\mathbb{RP}^1$ ), $\mathfrak{aff}(1)$ ;Hom( $\mathcal{F}_{\lambda}\otimes,\mathcal{F}_{\mu},\mathcal{F}_{\nu}$ )).

# $Vect(\mathbb{R})$ -Module Structures on the Space of Bilinear Differential Operators

Consider the standard (local) action of  $\mathfrak{aff}(1)$  on  $\mathbb R$  by linear-fractional transformations. Although the action is local, it generates global vector fields

$$\{\frac{d}{dx}, x\frac{d}{dx}\},\$$

that form a Lie subalgebra of  $Vect(\mathbb{R})$  isomorphic to the Lie algebra  $\mathfrak{aff}(1)$ . This realization of  $\mathfrak{aff}(1)$  is understood throughout this paper.

#### The space of tensor densities on $\mathbb{RP}^1$

The space of tensor densities of weight  $\lambda$  (or  $\lambda$ -densities) on  $\mathbb{RP}^1$ , denoted by:

 $\mathcal{F}_{\lambda} = \{ f(dx)^{\lambda} | f \in C^{\infty}(\mathbb{R}) \}, \, \lambda \in \mathbb{R},$ 

is the space of sections of the line bundle  $(T^*\mathbb{RP}^1)^{\otimes^{\lambda}}$ . This space coincides with the space of functions and differential forms for  $\lambda=0$  and for  $\lambda=1$ , respectively. The Lie algebra  $\operatorname{Vect}(\mathbb{RP}^1)$  acts on  $\mathcal{F}_{\lambda}$  by the Lie derivative. For all  $X \in \operatorname{Vect}(\mathbb{RP}^1)$  and for all  $\varphi \in \mathcal{F}_{\lambda}$ :

 $L_{x}(\varphi(dx)^{\lambda}) = X\varphi' + \lambda\varphi X', \qquad (1)$ 

where the superscript ' stands for d/dx.

## The space of bilinear differential operators as a $Vect(\mathbb{RP}^1)$ -module

We are interested in defining a three-parameter family of  $Vect(\mathbb{RP}^1)$ -modules on the space of bilinear differential operators. The counterpart  $Vect(\mathbb{RP}^1)$ -modules of the space of linear differential operators is a classical object [4].

Consider bilinear differential operators that act on tensor densities:

$$A:\mathcal{F}_{\lambda}\otimes\mathcal{F}_{\mu}\to\mathcal{F}_{\nu} \tag{2}$$

The Generalized Lie algebra Vect( $\mathbb{RP}^1$ ) acts on the space of bilinear differential operators as follows. For all  $\varphi \in \mathcal{F}_1$  and for all  $\psi \in \mathcal{F}_n$ :

$$L_{\chi}^{\lambda,\mu;\nu}(A)(\phi,\psi) = L_{\chi}^{\nu} \circ A(\phi,\psi) - A(L_{\chi}^{\lambda}\phi,\psi) - A(\phi,L_{\chi}^{\mu}\psi)$$
(3)

where  $L^{\lambda}_{\chi}$  is the action (1). We denote by  $\mathcal{D}_{\lambda\mu\nu\nu}$  the space of bilinear differential operators (2) endowed with the defined Vect( $\mathbb{RP}^1$ )-module structure (3).

#### **Relative Cohomology**

Let us first recall some fundamental concepts from cohomology theory [1]. Let g be a Lie algebra acting on a vector space V and let h be a sub- algebra of g. (If h is omitted it assumed to be {0}.) The space of h-relative n-cochains of g with values in V is the g-module

 $C^n(g,h;V)$ : =Hom<sub>*h*</sub>( $\Lambda^n(g/h);V$ )

The coboundary operator  $\delta_n: C^n(g,h;V) \to C^{n+1}(g,h;V)$  is a g-map satisfying  ${}_n \circ \delta_{n1} = 0$ . The kernel of  $\delta_n$ , denoted  $Z^n(g,h;V)$ , is the space of *h*-relative *n*- cocycles, among them, the elements in the range of  $\delta_{n-1}$  are called *h*-relative *n*- coboundaries. We denote  $B^n(g,h;V)$  the space of *n*-coboundaries.

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By definition, the  $n^{th}$  h-relative cohomolgy space is the quotient space

 $H^{n}(\mathfrak{g},h;V)=Z^{n}(\mathfrak{g},h;V)/B^{n}(\mathfrak{g},h;V).$ 

We will only need the formula of  $\delta_{\mu}$  (which will be simply denoted  $\delta$ ) in degrees 0,1 and 2: for  $v \in C^0(g,h;V) = V^h$ ,  $\delta v(g) := (-1)^{|g| \cdot v} g.v$ , where

 $V^{h} = \{v \in V | h.v = 0 \text{ for all } h \in h\},\$ 

and for  $\Upsilon \in C^1(g,h;V)$ ,

 $\delta(\Upsilon)(x,y) := x \cdot \Upsilon(y) - y \cdot \Upsilon(x) - \Upsilon([x,y])$  for any  $x,y \in g$ .

#### aff(1)-Invariant Differential Operators

The following steps to compute the relative cohomology has intensively been used in refs. [3,5-8]. First, we classify aff(1)-invariant differential operators, then we isolate among them those that are 1-cocycles. To do that, we need the following Lemma.

#### Lemma 4.1

Any 1-cocycle vanishing on the subalgebra  $\mathfrak{aff}(1)$  of  $\operatorname{Vect}(\mathbb{R})$  is aff(1)-invariant.

The 1-cocycle condition of  $\Upsilon$  reads:

$$X \cdot \Upsilon(Y) - Y \cdot \Upsilon(X) - \Upsilon([X, Y]) = 0, \tag{4}$$

where  $X, Y \in \text{Vect}(\mathbb{RP}^1)$ . Thus, if  $\Upsilon(X)=0$  for all  $X \in aff(1)$ , eqn. (4) becomes

 $\Upsilon([X,Y]) = X \cdot \Upsilon(Y)$ 

expressing the  $\mathfrak{aff}(1)$ -invariance property of  $\Upsilon$ .

As our 1-cocycles vanish on aff(1), we will investigate aff(1)invariant linear differential operators that vanish on aff(1).

Proposition 4.2: There exist aff(1)-invariant bilinear differential operators  $J_k^{\lambda,\mu}$ :  $\mathcal{F}_{\lambda} \otimes \mathcal{F}_{\mu} \to \mathcal{F}_{\lambda+\mu+k}$  given by:

$$J_{k}^{\lambda,\mu}(\varphi dx^{\lambda}, \phi dx^{\mu}) = \sum_{i+j=k} \gamma_{i,j} \varphi^{(i)} \phi^{(j)} dx^{\lambda+\mu+k}$$
(5)

where  $k \in \mathbb{N}$  and the coefficients  $\gamma_{i,i}$  are constants.

*Proof.* Any differential operator  $J_k^{\lambda,\mu}$ :  $\mathcal{F}_{\lambda} \otimes \mathcal{F}_{\mu} \to \mathcal{F}_{\nu}$  is of the form

$$J_k^{\lambda,\mu}(fdx^{\lambda},gdx^{\mu}) = \sum_{n=0}^m \sum_{i+j=k} \gamma_{i,j} f^{(i)} g^{(j)} dx^{\nu}, \quad m \in \mathbb{N}$$

The osp(1|2) -invariant property of the operators  $J_k^{\lambda,\mu}$  with respect to the vector field  $X = x \frac{d}{dx}$  yields:

$$\frac{d}{dx}\gamma_{i,j} = 0$$
 and  $v - \lambda - \mu = k$  with  $k = i + j$ .

So, we see that the corresponding operator can be expressed as (5).

Proposition 4.3: There exist aff(1)-invariant trilinear differential operators  $K_k^{\tau,\lambda,\mu}: \mathcal{F}_\tau \otimes \mathcal{F}_\lambda \otimes \mathcal{F}_\mu \to \mathcal{F}_{\tau+\lambda+\mu+k}$  given by:

$$K_{k}^{\tau,\lambda,\mu}(\varphi,\phi,\psi) = \sum_{i+j+l=k} \gamma_{i,j,l} \varphi^{(i)} \phi^{(j)} \psi^{(l)}.$$
(6)

where i+j+l=k and the coefficients  $\gamma_{i,j,l}$  are constants.

If  $\tau$ ,  $\lambda$  and  $\mu$  are generic, then the space of solutions is  $\frac{1}{2}(k+1)(k+2)$ -dimensional.

**Proposition 4.4:** There exist aff(1)-invariant trilinear differential

operators  $K_k^{\lambda,\mu}$ : Vect $(\mathbb{RP}^1) \otimes \mathcal{F}_{\lambda} \otimes \mathcal{F}_{\mu} \to \mathcal{F}_{\lambda+\mu+k-1}$  that vanishe on aff(1) given by:

$$K_{k}^{\lambda,\mu}(X,\phi,\psi) = \sum_{i+j+l=k} \gamma_{i,j,l} X^{(i)} \phi^{(j)} \psi^{(l)}.$$
(7)

where i+j+l=k and the coefficients  $\gamma_{i,j,l}$  are constants but  $\gamma_{0,j,k-j}=\gamma_{1,j,k-j}$ =0. Moreover, the space of solutions is  $\frac{1}{2}k(k-1)$  -dimensional, for all  $\lambda$  and  $\mu$ .

Proof of Proposition 4.3 and 4.4: We are going to prove Proposition 4.3 and 4.4 simultaneously. Any differential operator  $K_k^{\tau,\lambda,\mu}: \mathcal{F}_\tau \otimes \mathcal{F}_\lambda \otimes \mathcal{F}_\mu \to \mathcal{F}_{\tau+\lambda+\mu+k}$  is of the form

$$K_k^{\tau,\lambda,\mu}(\varphi,\phi,\psi) = \sum_{i+j+l=k} \gamma_{i,j,l} \varphi^{(i)} \phi^{(j)} \psi^{(l)}.$$
(8)

where  $\gamma_{i,il}$  are functions. The  $\mathfrak{aff}(1)$  -invariant property of the operators  $K_k^{\tau,\lambda,\mu}$  reads as follows.

$$L_{X}^{\nu}K_{k}^{\tau,\lambda,\mu}(\phi,\phi,\psi,) = K_{k}^{\tau,\lambda,\mu}(L_{X}^{\tau}\phi,\phi,\psi) + K_{k}^{\tau,\lambda,\mu}(\phi,L_{X}^{\lambda}\phi,\psi) + K_{k}^{\tau,\lambda,\mu}(\phi,\phi,L_{X}^{\mu}\psi).$$
(9)

The invariant property with respect to the vector field  $X = \frac{a}{1}$ implies that  $\gamma'_{i,j,l} = 0$ . On the other hand, the invariant property with respect to the vector fields  $X = x \frac{d}{dx}$  implies that  $v = \tau + \lambda + \mu + k$ . If  $\tau$ ,  $\lambda$  and  $\mu$  are generic, then the space of solutions is  $\frac{1}{2}(k+1)(k+2)$ -dimensional, spanned by

Now, the proof of Proposition 4.4 follows as above by putting  $\tau$ -1. In this case, the space of solutions is  $\frac{1}{2}k(k-1)$  -dimensional, spanned by

$$\begin{array}{l}
\gamma_{2,0,k-2}, \gamma_{2,1,k-3}, \cdots, \gamma_{2,k-2,0}, \\
\gamma_{3,0,k-3}, \gamma_{3,1,k-4}, \cdots, \gamma_{3,k-3,0}, \\
\vdots \\
\gamma_{k-1,0,1}, \gamma_{k-1,1,0}, \\
\gamma_{k,0,0}.
\end{array}$$
(11)

#### Cohomology of Vect( $\mathbb{RP}^1$ ) acting on $\mathcal{D}_{\lambda,\mu;\nu}$

In this section, we will compute the first cohomology group of Vect( $\mathbb{RP}^1$ ) with values in  $\mathcal{D}_{\lambda uv}$ , vanishing on  $\mathfrak{aff}(1)$ . Our main result is the following:

#### Theorem 5.1

(i) For  $\nu - \mu - \lambda \le 11$ , the space  $H^1_{diff}(\text{Vect}(\mathbb{RP}^1), \mathfrak{aff}(1); \mathcal{D}_{\lambda, u; \nu})$  has the following structure:

(1) If  $v - \mu - \lambda = 1$ , then

$$H^{1}_{diff}(\operatorname{Vect}(\mathbb{RP}^{1}),\mathfrak{aff}(1);\mathcal{D}_{\lambda,\mu,\nu}) \simeq \begin{cases} \mathbb{R} & \text{if } 0.2cm(\lambda,\mu) = (0,0), \\ 0 & otherwise. \end{cases}$$
(12)

$$H^{1}_{diff}(\operatorname{Vect}(\mathbb{RP}^{1}), \mathfrak{aff}(1); \mathcal{D}_{\lambda,\mu\nu}) \simeq \begin{cases} 0.2cm \quad \mathbb{R} \quad \text{if } (\lambda, \mu) \in \{(0,0), (0, -\frac{1}{2}), (-\frac{1}{2}, 0)\}, 0.2cm \\ 0 \quad otherwise. \end{cases}$$
(13)

(3) If  $\nu - \mu - \lambda = 3$ , then

(2) If  $\nu - \mu - \lambda = 2$ , then

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$$H_{diff}^{1}(\operatorname{Vect}(\mathbb{RP}^{1}), \mathfrak{aff}(1); \mathcal{D}_{\lambda,\mu,\nu}) \simeq \begin{cases} 0.2cm \quad \mathbb{R} \quad \text{if } (\lambda,\mu) \in \left\{ (0,-1), (-1,0), (0,-\frac{1}{3}), \\ (-\frac{1}{3}, 0), (-\frac{1}{2},-\frac{1}{2}) \\ 0 \quad otherwise \end{cases} \right\}, 0.2cm \quad (14)$$

(4) If  $\nu - \mu - \lambda = 4$ , then

$$H^{1}_{\text{deff}}(\text{Vect}(\mathbb{RP}^{1}),\mathfrak{aff}(1);\mathcal{D}_{\lambda,\mu^{\gamma}}) \simeq \begin{cases} 0.2cm \quad \mathbb{R} \quad \text{if } (\lambda,\mu) \in \begin{cases} (-\frac{2}{3},0),(0,-\frac{2}{3}),(-\frac{3}{2},0),(0,-\frac{3}{2}), \\ (-\frac{1}{2},-1),(-1,-\frac{1}{2}),(-\frac{1}{3},-\frac{1}{3}) \end{cases} \end{cases}, \quad (15)$$

#### (5) If $v - \mu - \lambda = 5$ , then

$$H_{\text{aff}}^{i}(\text{Vect}(\mathbb{RP}^{1}), \mathfrak{aff}(1); \mathcal{D}_{\lambda,\mu\gamma}) \simeq \left\{ \begin{array}{c} 0.2cm \quad \mathbb{R} \quad \text{if } (\lambda,\mu) \in \left\{ (-2,0), (0,-2), (-1,0), (0,-1), \\ (-1,-1), (-\frac{2}{3}, -\frac{1}{3}), (-\frac{1}{3}, -\frac{2}{3}), \\ (-\frac{3}{2}, -\frac{1}{2}) \end{array} \right\}, \quad (16)$$

(6) If  $v - \mu - \lambda = 6$ , then

$$H_{\text{aff}}^{i}(\text{Vect}(\mathbb{RP}^{i}), \mathfrak{aff}(1); \mathcal{D}_{\lambda,\mu\nu}) \simeq \begin{cases} 0.2cm \quad \mathbb{R} \quad \text{if } (\lambda, \mu) \in \left\{ (-\frac{1}{2}, -\frac{1}{4}), (0, -\frac{3}{4}), (-\frac{3}{2}, -\frac{3}{2}), (-\frac{1}{2}, -\frac{1}{2}), (-\frac{1}{2}, -\frac{1}{2}), (-\frac{1}{2}, -\frac{1}{2}), (-\frac{1}{2}, -2), (-\frac{1}{2}, -\frac{1}{2}), (-\frac{1}{2}, -2), (-\frac{1}{2$$

(7) If  $v - \mu - \lambda = 7$ , then

$$H_{aff}^{i}(\operatorname{Vect}(\mathbb{RP}^{i}), \mathfrak{aff}(1); \mathcal{D}_{i,\mu;\nu}) \simeq \begin{cases} 0.2cm \quad \mathbb{R} \quad \text{if } (\lambda, \mu) \in \begin{cases} (0,0), (-\frac{1}{7}, 0), (0, -\frac{1}{7}), \\ (-\frac{3}{5}, 0), (0, -\frac{3}{5}), (0, -1), \\ (-2, -1), (-\frac{1}{2}, -\frac{5}{2}), (0, -3) \end{cases} \end{cases}, \quad (18)$$

(8) If  $\nu - \mu - \lambda = 8$ , then

$$H_{\text{deff}}^{1}(\text{Vect}(\mathbb{RP}^{1}), \mathfrak{aff}(1); \mathcal{D}_{\lambda, \mu; \nu}) \simeq \begin{cases} 0.2cm \quad \mathbb{R} \quad \text{if } (\lambda, \mu) \in \begin{cases} (0, 0), (0, -\frac{1}{8}), (-\frac{7}{7}, 0), \\ (0, -\frac{2}{7}), (-\frac{1}{2}, 0), (0, -\frac{1}{2}), \\ (-\frac{3}{5}, -\frac{1}{5}), (0, -2), (0, -\frac{7}{2}) \end{cases} \end{cases}, \quad (19)$$

(9) If  $v - \mu - \lambda = 9$ , then

$$H_{\text{aff}}^{1}(\text{Vect}(\mathbb{RP}^{1}),\mathfrak{aff}(1);\mathcal{D}_{\lambda,\mu;\nu}) \simeq \left\{ \begin{array}{c} 0.2cm \quad \mathbb{R} \quad \text{if } (\lambda,\mu) \in \left\{ \begin{array}{c} (0,0),(0,-\frac{1}{9}),(-\frac{1}{4},0), \\ (-\frac{2}{3},0),(0,-\frac{2}{3}),(-\frac{3}{5},-\frac{2}{5}), \\ (0,-\frac{7}{3}),(0,-4) \end{array} \right\}, \quad (20)$$

(10) If  $\nu - \mu - \lambda = 10$ , then

$$H_{\text{diff}}^{1}(\text{Vect}(\mathbb{RP}^{1}), \mathfrak{aff}(1); \mathcal{D}_{\lambda,\mu;\nu}) \simeq \begin{cases} 0.2cm \quad \mathbb{R} \quad \text{if } (\lambda,\mu) \in \begin{cases} (0,0), (0, -\frac{1}{10}), (-\frac{2}{9}, 0), \\ (-\frac{5}{6}, 0), (0, -\frac{5}{6}), (-\frac{3}{5}, -\frac{3}{5}), \\ (0, -\frac{8}{3}), (0, -\frac{9}{2}) \end{cases} \end{cases}, \quad (21)$$

(11) If  $v - \mu - \lambda = 11$ , then

$$H^{1}_{aar}(\operatorname{Vect}(\mathbb{RP}^{1}), \mathfrak{aff}(1); \mathcal{D}_{\lambda,\mu\nu}) \simeq \begin{cases} 0.2cm \quad \mathbb{R} \quad \text{if } (\lambda,\mu) \in \begin{cases} (0,0), (0, -\frac{1}{11}), (0, -\frac{1}{5}), \\ (0, -\frac{1}{3}), (0, -\frac{5}{5}), (-\frac{7}{5}, 0), \\ (0, -\frac{7}{5}), (-\frac{1}{2}, -\frac{9}{2}) \end{cases} \end{cases}, \quad (22)$$

(ii) If  $\nu - \mu - \lambda$  is semi-integer but  $\lambda$  and  $\mu$  are generic then,

 $\mathrm{H}^{1}_{\mathrm{diff}}(\mathrm{Vect}(\mathbb{RP}^{1}),\mathfrak{aff}(1);\mathcal{D}_{\lambda,\mu;\nu})=0.$ 

**Proof of Theorem 5.1:** To proof Theorem (5.1) we proceed by following the three steps:

• We will investigate the dimension of the space of operators that satisfy the 1-cocycle condition. By Proposition (4.4), its dimension is at most  $\frac{1}{2}k(k-1)$ , where  $k=\nu-\mu-\lambda+1$ , since any 1-cocycle that vanishes on aff(1) is certainly aff(1)-invariant.

• We will study all trivial 1-cocycles, namely, operators of the form  $L_x B$ ,

where *B* is a bilinear operator. As our 1-cocycles vanish on the Lie algebra  $\mathfrak{aff}(1)$ , it follows that the operator *B* coincides with the transvectant  $J_k^{\lambda,\mu}$ .

• By taking into account Part 1 and Part 2 and depending on  $\lambda$  and  $\mu$  the dimension of the cohomology group  $H^1_{diff}(Vect(\mathbb{RP}^1),\mathfrak{aff}(1);\mathcal{D}_{\lambda,\mu;\nu})$  will be equal to

dim(operators that are 1-cocycles) – dim(operators of the form  $L_X J_k^{\lambda,\mu}$ ).

Now, clearly the coboundary  $L_{\chi}J_{k}^{\lambda,\mu}$  has the following form:

$$L_{X}J_{k}^{\lambda,\mu}(X,\phi,\psi) = \sum_{i+j+l=k+1} \beta_{i,j,l} X^{(i)} \phi^{(j)} \psi^{(l)},$$
(23)

where

 $\beta_{0,j,l} = \beta_{1,j,l} = 0.$ 

The following Lemma is proved directly which will be useful in the proof of Theorem 5.1.

#### Lemma 5.2

For  $\lambda, \mu \in \mathbb{R}$ 

$$\beta_{\alpha,\beta,k-\alpha-\beta+1} = - \left( \binom{\alpha+\beta-1}{\alpha} + \lambda \binom{\alpha+\beta-1}{\alpha-1} \right) \gamma_{\alpha+\beta-1,k-\alpha-\beta+1} - \left( \binom{k-\beta}{\alpha} + \mu \binom{k-\beta}{\alpha-1} \right) \gamma_{\beta,k-\beta},$$

where  $\alpha \ge 2$  and  $\beta \ge 0$ .

We need also the following Lemma.

#### Lemma 5.3

Every 1-cocycle on Vect( $\mathbb{RP}^1$ ) with values in  $\mathcal{D}_{\lambda,\mu,\nu}$ ) is differentiable Proof [7].

Now we are in position to prove Theorem (5.1). By Lemma (5.3), any 1-cocycle on Vect( $\mathbb{RP}^1$ ) should retains the following general form:

$$C(X,\phi,\psi) = \sum_{i+j+l=k} c_{i,j,l} X^{(i)} \phi^{(j)} \psi^{(l)}, \qquad (24)$$

where  $c_{ij,l}$  are constants. The fact that this 1-cocycle vanishes on  $\mathfrak{aff}(1)$  implies that

 $c_{0,i,l} = c_{1,i,l} = 0.$ 

The 1-cocycle condition reads as follows: for all  $\varphi \in \mathcal{F}_{\lambda}$ , for all  $\psi \in \mathcal{F}_{\mu}$ and for all  $X \in \text{Vect}(\mathbb{RP}^1)$ , one has

 $c([X,Y],\phi,\psi) - L_X^{\lambda,\mu;\nu}B(Y,\phi,\psi) + L_Y^{\lambda,\mu;\nu}B(X,\phi,\psi) = 0.$ 

The case where  $\nu - \mu - \lambda = 1$ : In this case, according to Proposition 4.4, the 1-cocycle (24) can be expressed as follows:

$$\Upsilon(X,\phi,\psi) = c_{2,0,0} X'' \phi \psi.$$

By a direct computation, we can see that the 1-cocycle condition is always satisfied. Let us study the triviality of this 1-cocycle. A direct computation proves that

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 $L_{X}J_{1}^{\lambda,\mu} = \beta_{2,0,0}X''\phi\psi = -(\lambda\gamma_{1,0} + \mu\gamma_{0,1})X''\phi\psi.$ 

So, for ( $\lambda,\mu)$ =(0,0), the coeffcient  $c_{_{2,0,0}}$  cannot be eliminated by adding a coboundary. Hence, the cohomology space is one-dimensional. While for ( $\lambda,\mu)$ =(0,0), we can see that the coeffcient  $c_{_{2,0,0}}$  can be eliminated because  $\beta_{_{2,0,0}}$ =0. Hence, the cohomology is zero-dimensional.

**The case where**  $\nu - \mu - \lambda = 2$ **:** In this case, according to Proposition 4.4, the 1-cocycle (24) can be expressed as follows:

 $\Upsilon(X,\phi,\psi) = c_{3,0,0} X^{''} \phi \psi + c_{2,1,0} X^{''} \phi' \psi + c_{2,0,1} X^{''} \phi \psi'.$ 

By a direct computation, we can see that the 1-cocycle condition is always satisfied. Let us study the triviality of this 1-cocycle. A direct computation proves that

$$L_X J_2^{\lambda,\mu} = \beta_{3,0,0} X''' \phi \psi + \beta_{2,1,0} X'' \phi' \psi + \beta_{2,0,1} X'' \phi \psi'.$$
  
where  
$$\beta_{3,0,0} = -\lambda \gamma_{2,0} - \mu \gamma_{0,2}; \ \beta_{2,1,0} = -(2\lambda + 1)\gamma_{2,0} - \mu \gamma_{1,1} \ and \ \beta_{2,0,1} = -\lambda \gamma_{1,1} - (2\mu + 1)\gamma_{0,2}.$$

So, for  $(\lambda, \mu) = (0, 0), (-\frac{1}{2}, 0), (0, -\frac{1}{2})$ , the cohomology space

is one-dimensional, since only one of the coefficients  $c_{_{3,0,0}}, c_{_{2,1,0}}$  or  $c_{_{2,0,1}}$  cannot be eliminated by adding a coboundary. While for  $(\lambda,\mu) \neq (0,0), (-\frac{1}{2},0), (0,-\frac{1}{2})$ , the coeffcient  $c_{_{3,0,0}}, c_{_{2,1,0}}$  and  $c_{_{2,0,1}}$  can be eliminated because  $\beta_{_{3,0,0}}, \beta_{_{2,1,0}}$  and  $\beta_{_{2,0,1}}$  are nonzero. Hence, the cohomology space is zero-dimensional.

The case where  $\nu - \mu - \lambda \ge 3$ : In this case, the 1-cocycle condition is equivalent to the system:

$$\begin{pmatrix} \begin{pmatrix} \alpha+\beta-1\\ \alpha \end{pmatrix} - \begin{pmatrix} \alpha+\beta-1\\ \alpha-1 \end{pmatrix} \end{pmatrix} c_{\alpha+\beta-1,\gamma,a} + \begin{pmatrix} \begin{pmatrix} \alpha+\gamma-1\\ \alpha \end{pmatrix} + \lambda \begin{pmatrix} \alpha+\gamma-1\\ \alpha-1 \end{pmatrix} \end{pmatrix} c_{\beta,\alpha+\gamma-1,a} - \begin{pmatrix} \begin{pmatrix} \beta+\gamma-1\\ \beta \end{pmatrix} + \lambda \begin{pmatrix} \beta+\gamma-1\\ \beta-1 \end{pmatrix} \end{pmatrix} c_{\alpha,\beta+\gamma-1,a} + \begin{pmatrix} \begin{pmatrix} \alpha+a-1\\ \alpha \end{pmatrix} + \mu \begin{pmatrix} \alpha+a-1\\ \alpha-1 \end{pmatrix} \end{pmatrix} c_{\beta,\gamma,\alpha+a-1}$$
(25)  
 -  $\begin{pmatrix} \begin{pmatrix} \beta+a-1\\ \beta \end{pmatrix} + \lambda \begin{pmatrix} \beta+a-1\\ \beta-1 \end{pmatrix} \end{pmatrix} c_{\alpha,\gamma,\beta+a-1} = 0,$ 

where  $\alpha + \beta + \gamma + a = k+1$ ,  $\alpha > \beta \ge 2$ ,  $\alpha > \gamma$  and  $\alpha > a$ , obtained from the coefficient of  $X^{(\alpha)} Y^{(\beta)(\gamma)(a)}$ .

This system can be deduced by a simple computation. Of course, such a system has at least one solution in which the solutions  $c_{i,j,l}$  are just the coefficients  $\beta_{i,j,l}$  of the coboundaries (23).

**The case where**  $\nu - \mu - \lambda = 3$ : In this case, according to Proposition 4.4, the space of solutions is spanned by:

 $c_{4,0,0,}\ c_{3,1,0,}\ c_{3,0,1,}\ c_{2,2,0,}\ c_{2,1,1,}\ c_{2,0,2.}$ 

Moreover, by formula (25), we readily obtain:

$$-2c_{2,0,0} + \lambda c_{2,0,0} \lambda c_{2,0,0} + \mu c_{2,0,0} - \mu c_{2,0,0} = 0.$$

Thus, we have just proved that the coefficients of every 1-cocycle is expressed in terms of

 $C_{3,1,0,} C_{3,0,1,} C_{2,2,0,} C_{2,1,1,} C_{2,0,2.}$ 

A direct computation proves that

$$\begin{split} L_{x}J_{3}^{\lambda,\mu} &= \beta_{2,1,1}X^{*}\phi'\psi' + \beta_{2,2,0}(\frac{\lambda}{2}X^{4}\phi\psi + X^{*}\phi^{*}\psi) + \beta_{2,0,2}(\frac{\mu}{2}X^{(4)}\phi\psi + X^{*}\phi\psi^{*}) \\ &+ \beta_{3,1,0}(-\frac{\lambda}{2}X^{(4)}\phi\psi + X^{*}\phi'\psi) + \beta_{3,0,1}(-\frac{\mu}{2}X^{(4)}\phi\psi + X^{*}\phi\psi'). \end{split}$$

where

$$\begin{split} \beta_{3,1,0} &= -(3\lambda+1)\gamma_{3,0} - \mu\gamma_{1,2} &, \quad \beta_{3,0,1} &= -\lambda\gamma_{2,1} - (3\mu+1)\gamma_{0,3}, \\ \beta_{2,2,0} &= -3(\lambda+1)\gamma_{3,0} - \mu\gamma_{2,1} &, \quad \beta_{2,0,2} &= -\lambda\gamma_{1,2} - 3(\mu+1)\gamma_{0,3} \\ \beta_{2,1,1} &= -(2\lambda+1)\gamma_{2,1} - (2\mu+1)\gamma_{1,2}. \end{split}$$

So, for  $(\lambda, \mu) = (-\frac{1}{3}, 0), (0, -\frac{1}{3}), (-1, 0), (0, -1), (-\frac{1}{2}, -\frac{1}{2})$ , the cohomology space is one-dimensional, since only one of the coefficients  $c_{3,1,0}$ ,  $c_{3,0,1}$ ,  $c_{2,2,0}$ ,  $c_{2,1,1}$  or  $c_{2,0,2}$  cannot be eliminated by adding a coboundary. While for  $(\lambda, \mu) \neq (-\frac{1}{3}, 0), (0, -\frac{1}{3}), (-1, 0), (0, -1), (-\frac{1}{2}, -\frac{1}{2})$ , the coeffcient  $c_{3,1,0}$ ,  $c_{3,0,1}$ ,  $c_{2,2,0}$ ,  $c_{2,1,1}$  and  $c_{2,0,2}$  can be eliminated because  $\beta_{3,1,0}$ ,  $\beta_{3,0,1}$ ,  $\beta_{2,2,0}$ ,  $\beta_{2,1,1}$  and  $\beta_{2,0,2}$  are nonzero. Hence, the cohomology space is zero-dimensional.

**The case where**  $\nu - \mu - \lambda = 4$ **:** In this case, according to Proposition 4.4, the space of solutions is spanned by:

$$c_{5,0,0}, c_{4,1,0}, c_{4,0,1}, c_{3,2,0}, c_{3,1,1}, c_{3,0,2}, c_{2,3,0}, c_{2,2,1}, c_{2,1,2}, c_{2,0,3}.$$

Moreover, by formula (25), we readily obtain:

$$\begin{aligned} -2c_{4,1,0} + (3\lambda + 1)c_{2,3,0} - (2\lambda + 1)c_{3,2,0} + \mu c_{2,1,2} - \mu c_{3,1,1} &= 0, \\ -2c_{4,0,1} + \lambda c_{2,2,1} - \lambda c_{3,1,1} + (3\mu + 1)c_{2,0,3} - (2\mu + 1)c_{3,0,2} &= 0, \\ -5c_{5,0,0} + \lambda c_{2,3,0} - \lambda c_{4,1,0} + \mu c_{2,0,3} - \mu c_{4,0,1} &= 0. \end{aligned}$$

Thus, we have just proved that the coefficients of every 1-cocycle is expressed in terms of

$$C_{3,2,0}, C_{3,1,1}, C_{3,0,2}, C_{2,3,0}, C_{2,2,1}, C_{2,1,2}, C_{2,0,3}.$$

A direct computation confirms that, the coefficients of  $L_{\chi}J_4^{\lambda,\mu}$  are expressed in terms of:

$$\begin{aligned} \beta_{2,2,1} &= -3(\lambda+1)\gamma_{3,1} - (2\mu+1)\gamma_{2,2} \quad , \quad \beta_{2,1,2} &= -(2\lambda+1)\gamma_{2,2} - 3(\mu+1)\gamma_{1,3}, \\ \beta_{3,2,0} &= -2(3\lambda+2)\gamma_{4,0} - \mu\gamma_{2,2} \quad , \quad \beta_{3,0,2} &= -\lambda\gamma_{2,2} - 2(3\mu+2)\gamma_{0,4}, \\ \beta_{2,3,0} &= -2(2\lambda+3)\gamma_{4,0} - \mu\gamma_{3,1} \quad , \quad \beta_{2,0,3} &= -\lambda\gamma_{1,3} - 2(2\mu+3)\gamma_{0,4}, \\ \beta_{3,1,1} &= -(3\lambda+1)\gamma_{3,1} - (3\mu+1)\gamma_{1,3}. \end{aligned}$$
So, for  $(\lambda,\mu) = (-\frac{2}{2},0), (0,-\frac{2}{2}), (-\frac{3}{2},0), (0,-\frac{3}{2}), (-\frac{1}{2},-1), (-1,-\frac{1}{2}), (-\frac{1}{2},-\frac{1$ 

So, for  $(\lambda, \mu) = (-\frac{2}{3}, 0), (0, -\frac{2}{3}), (-\frac{3}{2}, 0), (0, -\frac{3}{2}), (-\frac{1}{2}, -1), (-1, -\frac{1}{2}), (-\frac{1}{3}, -\frac{1}{3}),$ the cohomology space is one-dimensional, since only one of the coefficients  $C_{3,2,0}, C_{3,1,1}, C_{3,0,2}, C_{2,3,0}, C_{2,2,1}, C_{2,1,2}$  Of  $C_{2,0,3}$  cannot be eliminated by adding a coboundary. While for  $(\lambda, \mu) \neq (-\frac{2}{3}, 0), (0, -\frac{2}{3}), (-\frac{3}{2}, 0), (0, -\frac{3}{2}), (-\frac{1}{2}, -1), (-1, -\frac{1}{2}), (-\frac{1}{3}, -\frac{1}{3}),$  the coefficient  $c_{3,2,0}, c_{3,1,1}, c_{3,0,2}, c_{2,3,0}, c_{2,2,1}, c_{2,1,2}$  and  $c_{2,0,3}$  can be eliminated because  $\beta_{3,2,0}, \beta_{3,1,1}, \beta_{3,0,2}, \beta_{2,3,0}, \beta_{2,2,1}, \beta_{2,1,2}$  and  $\beta_{2,0,3}$  are nonzero. Hence, the cohomology space is zero-dimensional.

The case where  $\nu - \mu - \lambda = 5$ : In this case, according to Proposition 4.4, the space of solutions is spanned by:

$$\begin{split} & \mathcal{C}_{6,0,0}, \mathcal{C}_{5,1,0}, \mathcal{C}_{5,0,1}, \mathcal{C}_{4,2,0}, \mathcal{C}_{4,1,1}, \mathcal{C}_{4,0,2}, \mathcal{C}_{3,3,0}, \mathcal{C}_{3,2,1}, \\ & \mathcal{C}_{3,1,2}, \mathcal{C}_{3,0,3}, \mathcal{C}_{2,4,0}, \mathcal{C}_{2,3,1}, \mathcal{C}_{2,2,2}, \mathcal{C}_{2,1,3}, \mathcal{C}_{2,0,4}. \end{split}$$

Moreover, by formula (25), we readily obtain:

 $\begin{aligned} &-2c_{4,1,1}+(3\lambda+1)c_{2,3,1}-(2\lambda+1)c_{3,2,1}+(3\mu+1)c_{2,1,3}-(2\mu+1)c_{3,1,2}=0,\\ &-5c_{5,1,0}+(4\lambda+1)c_{2,4,0}-(2\lambda+1)c_{4,2,0}+\mu c_{2,1,3}-\mu c_{4,1,1}=0,\\ &-5c_{5,0,1}+\lambda c_{2,3,1}-\lambda c_{4,1,1}+(4\mu+1)c_{2,0,4}-(2\mu+1)c_{4,0,2}=0,\\ &-2c_{4,2,0}+2(3\lambda+2)c_{2,4,0}-3(\lambda+1)c_{3,3,0}+\mu c_{2,2,2}-\mu c_{3,2,1}=0,\\ &-2c_{4,0,2}+\lambda c_{2,2,2}-\lambda c_{3,1,2}+2(3\mu+2)c_{2,0,4}-3(\mu+1)c_{3,0,3}=0,\\ &-9c_{6,0,0}+\lambda c_{2,4,0}-\lambda c_{5,1,0}+\mu c_{2,0,4}-\mu c_{5,0,1}=0,\\ &-5c_{6,0,0}+\lambda c_{3,3,0}-\lambda c_{4,2,0}+\mu c_{3,0,3}-\mu c_{4,0,2}=0. \end{aligned}$ 

Thus, we have just proved that the coefficients of every 1-cocycle is expressed in terms of

$$C_{2,4,0}, C_{2,0,4}, C_{3,3,0}, C_{3,0,3}, C_{3,2,1}, C_{3,1,2}, C_{2,3,1}, C_{2,2,2}.$$

A direct computation confirms that, the coefficients of  $L_{\chi}J_5^{\lambda,\mu}$  are expressed in terms of:

$$\begin{array}{ll} \beta_{3,2,1} = -2(3\lambda+2)\gamma_{4,1} - (3\mu+1)\gamma_{2,3} &, & \beta_{3,1,2} = -(3\lambda+1)\gamma_{3,2} - 2(3\mu+2)\gamma_{1,4}, \\ \beta_{2,3,1} = -2(2\lambda+3)\gamma_{4,1} - (2\mu+1)\gamma_{3,2} &, & \beta_{2,2,2} = -3(\lambda+1)\gamma_{3,2} - 3(\mu+1)\gamma_{2,3}, \\ \beta_{3,3,0} = -10(\lambda+1)\gamma_{5,0} - \mu\gamma_{3,2} &, & \beta_{3,0,3} = -\lambda\gamma_{2,3} - 10(\mu+1)\gamma_{0,5}, \\ \beta_{2,4,0} = -5(\lambda+2)\gamma_{5,0} - \mu\gamma_{4,1} &, & \beta_{2,0,4} = -\lambda\gamma_{1,4} - 5(\mu+2)\gamma_{0,5}. \end{array}$$

So, for  $(\lambda,\mu) = (-\frac{2}{3}, -\frac{1}{3}), (-\frac{1}{3}, -\frac{2}{3}), (-\frac{3}{2}, -\frac{1}{2}), (-1, -1), (-1, 0), (0, -1), (-2, 0), (0, -2))$ , the cohomology space is one-dimensional, since only one of the coefficients  $C_{2,4,0}, C_{2,0,4}, C_{3,3,0}, C_{3,0,3}, C_{3,2,1}, C_{3,1,2}, C_{2,3,1}$  Of  $C_{2,2,2}$  cannot be eliminated by adding a coboundary. While for  $(\lambda, \mu) \neq (-\frac{2}{3}, -\frac{1}{3}), (-\frac{1}{3}, -\frac{2}{3}), (-\frac{3}{2}, -\frac{1}{2}), (-1, -1), (-1, 0), (0, -1), (-2, 0), (0, -2),$  the coeffcient  $c_{2,4,0}, c_{2,0,4}, c_{3,3,0}, c_{3,0,3}, c_{3,2,1}, c_{3,1,2}, c_{2,3,1}$  and  $c_{2,2,2}$  can be eliminated because  $\beta_{2,4,0}, \beta_{2,0,4}, \beta_{3,3,0}, \beta_{3,0,3}, \beta_{3,2,1}, \beta_{3,1,2}, \beta_{2,3,1}$  and  $\beta_{2,2,2}$  are nonzero. Hence, the cohomology space is zero-dimensional.

The case where  $\nu - \mu - \lambda = 6$ : In this case, according to Proposition 4.4 together with formulas (25), we check that the coefficients of every 1-cocycle are expressed in terms of

$$c_{4,2,1}, c_{4,0,3}, c_{3,2,2}, c_{3,1,3}, c_{3,0,4}, c_{2,5,0}, c_{2,4,1}, c_{2,3,2}, c_{2,1,4}.$$

A direct computation confirms that, the coefficients of  $L_X J_6^{\lambda,\mu}$  are expressed in terms of:

 $\beta_{4,2,1}, \beta_{4,0,3}, \beta_{3,2,2}, \beta_{3,1,3}, \beta_{3,0,4}, \beta_{2,5,0}, \beta_{2,4,1}, \beta_{2,3,2}, \beta_{2,1,4}.$ 

So, in the same way as before, by Lemma 5.2, we can see, with the help of the maple, that the cohomology space is given as in (17).

**The case where**  $\nu - \mu - \lambda = 7$ **:** In this case, according to Proposition 4.4 together with formulas (25), we check that the coefficients of every 1-cocycle are expressed in terms of

 $C_{8,0,0}, C_{7,1,0}, C_{7,0,1}, C_{5,3,0}, C_{5,0,3}, C_{4,0,4}, C_{2,4,2}, C_{2,1,5}, C_{2,0,6}.$ 

A direct computation confirms that, the coefficients of  $L_{\chi}J_{\gamma}^{\lambda,\mu}$  are expressed in terms of:

$$\beta_{8,0,0},\beta_{7,1,0},\beta_{7,0,1},\beta_{5,3,0},\beta_{5,0,3},\beta_{4,0,4},\beta_{2,4,2},\beta_{2,1,5},\beta_{2,0,6}.$$

So, in the same way as before, by Lemma 5.2, we can see, with the help of the maple, that the cohomology space is given as in (18).

**The case where**  $v - \mu - \lambda = 8$ : In this case, according to Proposition 4.4 together with formulas (25), we check that the coefficients of every 1-cocycle are expressed in terms of:

 $C_{9,0,0}, C_{8,0,1}, C_{7,2,0}, C_{7,0,2}, C_{6,3,0}, C_{6,0,3}, C_{5,3,1}, C_{3,0,6}, C_{2,0,7}.$ 

A direct computation confirms that, the coefficients of  $L_{\chi}J_8^{\lambda,\mu}$  are expressed in terms of:

$$\beta_{9,0,0},\beta_{8,0,1},\beta_{7,2,0},\beta_{7,0,2},\beta_{6,3,0},\beta_{6,0,3},\beta_{5,3,1},\beta_{3,0,6},\beta_{2,0,7}.$$

So, in the same way as before, by Lemma 5.2, we can see, with the help of the maple, that the cohomology space is given as in (19).

The case where  $v - \mu - \lambda = 9$ : In this case, according to Proposition 4.4 together with formulas (25), we check that the coefficients of every 1-cocycle are expressed in terms of

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 $C_{10,0,0}, C_{9,0,1}, C_{8,2,0}, C_{6,4,0}, C_{6,0,4}, C_{5,3,2}, C_{3,0,7}, C_{2,0,8}.$ 

A direct computation confirms that, the coefficients of  $L_X J_9^{\lambda,\mu}$  are expressed in terms of:

$$\beta_{10,0,0}, \beta_{9,0,1}, \beta_{8,2,0}, \beta_{6,4,0}, \beta_{6,0,4}, \beta_{5,3,2}, \beta_{3,0,7}, \beta_{2,0,8}.$$

So, in the same way as before, by Lemma 5.2, we can see, with the help of the maple, that the cohomology space is given as in (20).

**The case where**  $\nu - \mu - \lambda = 9$ : In this case, according to Proposition 4.4 together with formulas (25), we check that the coefficients of every 1-cocycle are expressed in terms of

 $C_{11,0,0}, C_{10,0,1}, C_{9,2,0}, C_{6,5,0}, C_{6,0,5}, C_{5,3,3}, C_{3,0,8}, C_{2,0,9}.$ 

A direct computation confirms that, the coefficients of  $L_{\chi}J_{10}^{\lambda,\mu}$  are expressed in terms of:

 $\beta_{11,0,0}, \beta_{10,0,1}, \beta_{9,2,0}, \beta_{6,5,0}, \beta_{6,0,5}, \beta_{5,3,3}, \beta_{3,0,8}, \beta_{2,0,9}.$ 

So, in the same way as before, by Lemma 5.2, we can see, with the help of the maple, that the cohomology space is given as in (21).

The case where  $v - \mu - \lambda = 11$ : In this case, according to Proposition 4.4 together with formulas (25), we check that the coefficients of every 1-cocycle are expressed in terms of

$$C_{12,0,0}, C_{11,0,1}, C_{10,0,2}, C_{9,0,3}, C_{7,0,5}, C_{5,7,0}, C_{5,0,7}, C_{2,1,9}.$$

A direct computation confirms that, the coefficients of  $L_{\chi} J_{11}^{\lambda,\mu}$  are expressed in terms of:

 $\beta_{12,0,0}, \beta_{11,0,1}, \beta_{10,0,2}, \beta_{9,0,3}, \beta_{7,0,5}, \beta_{5,7,0}, \beta_{5,0,7}, \beta_{2,1,9}.$ 

So, in the same way as before, by Lemma 5.2, we can see, with the help of the maple, that the cohomology space is given as in (22). This completes the proof.

#### Conjecture 5.1

For  $\nu - \mu - \lambda \in \mathbb{N} + 12$ ,  $\lambda$  and  $\mu$  are generic, one hase

 $H^{1}_{diff}(\operatorname{Vect}(\mathbb{RP}^{1}),\mathfrak{aff}(1);\mathcal{D}_{\lambda,\mu;\nu})=0.$ 

#### Conclusion

In this paper, we classify  $\mathfrak{aff}(1)$  -invariant linear differential operators from  $\operatorname{Vect}(\mathbb{RP}^1)$  to  $\mathcal{D}_{\mu,\nu}$  vanishing on  $\mathfrak{aff}(1)$ , where  $\mathcal{D}_{\mu,\nu}$ :=Homdiff( $\mathcal{F}_{\lambda}\otimes \mathcal{F};\mathcal{F}_{\nu}$ ) is the space of bilinear differential operators acting on weighted densities. This result allows us to compute the first differential  $\mathfrak{aff}(1)$ -relative cohomology of  $\operatorname{Vect}(\mathbb{RP}^1)$  with coefficients in  $\mathcal{D}_{\lambda,\mu,\nu}$ .

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