

# Speakable and Unspeakable in Special Relativity: the Ageing of the Twins in the Paradox

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## Abstract

In previous papers, we have presented a general formulation of special relativity, based on a weaker statement of the postulates. In this work, the paradigmatic example of the twin paradox is discussed in detail. Within the present formulation of special relativity, a “non-paradoxical” interpretation of the asymmetric ageing of the twins emerges. It is based exclusively on the *rhythms* of the clocks, which are not related by the standard textbook expressions and shall not be confused with clock *time readings*. Moreover, the current approach exposes the irrelevance of the acceleration of the returning twin in the discussion of the paradox.

**Keywords:** Special relativity; Lorentz transformation; Einstein speed; Twin paradox; IST transformation; Clock rhythms

## 1. Introduction

In previous works we undertook a reflection on the foundations of special relativity [1-6] where we present a general formulation of special relativity and reconcile the ideas of Lorentz and Poincaré of the existence of a “preferred reference frame” and Einstein’s “equivalence of all inertial frames.” A comprehensive overview of the theory is given in our last paper [6], hereafter denoted as I, where it is shown that it is not correct to speak about two philosophies, as they are different aspects of one and the same theory. Herein we extend our previous results and illustrate the power of the present formulation with a detailed discussion of the twin paradox.

The analysis of the “paradoxes” of a physical theory is an important matter, since it questions and assigns physical meaning to the statements of the theory. Following Feynman et al. [7], “a paradox is a situation which gives one answer when analyzed one way, and a different answer when analyzed another way, so that we are left in somewhat of a quandary as to actually what should happen. Of course, in physics there are never any real paradoxes because there is only one correct answer; at least we believe that nature will act in only one way (and that is the *right way*, naturally). So in physics a paradox is only a confusion in our own understanding”.

The famous twin paradox emerges in the standard interpretation of special relativity, which says that each twin ages slower than the other during any part of the to-and-fro traject, but when they meet one of them is older. This is of course not possible if “ageing slower” refers exclusively to rhythms of clocks along the trip and nothing else happens. Although the twin paradox is addressed and “solved” in any introductory course on special relativity, it was at the origin of more than 25,000 articles in the literature [8] since it was launched by Langevin in 1911 [9], and new publications arise regularly (the references given are therefore merely indicative [8,10-31]. Therefore, one can only suspect that “perhaps the last word on the twin paradox has yet to be said” [18].

In this paper, we argue that the key point to understand the twin paradox remains misinterpreted in its usual “solutions.” We start with a brief overview of the problem and then criticise the two main lines used in its standard “resolution:” the “simultaneity” and the “general relativity” arguments. In fact, despite giving the correct final answer, the usual explanations have a severe interpretation mistake, as they fail

both to make the distinction between clock rhythms and clock time readings [6] and to recognise the indeterminacy of special relativity [4,6].

The structure of this paper is the following. In the next section we very briefly review some of the main findings from our former work [1-6], with emphasis on the relation of rhythms between two clocks in relative motion and its connection with the usual time dilation expressions. The twin paradox is then presented and examined in the subsequent sections. Section 3 presents a short historical overview of the problem. Critical reviews of the standard “solutions” using either only special relativity and no reference to acceleration or making use of general relativity and acceleration are given in sections 4 and 5, respectively. Section 6 contains the general analysis made in the context of the present formulation of special relativity. Finally, section 7 summarises our main findings.

## 2. Brief Overview of the Theory

In this section we very quickly review some results from I [6]. We expected the reader to be familiar with that paper and call a special attention to its sections 5, 6 and 8, of particular relevance to the ensuing discussion.

### 2.1 The IST transformation, Lorentzian time and Einstein speed

Our formulation is based on the definition of the rest system – as the system where the one-way speed of light in vacuum is isotropic – and on the IST transformation (Inertial [32-34]–Synchronized [2,4,35-37] Tangherlini [38]). In the usual configuration where the axis of the rest frame  $S$  and a moving frame  $S'$  are aligned, the origin of  $S'$  moves along the  $x$ -axis of  $S$  with speed  $v$  in the positive direction, and the reference event is the overlapping of the origins of both frames at time zero, the IST transformation is given by Guerra and Abreu [2]

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$$x' = \gamma(x - vt)$$

$$t' = \frac{t}{\gamma}, \quad (1)$$

with

$$\gamma = \frac{1}{\sqrt{1 - \frac{v^2}{c^2}}}. \quad (2)$$

It is important to note that the Lorentz transformation is readily obtained from the IST transformation by introducing the offset factor

$$t'_L = t' - \frac{v}{c^2}x' \quad (3)$$

and substituting  $t'$  in (1) [2,4]. Therefore, any phenomenon that can be described by the Lorentz transformation can also be described by the IST transformation. We denote by Lorentzian clocks the clocks “synchronized” (i.e., “adjusted” or “set”) according to the Lorentz transformation, and by Lorentzian times the time readings exhibited by the Lorentzian clocks (3). Similarly, we denote by synchronized clocks the clocks adjusted according to the IST transformation and by synchronized times or simply “times” the time readings of the synchronized clocks.

The Lorentz transformation and the IST transformation can be associated with specific forms of clock “synchronization” [2]. However, as stated in I, one simple – and yet critical – issue that has to be clarified before engaging any discussion on the interpretation of special relativity or of special relativity results is to state the difference between clock *rhythms* (or clock tick rates) and clock *time readings* (or time coordinates), as further discussed below. Contrary to time readings, clock rhythms do *not* depend on any particular form of “synchronization.” Surprisingly, the failure to make this basic distinction is at the origin of several misunderstandings surrounding the theory, including the discussion of the twin paradox.

The velocity addition formula can be obtained easily from (1) [2]. If an object is moving at speed  $w$  in  $S$ , then its speed in the inertial frame  $S'$ ,  $w'_v$ , is simply

$$x' = w'_v t', \quad (4)$$

which can be written in the form

$$w'_v = \gamma_v^2 (w - v) = \frac{w - v}{1 - \frac{v^2}{c^2}}. \quad (5)$$

It is also essential to distinguish *speed* (or “synchronized speed”) from *Einstein speed* [2,4,5]. Speed is defined as in the previous equations and can be calculated from  $w'_v = \Delta x' / \Delta t'$ . In turn, the Einstein speed is defined from

$$w'_E = \frac{\Delta x'}{\Delta t'_L}, \quad (6)$$

i.e., its value is calculated with the difference of the time readings of “Lorentzian” clocks. Substituting (3) into (4), it is straightforward to show [2] that the Einstein velocity,  $w'_E$ , measured in a frame moving with speed  $v$  (in  $S$ ) of an object which has speed  $w$ , is

$$w'_E = \frac{w'_v}{1 - \frac{vw'_v}{c^2}}. \quad (7)$$

Replacing  $w'_v$  and using (5) this last expression can be rewritten as

$$w'_E = \frac{w - v}{1 - \frac{vw}{c^2}}. \quad (8)$$

## 2.2 General expression for the rhythms of two clocks in relative motion

In prior work it was thoroughly debated that the indeterminacy of special relativity implies that just with Lorentzian clocks it is not possible to know in which inertial frame clocks are actually running slower [4]. This result was derived more formally in section 8 of I, where the general expression for the rhythms of two clocks in relative motion was deduced. Due to its importance for the discussion of the twin paradox, we reproduce a big part of it here.

Consider two inertial frames,  $S'$  and  $S''$ , moving in the  $x$  direction respectively with speeds  $v$  and  $w$  in the rest system,  $S$ . Clock 1 is at the origin of  $S'$  and clock 2 is at the origin of  $S'$  and  $S''$ . The speed of clock 2 in  $S'$  is given by (5), the Einstein speed of clock 2 in  $S'$  is given by (8). Let us further define the *proper time*,  $\tau$ , in the usual way, as the time elapsed for one particular observer. Since the proper time is measured by a single clock, it is indeed associated with the clock *rhythm* and does not depend on the initial adjustment of distant clocks. The proper time of a clock in  $S''$  relates to the time elapsed in the rest system  $S$  by the time equation in (1), which can be written in the form

$$d\tau'' = \frac{dt}{\gamma_w}, \quad (9)$$

where  $dt$  are the differential times marked by the different clocks in the rest system  $S$  that are co-punctual with clock 2 at each instant. The relation of the rhythm of a clock 1 in  $S'$  with the rhythms of clocks in  $S$  is given by a similar expression, namely

$$d\tau' = \frac{dt}{\gamma_v}. \quad (10)$$

Therefore, the relation of the proper times of clocks in  $S''$  and in  $S'$  is

$$d\tau'' = d\tau' \frac{\gamma_v}{\gamma_w}. \quad (11)$$

Equation (11) establishes the relation of clock *rhythms* in two inertial frames. However, one final step is still missing. From (10) and (1) it directly follows

$$d\tau' = dt'. \quad (12)$$

Here, the equality is completely general and is valid whether or not  $dx'=0$  as  $x'$  does not appear in the second equation (1). In turn, as a consequence of (3), in general

$$d\tau' = dt' \neq dt'_L, \quad (13)$$

although, if  $dx'=0$ , then  $d\tau' = dt' = dt'_L$ . Similarly, we can write  $d\tau'' = dt'' \neq dt''_L$ , except when  $dx''=0$ . Substituting (12) in (11), we get

$$d\tau'' = dt' \frac{\gamma_v}{\gamma_w}. \quad (14)$$

The inverse relation is simply  $d\tau' = dt'' (\gamma_w / \gamma_v)$ .

As it is well-known from standard special relativity – and it is extremely easy to deduce from the Lorentz transformation – the proper time of clock 2 relates with the differential Lorentzian times through the Einstein speed,

$$d\tau'' = \frac{dt'_L}{\gamma_E}, \quad (15)$$

where  $\gamma_E$  is the Lorentz factor associated with the *relative* Einstein speed

between both moving frames  $v_E$  given by equation (8),  $\gamma_E = 1 / \sqrt{1 - \frac{v_E^2}{c^2}}$ . Contrary to (11), this expression corresponds to comparisons of *time readings* of Lorentzian clocks in  $S'$  and does *not* correspond to a relation with the clock rhythms in  $S'$ .

Finally, by noting that  $dx' = w'_E dt'_L$  and  $dx' = w'_v dt'$ , and using (13), equation (15) can be rewritten as

$$d\tau'' = d\tau' \frac{w'_v}{w'_E} \frac{1}{\gamma_E}. \quad (16)$$

This result confirms that Lorentzian clocks and Einstein speeds are not enough to determine in which frame clocks are running faster, since the relation between rhythms additionally involves the speed  $w'_v$ .

The indeterminacy of special relativity [4] has thus been expressed in an alternative way. If the rest system is inaccessible, then we *do not know* the value of the one-way speed of light in one inertial frame, nor can we know the value of  $w'_v$  in (16), and so *we cannot know in which of two inertial frames clocks are ticking slower*. This result, evidently, has profound implications in the analysis of the twin paradox.

### 3. Brief Overview

The twin paradox is one of the most famous paradoxes in physics. It emerges in the context of the standard interpretation of special relativity. Very briefly, consider twins Andrew and Bob [39]. Andrew stays on the Earth, while Bob flies to a distant star and back. Within the standard interpretation of special relativity, from Andrew's point of view Bob's clock is running slower both on the outward and return parts of the trip. Therefore, Bob is younger than Andrew when they meet up again. The "paradox" part of the story comes from the following alternate reasoning. Within the standard interpretation of special relativity, in Bob's frames it is Andrew's clock which is running slower, and so it is Andrew who is younger than Bob when they meet up again. Of course this is self-contradictory, as Andrew cannot be younger and older than Bob.

About 10 years ago Peter Pesic wrote a very interesting review article on the twin paradox [23]. In particular, he points out that, facing the paradox, Einstein addresses two arguments. The first one is the "simultaneity argument." It uses only special relativity and does not attribute importance to acceleration other than to break the symmetry of the situation of both twins [40]. This is indeed correct, as recognized by many physicists. For instance, Morin [39] states our time-dilation result holds only from the point of view of an inertial observer. The symmetry in the problem is broken by the acceleration. If both A and B are blindfolded, they can still tell who is doing the traveling, because B will feel the acceleration at the turnaround (...). For the entire outward and return parts of the trip, B does observe A's clock running slow, but enough strangeness occurs during the turning-around period to make A end up older. Note, however, that a discussion of acceleration is not required to quantitatively understand the "paradox".

The somewhat uncomfortable but inescapable remark is that the change of inertial reference frames by Bob can be made in a very short time, in the limit even instantaneously. And it is this very short period, as compared to the trip duration, which is responsible for making Bob younger. Moreover, if Bob makes a longer trip and if he does the turn

back exactly in the same way as before, he will return even younger than before. Thus, the effect of the very short turn-around period has to be higher now, in order to compensate for the longer period in which Bob sees Andrew ageing slower. The "solution" proposed by the standard interpretation of special relativity in this case is related to the "loss of simultaneity" exhibited by the time readings of Lorentzian clocks in different inertial frames [23,40]. If this is correct we have to acknowledge that "at the root of the twin paradox is the problem of synchronization of clocks" [19]. However, this raises another major alarm. In fact, as repeated several times along in this work, physical results cannot depend on the way in which one decides to set his own clocks. This "solution" is analyzed in section 4.

In 1918 Einstein has advanced the second argument, the "general relativity argument" [41]. It is now affirmed that we must consider general relativity in order to solve the paradox from the point of view of Bob. Since before and after the acceleration the argumentation of special relativity subsists, this comes back to attributing an important physical role to the acceleration. During the acceleration corresponding to the turn around period Bob would see Andrew age very quickly, because, according to the equivalence principle, Bob would see Andrew at a high gravitational potential. This second "solution" is discussed in section 5.

### 4. "Solution" without Acceleration

From the presentation of the broader view of special relativity proposed in I, in particular with the distinction between clock rhythms and clock time readings, together with the indeterminacy of special relativity, which does not allow to know in which inertial frame clocks are actually running slower, the meaningful solution to the paradox is already clear. The symmetry in the *description* of the outward trip between Andrew and Bob when we refer to Lorentzian times, repeated also for the return trip, does *not* correspond to a symmetry in the ageing (proper times) of the twins. Regarding proper times, we know that during the outward trip either Andrew or Bob is ageing slower and, without a reference to the rest system, we do not know which of them. And it may even happen that both are ageing at the same rhythm. The same occurs during the return trip. It is possible that it is always one of the twins who is ageing slower, both on the onward and on the return trip, or that one of the twins ages slower during the onward trip and the other during the return trip, or even that in one of the trips they are ageing at the same rate. However, we do know that when they meet it is Bob who is younger, and by which factor. Of course that all calculations can be made both with Lorentzian times and with synchronized times and from the point of view of each of the twins: all these calculations must give the same final result.

One quite convincing standard solution of the paradox using the "simultaneity argument" is obtained from the analysis of Minkowski diagrams [10,20,39]. Let us assume that Bob travels along the  $x'$  axis with *Einstein speed* of absolute value  $v_E$  in relation to Andrew, both in the outward and on the return trips. There are various Lorentzian clocks along the  $x'$  axis with photo cameras. At regular intervals  $\Delta t'_L = \Delta t'_A$  in Andrew's frame the cameras shoot. For each set of photos one and only one of the cameras has the photo of Bob. The situation is depicted in Figure 1a). The universe line of Andrew is the  $ct'$  axis ( $x' = 0$ ), whereas the universe line of Bob is represented in red. Obviously, Andrew and Bob agree on the number  $N$  of photos taken during the all trip. It is straightforward to obtain the conversion of lengths in the diagram to the Lorentzian times in each frame [(one unit in the  $ct''$

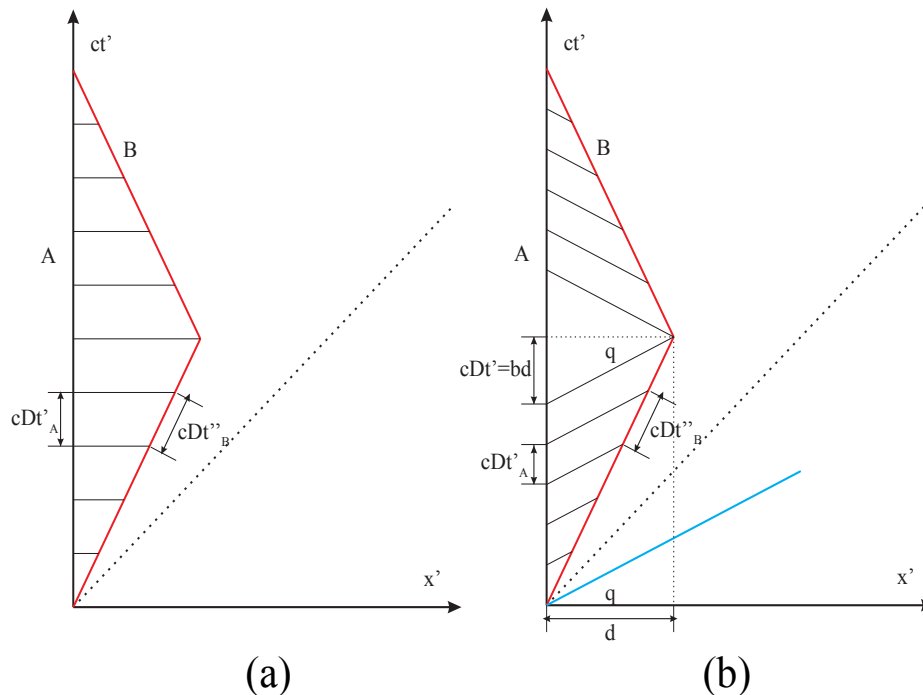


Figure 1: Minkowski diagram for the twin paradox from Andrew's point of view (a) and from Bob's point of view (b).

axis)/(one unit in the  $ct'$  axis) =  $\sqrt{(1 + \beta^2) / (1 - \beta^2)}$ , with  $\beta = v/c$  [39] and from simple trigonometry to get  $\Delta t_B'' = \Delta t_A' \sqrt{1 - \beta^2} = \Delta t_A' / \gamma_E$ . Notice that  $\Delta t_A'$  refers to a comparison of the time readings of Lorentzian clocks in Andrew's frame while  $\Delta t_B''$  corresponds to Bob's proper time, so that this result is simply the symmetric description of time dilation obtained with Lorentzian clocks (cf. Figure 2 in I and its respective discussion). The total time elapsed for Andrew is  $t_A' = N \Delta t_A'$ , while for Bob the time lapse is  $t_B'' = N \Delta t_B''$ , so that

$$t_B'' = t_A' / \gamma_E. \quad (17)$$

Therefore, Bob is younger than Andrew when the twins meet

To complete the analysis we have now to use the reverse reasoning, taking photos of Andrew from cameras in Bob's frame, shot when the clocks associated with the cameras display the same Lorentzian time. Only one of each set of photos has Andrew's photo. This new situation is represented in Figure 1b. There is not a single inertial frame associated with Bob, as the inertial frames for the onward and return trips are distinct. While Bob sees himself photographed at equal time intervals,  $\Delta t_B''$  in Figure 1a), Andrew seems himself being photographed regularly in the beginning of the trip,  $\Delta t_A'$  in Figure 1b), then there is a time interval  $2\Delta t_L'$  (corresponding to the proximity of Bob with the star, the turn-around point) where no photo is taken, and then Andrew is again photographed regularly. As before, Andrew and Bob agree with the number of photos taken. Let  $d$  be the distance between the Earth and the star (in the Earth frame). The total time of the trip for Andrew is simply

$$t_A' = 2d / v_E. \quad (18)$$

The total time elapsed for Bob is

$$t_B'' = N \Delta t_B''. \quad (19)$$

Furthermore, in this case we have the reverse situation as compared to the previous figure concerning the intervals  $\Delta t_A'$  and  $\Delta t_B''$ , with  $\Delta t_A' = \Delta t_B'' / \gamma_E$ , as derived from the Lorentz transformation or directly from the Minkowski diagram. Of course  $\Delta t_B''$  refers to comparisons of Lorentzian times, while  $\Delta t_A'$  corresponds to Andrew proper times, so that  $\Delta t_A'$  and  $\Delta t_B''$  are not the same as in the Figure 1a). It is also not difficult to show that  $2c\Delta t_L' = 2dv_E / c$ . In other words, the time that passes for Andrew while he is not photographed, corresponding to Bob's turn-around, is

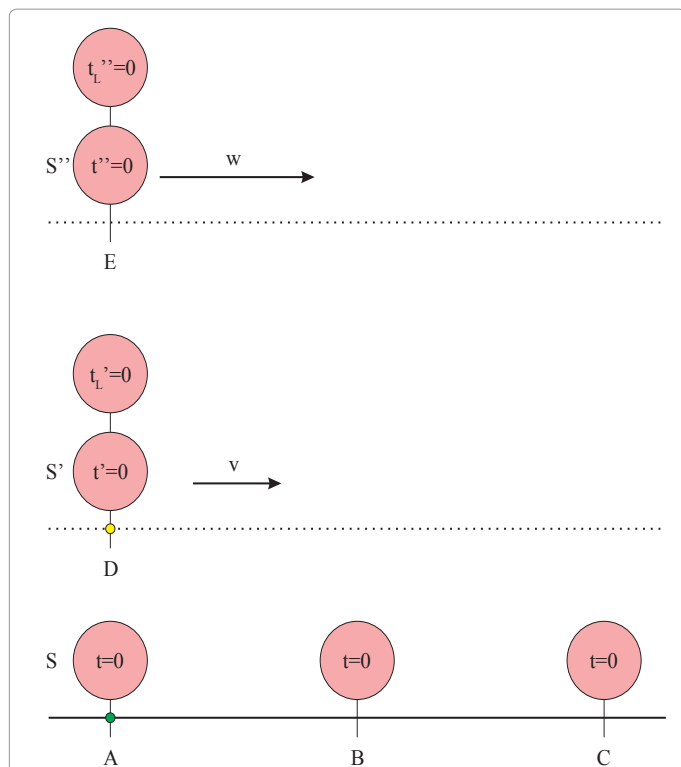
$$2\Delta t_L' = 2dv_E / c^2, \quad (20)$$

i.e., the offset factor (3). The final calculation for  $t_A'$  then gives

$$t_A' = N \Delta t_A' + (2\Delta t_L') = N \Delta t_B'' / \gamma + (2\Delta t_L') \quad (21)$$

from where, substituting (18–20), one recovers (17),  $t_A' = \gamma_E t_B''$ , as it had to be. There is no paradox, and a careful calculation made only within standard special relativity leads to the right result, both from Andrew's and from Bob's points of view. In turn, failing to consider the factor  $\Delta t_L'$  would lead to the paradoxical result  $t_A' = t_B'' / \gamma_E$ , contradictory with (17).

Despite the correctness of the result obtained with the Minkowski diagrams, the language used in the standard solution confuses "time" with Lorentzian times, making the interpretation somewhat awkward. In fact, without drawing the Minkowski diagrams and just using the Lorentz transformation, very few scientists would remember to add the offset factor  $2dv_E/c^2$  on a change of inertial frame corresponding to "Bob's point of view" to get the final calculation right. This calculation error comes from a correct intuition: if the Lorentzian time intervals would correspond to proper times, then one would not need to add, a bit artificially, some extra interval to the travel time of Andrew, which in addition depends on the distance, to correct for the instantaneous action of changing inertial frame. The need to add  $2dv_E/c^2$  to relate

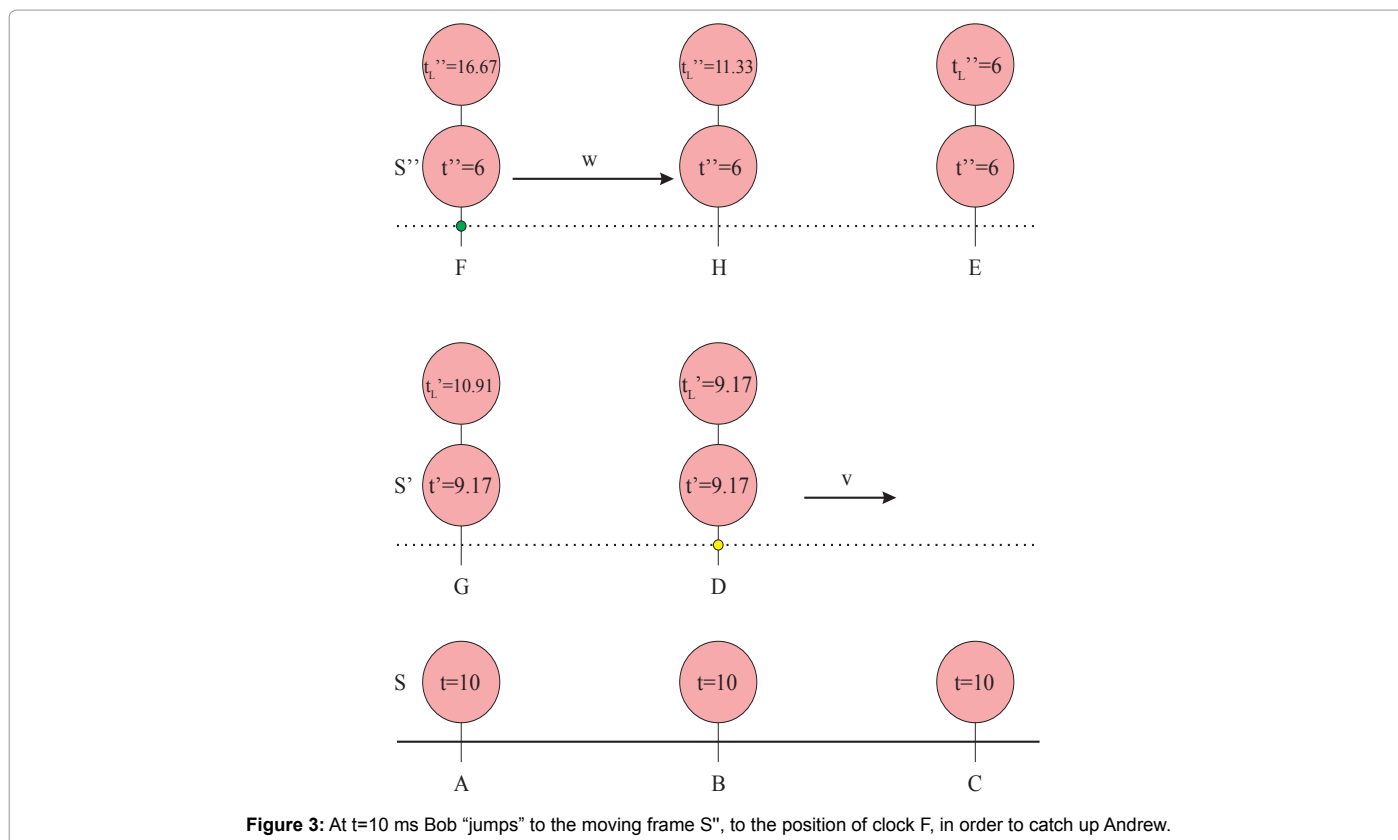


**Figure 2:** All times in the figures are expressed in milliseconds and rounded to two decimal places. Twins Andrew and Bob, initially in the moving frame  $S'$  separate at  $t=0$ ; Andrew (yellow dot) remains in frame  $S'$  with clock D, while Bob (green dot) "jumps" to the rest system  $S$  to the location of clock A.

the Lorentzian times and compensate the different offset factors (3) of Bob's two frames, together with the repeated remark that clock rhythms do not depend on the initial adjustment of the clocks, points the way to a clear and simple analysis of the twin paradox. Let us take one example to illustrate how does it work.

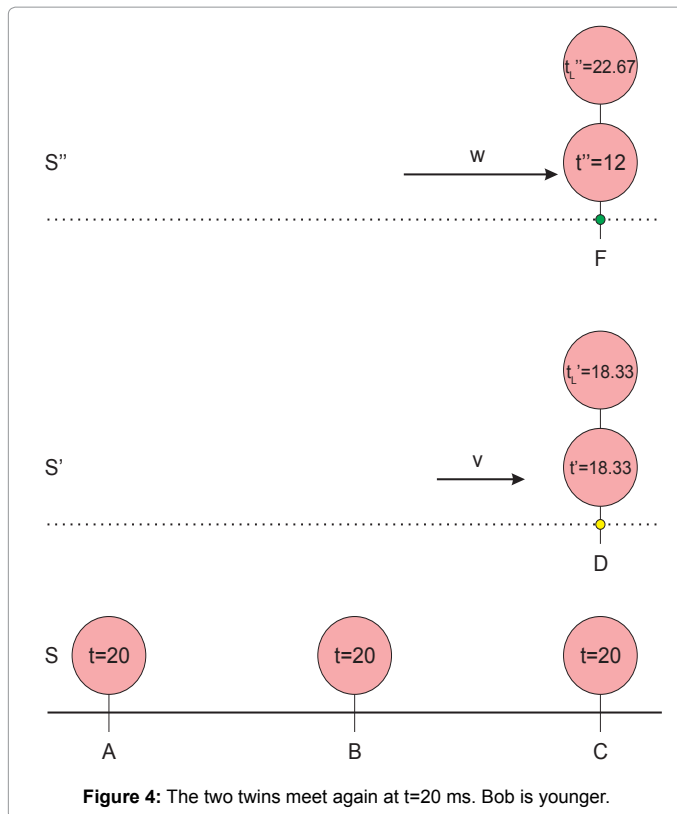
Consider the setup shown in Figure 2. Andrew and Bob are initially in a moving frame  $S'$ , which goes with speed  $v=0.4c$  in relation to the rest system  $S$ . A second moving frame  $S''$  goes with speed  $w=2v=0.8c$ . All frames are equipped both with synchronized and with Lorentzian clocks. The *same* situation can then be *described* with both types of clocks. As the physical reality is not affected by the description chosen, one can use both types of clocks and both descriptions at the same time no conflict at all [6]. An *apparent* conflict only arises if one confuses Lorentzian clocks with synchronized ones.

We start the analysis by looking at the time readings of the synchronized clocks, so that the time lapses correspond to proper times (section 2.2). To go to the distant star Bob "jumps" at  $t=0$  from  $S'$  into another frame, which, for simplicity, it happens to be the rest system  $S$ , staying there at rest for 10.00 milliseconds. From the "point of view of Andrew," Bob can be seen as "moving" outward, but still it is Bob who is ageing faster. Indeed, for Bob 10.00 ms have passed, while for Andrew the elapsed time was only about 9.17 ms, as shown by the synchronized clocks A and G in Figure 3, so that Bob is older by  $(10.00-9.17)$  ms = 0.83 ms. On frame  $S''$  only 6.00 ms have passed. Clock F from  $S''$  is now on top of Bob, who "jumps" into  $S''$  in order to rejoin Andrew. Both twins meet when 20.00 ms passed in the rest system, at the position of clock C as depicted in Figure 4. For Andrew the all process took approximately 18.33 ms. But for Bob only 16.00 ms have passed, 10.00 ms during the "outward" trip and  $(12.00-6.00)$  ms = 6.00 ms during the "return" trip. In this particular case, although



**Figure 3:** At  $t=10$  ms Bob "jumps" to the moving frame  $S''$ , to the position of clock F, in order to catch up Andrew.





Bob is ageing faster during the outward trip, he is ageing slower during the return trip. The latter journey compensates for the first part of the trip and it is actually Bob who is younger when the twins meet, by 2.33 ms. When speaking about proper times there is no difference in the descriptions made “from Andrew’s point of view” and “from Bob’s point of view,” they both agree with all the numbers presented above and describe the sequence of events exactly in the same way. Different variations of this example can be made, all leading to the same final result.

It is enlightening to crosscheck how the same situation is described from the comparison of the *time readings* of the Lorentzian clocks, within the reasoning of the standard interpretation of special relativity. Let us first see how Andrew argues. He considers himself at rest and will compare the time readings of different Lorentzian clocks in his frame, co-punctual with Bob, with the time readings of Bob’s clock. When Bob arrives at the turn-around point his clock (A) in S displays 10 ms (Figure 3). In turn, the Lorentzian clock in Andrew’s frame S’ (clock G) exhibits 10.91 ms. Thus, the difference in Lorentzian times is  $(10.91-10.00)$  ms=0.91 ms. In this way, one may be induced to think that Bob is younger by 0.91 ms (while in fact he is older by 0.83 ms, as shown before). During the return trip the comparison of Lorentzian times reveals that at the beginning of the return trip (Figure 3) Bob’s clock F is in advance of the S’ clock G by  $(16.67-10.91)$  ms=5.76 ms, while at the end of the trip (Figure 4) his clock F is in advance of the S’ clock D only by  $(22.67-18.33)$  ms=4.34 ms. Therefore, during the return trip one may be induced to think that Bob got younger by  $(5.76-4.34)$  ms=1.42 ms.

Just adding the differences of the Lorentzian time intervals we get the correct number and conclude that Bob is younger by  $(0.91+1.42)$  ms=2.33 ms, the same number obtained before, as it should be: the physical reality does not depend on its description. Adding

the Lorentzian time intervals provides the correct result, albeit an erroneous interpretation if we forget they are Lorentzian times and confuse them with clock rhythms.

Still within the reasoning of the standard interpretation of special relativity, Bob’s argumentation is much more striking and problematic. Bob considers himself to be at rest and so he expects to see Andrew ageing slower. He will compare the time readings of the Lorentzian clocks in his inertial frames (one frame during the onward trip, another one during the return trip) co-punctual with Andrew, with the time readings of Andrew’s clock. For the outward trip, Andrew got younger by  $(10.00-9.17)$  ms = 0.83 ms (clocks B and D in figure 3).

To keep the rules of the standard interpretation of special relativity, for the return trip we have to compare first the time readings of the Lorentzian clocks D and H in Figure 3, and then those of clocks D and F in Figure 4. Bob may then think that Andrew got even younger during the return trip, as for Andrew the difference in Lorentzian times is  $(18.33-9.17)$  ms=9.16 ms, whereas for Bob it is  $(22.67-11.33)$  ms=11.34 ms, so that one may think that Andrew got younger by  $(11.34-9.16)$  ms=2.18 ms. Therefore, not counting with any “strangeness” of the turn-around point, Bob may think that in total Andrew got younger by  $(0.83+2.18)$  ms=3.01 ms. However, Bob knows that at the turn-around point the Lorentzian clock H displays 11.33 ms Figure 3. However, his new Lorentzian clock F displays 16.67 ms! This difference of  $(16.67-11.33)$  ms=5.34 ms, which of course corresponds to the factor (20) resulting from the offset (3) of Lorentzian clocks, has to propagate to all clocks in S’, as we are interested only in time-differences.

This is all there is about the strangeness of the turn-around point. However, if we would want to keep the speech that “for the entire outward and return parts of the trip, Bob does observe Andrew’s clock running slow,” then we could not correct the clocks in S’ (as if we do this we conclude after all that Bob sees himself ageing slower). Instead, *we would have to see these 5.34 ms – which simply correspond to a constant offset factor between two distant clocks and have nothing to do with clock rhythms – as an additional ageing of Andrew.* Regarding the example with photographs of the Minkowski diagrams, it would mean that Andrew has to advance for 5.34 ms until he crosses a Lorentzian clock from S’ displaying 16.67 ms, where he restarts being photographed again, and that this would be an additional ageing that would have nothing to do with time dilation itself, as if while Andrew advances for 5.30 ms Bob would not age at all. Whatever discourse we decide to adopt, the final result is that Andrew is older, by  $(5.34-3.01)$  ms=2.33 ms, has it had to be. Nevertheless, hopefully this simple example makes it self-evident that the addition of off-set factors does not correspond to any ageing of either twin. The plain and straight interpretation of the paradox naturally appears when synchronized clocks are used to compare time readings, and is directly related to proper times and true clock rhythms.

The above example confirms the validity of the “simultaneity argument,” in the sense that general relativity is not necessary for an explanation of the twin paradox and a correct calculation can be made within special relativity. It is not too complicated to perform all the formal calculations with Lorentzian and synchronized times from the expressions presented in section 2 and to arrive at the correct result (17), whatever the description and the twin point of view is adopted. Some of these calculations are presented in section 6.

## 5. “Solution” with Acceleration

Let the total ageing in the round trip for Andrew be  $t'_A$ , while for

Bob  $t_B''$ . Any standard calculation made “from Andrew’s point of view” is trivial using time dilation and gives the result (17),  $t_A' = t_B'' \gamma_E$ , as shown in the previous section. It is the calculation made “from Bob’s point of view” which is more difficult. As we have just seen, the correct result can be obtained, e.g., from the analysis of Minkowski’s diagrams and only within the framework of special relativity. That being so, it should be clear that acceleration itself is not an essential ingredient in the solution of the paradox, other than to evince a lack of symmetry between the twins. However, it is also possible to use general relativity to explicitly account for the effect of acceleration. We take the standard argument from Tolman as presented by Grøn [27,42].

Consider a fixed coordinate non-inertial system with Bob. There are five different moments to take into account: 1) the acceleration of Bob up to speed  $v$ ; 2) the travel up to the distant star at constant speed  $v$ ; 3) the time required to reverse the speed at the distant star; 4) the back trip at constant speed  $-v$ ; 5) the deceleration of Bob to meet Andrew. In order to keep the analysis as simple as possible, let us assume that all acceleration times in Bob’s frame are very short. According to general relativity we can replace all accelerations by gravitational fields and calculate the times processes 1) – 5) take in each frame. For the processes involving acceleration the relation of the time lapses in each frame is given by Tolman [42,27].

$$\Delta t_A' = \Delta t_B'' \left( 1 + \frac{\Delta \phi}{c^2} \right), \quad (22)$$

where the gravitational potential difference  $\Delta \phi$  is

$$\Delta \phi = g d_{AB}, \quad (23)$$

$d_{AB}$  is the distance between Andrew and Bob and  $g$  is the proper acceleration of Bob.

For processes 1) and 5) Andrew and Bob are placed in the same location, so that  $d_{AB} = 0$ . As such, there is no difference in the gravitational potential and no difference in the ageing of the twins in these periods. In fact, the irrelevance of these periods is self-evident from the start, as Bob does not even have to accelerate nor to decelerate. Bob can simply fly past Andrew already at speed  $v$  at the beginning of the trip and compare his age with the one of Andrew when they cross, and then do exactly the same on the arrival. Next, during the onward and back trips, “Bob sees Andrew ageing slower.” As thoroughly debated along this paper, this assertion refers to the comparison of the displays of successive Lorentzian clocks in Bob’s frame co-punctual with Andrew, and not to clock rhythms. Thus, for processes 2) and 4) we have

$$\Delta t' = \Delta t_L'' \frac{1}{\gamma_E}, \quad (24)$$

where  $\Delta t'$  and  $\Delta t_L''$  are respectively the proper time interval for Andrew (as it is measured with a single clock, carried by Andrew) and the Lorentzian time interval measured by different clocks in Bob’s frames in processes 2) and 4). Finally, for the turn around period 3), Bob is at a lower gravitational potential than Andrew, so that Bob’s clock runs slower according to (22), with  $d_{AB} = d$ , the distance between Andrew and the distant star in Andrew’s frame. According to the standard argument, Bob sees Andrew moving freely upwards in a gravitational field until he stops, when he is at a distance  $d$  and then begins to fall down [27]. Denoting by  $\Delta t_3''$  the short duration of the acceleration leading to the velocity reversal in Bob’s frame, and noting that the proper acceleration is given by  $g = 2v / \Delta t_3''$ , we have

$$\Delta t_3' = \Delta t_3'' \left( 1 + \frac{gd}{c^2} \right) = \frac{2v}{g} + \frac{2dv}{c^2}. \quad (25)$$

Therefore, in the limit of short acceleration periods ( $\Delta t_3'' \rightarrow 0$ ,

$g \rightarrow \infty$ ),

$$\Delta t_3' = \frac{2dv}{c^2}. \quad (26)$$

Grøn [27] considers that the time interval  $\Delta t_L''$  in (24) corresponds to the actual ageing of Bob during processes 2) and 4), consistently with the assertion from the standard interpretation of special relativity that “for the entire outward and return parts of the trip, Bob does observe Andrews clock running slow.” Hence, it is claimed that Bob should calculate the time that passed for Andrew as

$$t_A' = \Delta t' + \Delta t_3' = t_B'' \frac{1}{\gamma_E} + \Delta t_3'. \quad (27)$$

This equation is the same as (21) and the additional factor (26) is precisely the same as the factor (20). Thus, the correct final result (17) directly follows and has now been obtained in an alternative way.

The standard interpretation of general relativity claims that while for Bob a very short period passes during the turn around period, during this time he sees Andrew ageing significantly because he is subject to an intense acceleration. Grøn [27] considers that special relativity leads to a paradox, as he assigns to the Lorentzian time interval in (24) the physical meaning of proper time. As such, he does not find a justification to the inclusion of the additional factor (26) within the context of special relativity and advocates a “general relativity solution.” He then presents two possible “solutions” to the question, an “elementary solution” and another “using relativistic Lagrangian dynamics.” He acknowledges the difficulty of the former, as “it involves an assumption that is not obvious when the travelling twin stipulates the distance of his brother.” Although this is incorrect, it exposes the difficulties with the speech surrounding the standard interpretation of special relativity. A deeper analysis of this question will be given separately in a future publication.

Equation (27) contains a correct final result. But its customary interpretation is exceedingly misleading. For a start, it confuses proper times and Lorentzian clock time readings in expression (24). Then it misinterprets as well equation (22). In reality, the general relativity approach merely hides a procedure to compensate for the offset factors (3) associated with the Lorentzian clocks, as explicitly proven in the next section. Therefore, it does not bring any new physics and the obtained effect during the turn-around period described by (26) is *not* related to differences in proper times.

To finish this section, let us still note that for the discussion of the paradox the acceleration in the turn-around point can be eliminated in the same way as the accelerations in processes 1) and 5): a third twin can cross Bob at the turn around point, take note of Bob’s age when they cross, and then compare his own ageing since the meeting with Bob until he meets Andrew. He will age exactly the same as Bob in the limit Bob makes an instantaneous change in reference frame. This reinforces the indication that it is not the acceleration itself which is important, only the fact that there is one more inertial frame involved, with the corresponding modification of the clock rhythms.

## 6. General Solution

In this section we show that the broader view of special relativity proposed here demystifies the standard interpretations of the paradox both within the frameworks of special and of general relativity. It unveils, without ambiguity, that the additional factor (20) or (26) required to solve the twin paradox in either approach does *not*

correspond to any physical effect on the ageing of either of the twins. It is simply a correction of the offset factors (3) resulting from Einstein's procedure of "synchronization," which are different for the two inertial frames associated with Bob in the roundtrip. This additional factor has no deeper meaning than to say that a person who suddenly advances his own clock by five minutes must then subtract five minutes to the time reading of his clock in order to know for how long he has been waiting. Let us then perform the calculations relevant to the twin paradox using the general expressions from section (2.2) and from the point of view of each of the twins. For simplicity, we again assume that Bob travels along the  $x'$  axis with *Einstein speed* of absolute value  $v_E$  in relation to Andrew, both in the outward and on the return trips.

The easiest calculation is the simple integration of equation (15), corresponding to using the standard expression "from Andrew's point of view." Keeping the same notation as in section 2.2, the actual ageing of the twins, given by the lapses in their proper times, is just

$$\tau_B'' = \oint d\tau'' = \oint \frac{dt_L'}{\gamma_E} = \frac{1}{\gamma_E} \oint dt_L' = \frac{\tau_A'}{\gamma_E}, \quad (28)$$

where, notwithstanding (13),  $\oint dt_L' = \Delta t_L' = \tau_A'$ . Indeed, although  $dt_L'$  is measured with different clocks from  $S'$ ,  $\Delta t_L'$  in the roundtrip is measured with a single clock.

A second calculation can be made using again the standard relation, but "from Bob's point of view." This corresponds to the use of equation (15), interchanging the primed and double primed variables. In what follows the subscripts "+" and "-" refer to the outward and return trips, respectively. Note that in the onward and the return trips Bob uses different inertial frames. We have, successively,

$$\begin{aligned} \tau_A' &= \oint \frac{dt_L''}{\gamma_E} \\ &= \int_+ \frac{dt_L''}{\gamma_E} + \int_- \frac{dt_L''}{\gamma_E} \\ &= \frac{1}{\gamma_E} (\Delta t_{L,+}'' + \Delta t_{L,-}'') \\ &= \frac{1}{\gamma_E} \left( \frac{\Delta x_+''}{v_E''} + \frac{\Delta x_-''}{-v_E''} \right) \\ &= \frac{1}{\gamma_E} \left( \frac{\Delta x_+''}{v_{w+}''} \frac{v_{w+}''}{v_E''} + \frac{\Delta x_-''}{v_{w-}''} \frac{v_{w-}''}{-v_E''} \right) \\ &= \frac{1}{\gamma_E} \left( \tau_+'' \frac{v_{w+}''}{v_E''} + \tau_-'' \frac{v_{w-}''}{-v_E''} \right) \end{aligned} \quad (29)$$

where we have used (6). At this point we can show that Bob's proper time lapses are the same on the onward and the return trips,  $\tau_+'' = \tau_-''$ . Take note that Bob's ageing rhythm is different in each of the trips, except in the peculiar case in which Andrew is in the rest system. However, although the Einstein speed is the same the (synchronized) speed is not, which compensates for the difference in the rhythms. From (15),

$$\tau_+'' = \frac{\Delta t_{L,+}'}{\gamma_E} = \frac{1}{\gamma_E} \frac{\Delta x'}{w_E'} \quad (30)$$

and

$$\tau_-'' = \frac{\Delta t_{L,-}'}{\gamma_E} = \frac{1}{\gamma_E} \frac{\Delta x'}{w_E'} = \tau_+'' \quad (31)$$

Since,  $\tau_B'' = \tau_+'' + \tau_-''$ , we have

$$\tau_+'' = \tau_-'' = \frac{1}{2} \tau_B'' \quad (32)$$

We now need to eliminate the speeds of Andrew in Bob's frame,  $v_{w+}''$  and  $v_{w-}''$ , from (29). To do so, we note that, solving (8) for  $w$ ,

$$w = \frac{w_E' + v}{1 + \frac{vw_E'}{c^2}} \quad (33)$$

In turn, exchanging the roles of  $v$  and  $w$  and the primed and double primed variables in (7),

$$v_E'' = \frac{v_{w+}''}{1 - \frac{wv_{w+}''}{c^2}} \quad (34)$$

Inverting for  $v_{w+}''$  and noting that  $v_E'' = -w_E'$ , we get

$$v_{w+}'' = \frac{v_E''}{1 + \frac{vw_E'}{c^2}} = -\frac{w_E'}{1 - \frac{wv_{w+}''}{c^2}} \quad (35)$$

Substituting (33) in this last expression,

$$v_{w+}'' = -w_E' \frac{1 + \frac{vw_E'}{c^2}}{1 - \frac{w_E'^2}{c^2}} = v_E'' \frac{1 + \frac{vw_E'}{c^2}}{1 - \frac{w_E'^2}{c^2}} \quad (36)$$

Hence,

$$\frac{v_{w+}''}{v_E''} = \frac{1 + \frac{vw_E'}{c^2}}{1 - \frac{w_E'^2}{c^2}} \quad (37)$$

Analogously, for the return trip we simply have to replace  $v_E''$  by  $-v_E''$ , so that

$$\frac{v_{w-}''}{-v_E''} = \frac{1 - \frac{vw_E'}{c^2}}{1 - \frac{w_E'^2}{c^2}} \quad (38)$$

Finally, substituting (32), (37) and (38) in (29),

$$\begin{aligned} \tau_A' &= \frac{1}{\gamma_E} \frac{\tau_B''}{2} \left( \frac{1 + \frac{vw_E'}{c^2}}{1 - \frac{w_E'^2}{c^2}} + \frac{1 - \frac{vw_E'}{c^2}}{1 - \frac{w_E'^2}{c^2}} \right) \\ &= \gamma_E \tau_B'' \end{aligned} \quad (39)$$

so that we get once more the relation between the proper times of the twins (28). This somewhat dry but illuminating exercise proves that the symmetric relations

$$\tau_B'' = \oint (1/\gamma_E) dt_L' \quad (40)$$

and

$$\tau_A' = \oint (1/\gamma_E) dt_L'' \quad (41)$$

are both valid and lead directly to the relation of rhythms,

$$\tau_B'' = \tau_A' / \gamma_E, \quad (42)$$

with no need to include any additional factor. It further demonstrates that when it is said that "for the entire outward and return parts of the trip,  $B$  does observe  $A$ 's clock running slow"[39] in the traditional symmetric account of time dilation, the expression "observing the



clock running slow” is simply the comparison of the Lorentzian times of different clocks in  $B$ 's frame and has nothing to do with true clock rhythms. If it would correspond to the true ageing of the twins the paradox would remain, as it could not be simultaneously true that  $\tau_B'' = \gamma \tau_A'$  and  $\tau_A' = \gamma \tau_B''$ . We would then have to add the factors (20) or (26) to Andrew's proper time during Bob's turn-around period when the calculation is made from Bob's point of view, in order to compensate for the supposedly slower ageing of Andrew in the remaining of the trip. Incidentally, one should note that in the standard interpretation of relativity Bob needs to invoke an extremely fast ageing of Andrew, who can be light years apart, taking place exclusively during his own nearly instantaneous turn-around period. This additional factor has a much simpler nature, namely the offset factor (3), as we show next.

For simplicity, and without loss of generality, let us assume that Andrew remains in the rest system,  $S$ . Bob travels to the distant star and returns, with speed of absolute value  $v$  both in the onward and on the return trips. In this case, the absolute value of Bob's Einstein speed in  $S$  coincides with the synchronized speed,  $v_E = v$  and when Bob reaches the distant star the proper time elapsed for each of the twins is half of the total time elapsed in the complete round-trip. The ageing of Andrew until Bob arrives to the distant star can be calculated, e.g., as in (29)

$$d\tau_A = \frac{dt'_{L,+}}{\gamma_v} \quad (43)$$

Here,  $dt'_{L,+}$  refers to Lorentzian times displayed by different clocks in Bob's frame,  $S'$ , co-punctual with Andrew, along the onward trip. Integrating along the onward trip,

$$\frac{1}{2} \tau_A = \frac{\Delta t'_{L,+}}{\gamma_v} \quad (44)$$

Now, instead of proceeding as before, we can relate the Lorentzian time with the synchronized time and the proper time by (3) and (13), from where

$$\Delta t'_{L,+} = \frac{1}{2} \tau_B' + \frac{v}{c^2} d' \quad (45)$$

Note that  $x'$  in (3) is negative and  $d' = |x'|$  is the position of Andrew in Bob's frame when he arrives to the star, i.e.,  $d'/\gamma_v = d$ . Substituting in (44),

$$\begin{aligned} \frac{1}{2} \tau_A &= \left( \frac{1}{2} \tau_B' + \frac{v}{c^2} d' \right) \frac{1}{\gamma_v} \\ &= \frac{1}{2} \frac{\tau_B'}{\gamma_v} + \frac{v}{c^2} d \end{aligned} \quad (46)$$

Finally, noting that  $d = v\tau_A/2$  have, successively,

$$\begin{aligned} \frac{1}{2} \tau_A &= \frac{1}{2} \frac{\tau_B'}{\gamma_v} + \frac{1}{2} \tau_A \frac{v^2}{c^2} \\ \tau_A \left( 1 - \frac{v^2}{c^2} \right) &= \frac{\tau_B'}{\gamma_v} \\ \tau_A &= \gamma_v \tau_B' \end{aligned} \quad (47)$$

which is again the true relation between the ageing of the twins.

The major and very consequential remark is to note that the second term in the r.h.s. of (46) is the additional factor (20) or (26). This expression is valid all along the one-way outward trip, in which Bob remains in the same inertial frame. It relates the proper times of both twins along the onward trip. It comes without saying that the

additional factor has *nothing* to do neither with any strangeness at the turn-around point related to the “simultaneity argument” nor with acceleration, general relativity and the equivalence principle.

To summarize, for the particular case just considered the following relations are all valid all along the one-way onward trip,

$$\begin{aligned} d\tau_A &= \frac{dt'_{L,+}}{\gamma_v} \\ &= \frac{d\tau_B'}{\gamma_v} + \frac{v}{c^2} dx \\ &= \gamma_v d\tau_B' \end{aligned} \quad (48)$$

and express the fact that Bob is ageing slower in this case. Understanding these three equalities clarifies the language and completely solves the twin paradox.

A straightforward analysis made with proper times and the IST transformation removes all difficulties surrounding the paradox at the onset. In particular, it is no longer said that each twin sees the other ageing slower during the length of the one-way trips. Instead, both twins agree that in each part of the trip one of them is ageing slower while the other is ageing faster (or are both ageing at the same rhythm). The simple addition of the proper times of the onward and back trips then gives directly the final result for both twins. There is no need to correct for any strangeness on the turn around period nor for any supposed effect of acceleration and equivalent gravitational fields. Specifically, there is no need to add any factor to Andrew's age due to a change in inertial frame of Bob, whatever point of view is considered.

As a final remark, the general calculations can be made as well using (16) as a starting point, both as it is written and as well interchanging the roles of  $v$  and  $w$  and the primed and double primed variables, in an additional consistency check of the present formalism.

## 7. Conclusion

In previous papers we have proposed a general formulation of special relativity, where the postulates are formulated in a weaker form than in the traditional presentation, while keeping fully compatible with all experimental evidence [1,2,3,4,5,6]. The theory assumes the existence of (at least) one reference frame where the one-way speed of light in vacuum is isotropic and equal to  $c$ , denoted as the “rest system.” It was shown that the theory is undetermined, unless the one-way speed of light in one reference frame is measured [4]. The somewhat evident but very important difference between “time readings of Lorentzian clocks” and proper times or clock rhythms is thoroughly discussed in I [6]. It is noted that although the *description* of time dilation made with Lorentzian clocks is symmetrical for two inertial observers in relative motion, as it is well-known from the standard interpretation of special relativity, the reciprocal relation does not relate the clock *rhythms* [6]. Actually, the time dilation relation between the clock rhythms of two inertial observers in relative motion is *not* symmetric. During each one-way trip one of the observers is actually ageing slower than the other, or it may even happen that both are ageing at the same rhythm. Furthermore, without reference to the rest system it is impossible to know which of them is actually ageing slower, in another way of stating the indeterminacy of special relativity [6].

Within this context, in this work we have discussed the twin paradox in detail, as an illustration of the power and simplicity of the general formulation of special relativity formerly presented. In a round-trip such as the one in the classical configuration of the twin

paradox, it is the returning twin who is younger when both twin meets, despite the impossibility of knowing which of the twins is younger at each phase of the trip. The result is due to the cumulative effect of the clock rhythms along the complete journey.

The total ageing of each twin is calculated directly from the *sum of their proper times* on the onward and the return trips. Contrary to what happens in the standard interpretation of special relativity and of general relativity, there is no need to consider any additional ageing factor of the resting twin as seen by the moving twin to account for the change in reference frame or the acceleration (sections 4 and 5). As a matter of fact, the factor invoked by these standard interpretations was deduced from the offset between Lorentzian and synchronized clocks (section 6). It was thus demonstrated that this factor has nothing to do with any modification with the clock rhythms and with the ageing of the twins during the turn-around period. It is merely a correction to a peculiar way of giving the initial adjustment (a so-called “synchronization”) to the clocks. As a consequence, it becomes clear that acceleration does not play any role in the twin paradox other than telling which of the twins is returning back.

Finally, we would like to underline the following. The standard interpretation of special relativity pretends to assign a physical meaning of real ageing to assertions like “during the onward trip Bob sees Andrew ageing slower, Andrew himself also sees Bob ageing slower, but the change in inertial frames corrects this symmetry, as a result of the relativity of simultaneity, and makes it in the end that Bob is younger when the twins meet.” The standard interpretation of general relativity pretends to assign a physical meaning to sentences like “during the turn around period Bob sees Andrew ageing very quickly because he sees him under the effect of a gravitational field at a higher gravitational potential.” These are erroneous interpretations of correct mathematical results. One message we want to convey regarding the twin paradox is that such discourse is no longer tolerable and should become “unspeakable:” it was proven that “seeing the other twin ageing slower” is meaningless in this context and corresponds to the symmetric description arising from the comparison of the *time readings* of Lorentzian clocks (4 and 5), whose roots lie in the indeterminacy of special relativity [4]. It does *not* correspond to the clock *rhythms* and to the ageing of the twins.

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